



USPAS June 2015

Rutgers

SRF MATERIALS OTHER THAN BULK NIOBIUM

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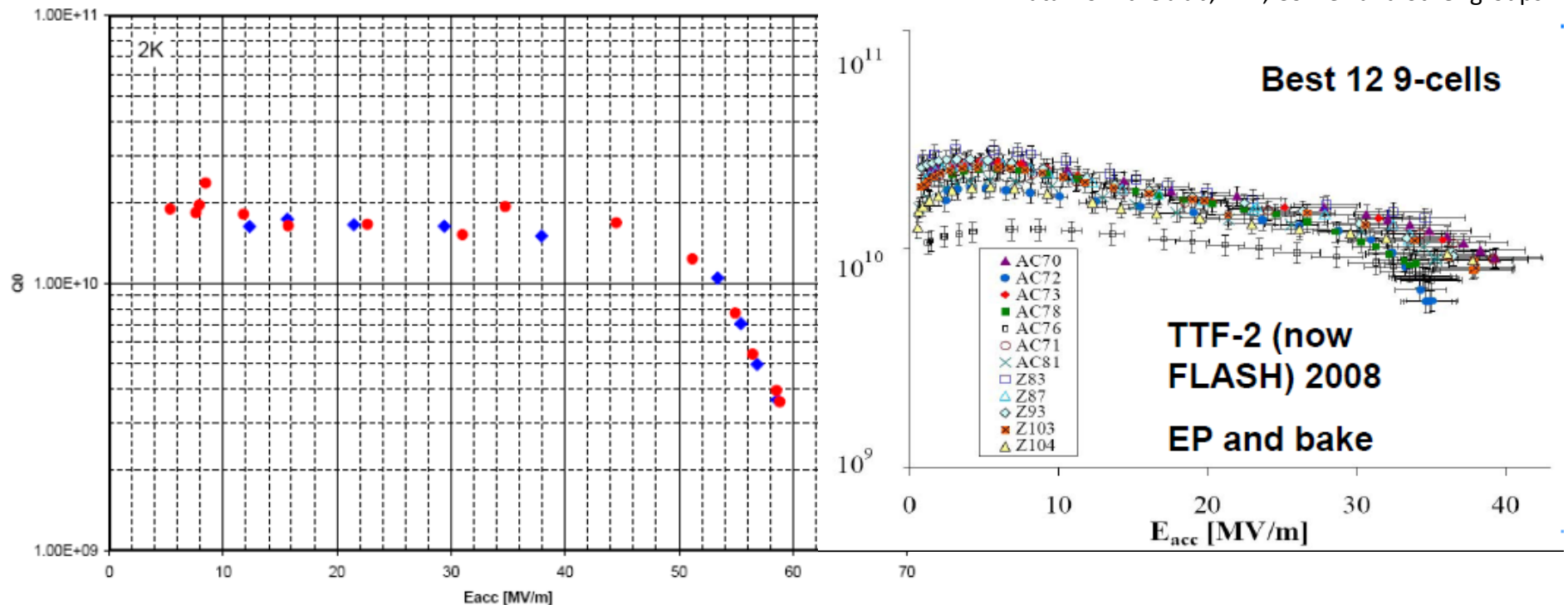
Thomas Jefferson National Accelerator Facility

Jefferson Science Associates LLC

Why looking beyond bulk Nb?

Nb has the highest critical temperature T_c ($=9.25\text{k}$) and the highest lower critical magnetic field H_{c1} ($\approx 180\text{ mT}$) of any elemental superconductor

Cornell 60 mm aperture re-entrant cavity LR1-3 March 14, 2007



Breakdown fields close to the de-pairing limit of 50 MV/m for Nb have been achieved

Best Nb cavities approaching their intrinsic limit at $H_{\max} = H_c$

For further improved cavity RF performance, innovation needed

Possibilities to use higher performance superconductors other than Nb?

Looking beyond Nb – Potential Benefits

➤ Reduced material costs

- Use of inexpensive, highly formable materials with higher thermal conductivity such as Cu or Al

➤ Simplified engineering, fabrication and assembly

- Separation of cavity structure from superconducting surface
- Maximum flexibility & largest variety of options in design of integrated cavity/cryostat structures

➤ Higher gradients

Increasing gradients reduces capital expense of cryomodules – potentially several \$100M savings & minimize conventional facilities

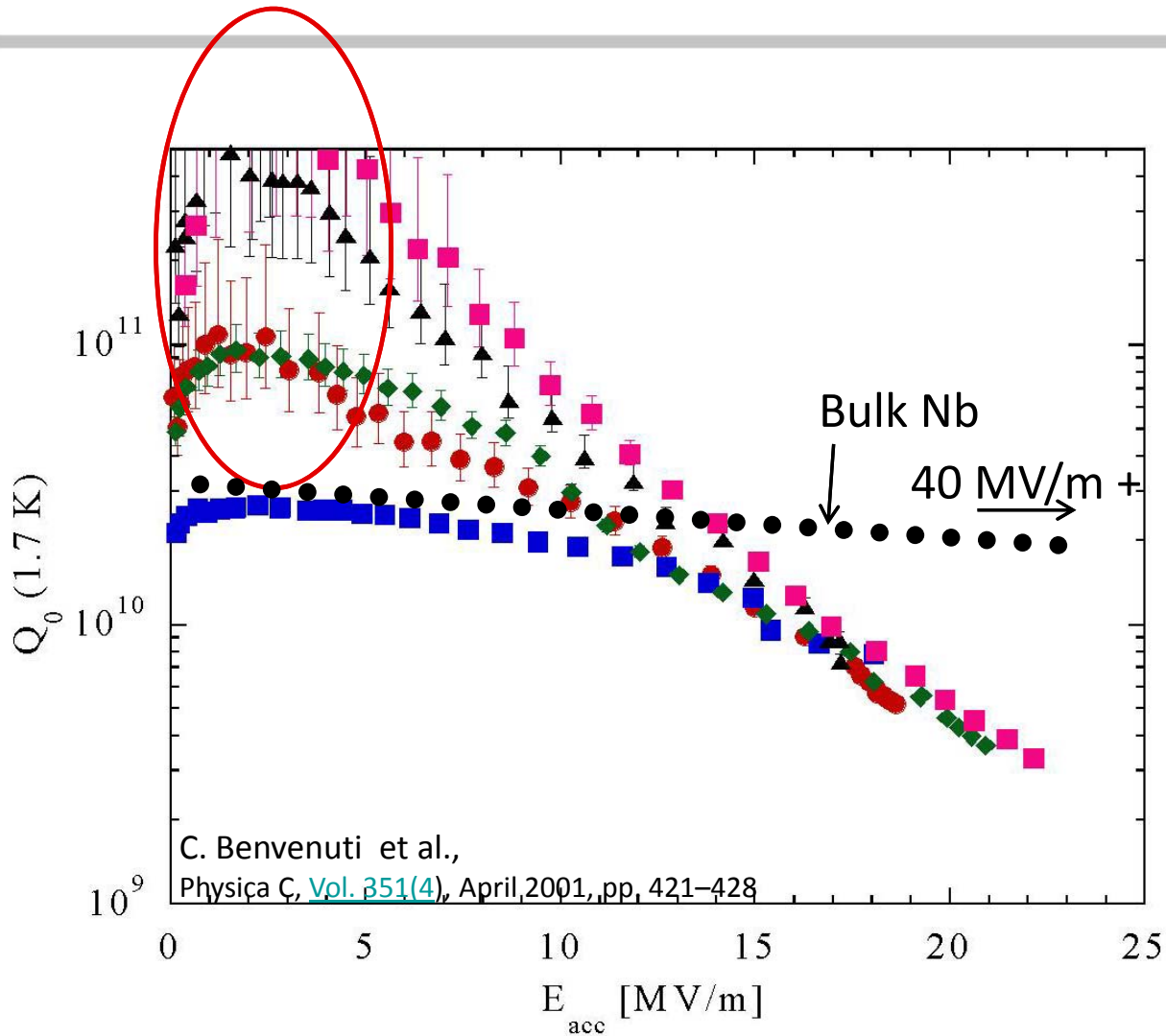
➤ Lower RF losses

Low loss (high Q) cavities reduce He costs, >\$10M potential capital savings, and several \$M/year in operating costs

➤ Potentially higher operating temperatures (>4.2K)

Nb Thin films

Thin Films for SRF - State of the Art



1.5 GHz Nb/Cu cavities, sputtered w/ Kr @ 1.7 K ($Q_0=295/R_s$)

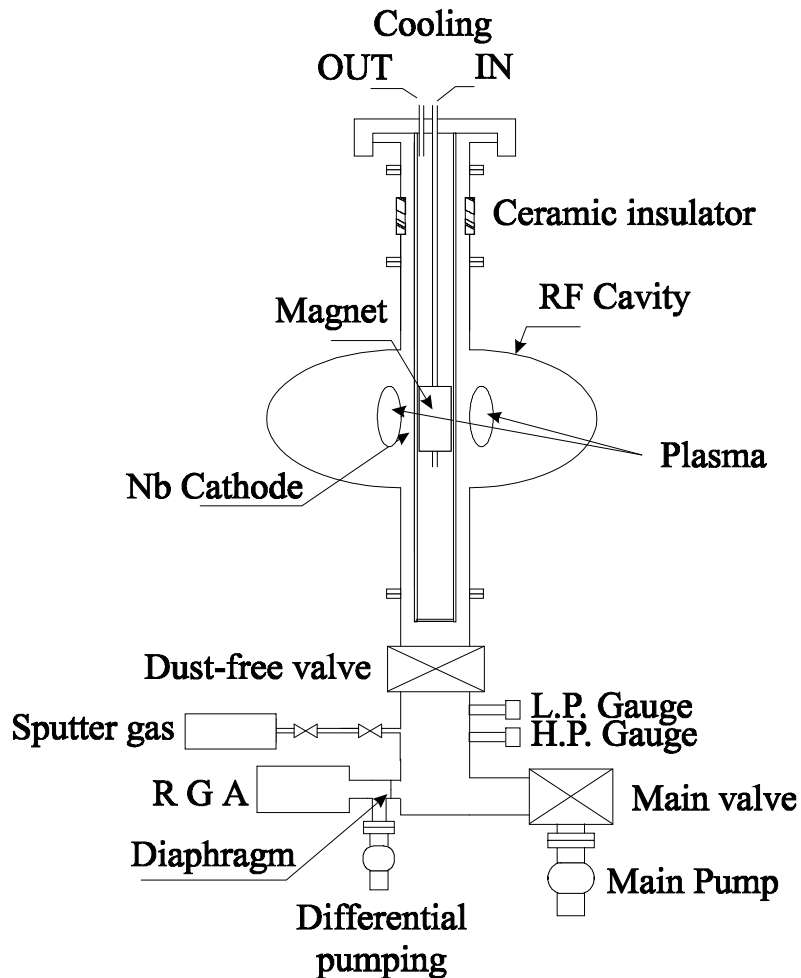
Thickness of interest for SRF applications = RF penetration depth, i.e. the very top 40 nm of the Nb surface.

High Q at low field
BUT strong Q-slope

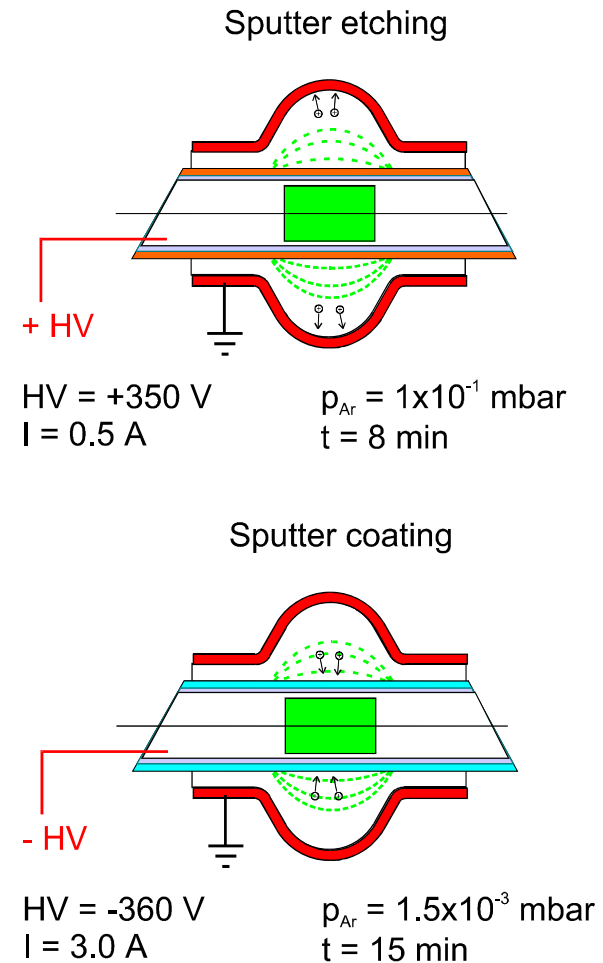
Nb/Cu Sputtered Films

Two Production Methods: standard and oxide-free coatings

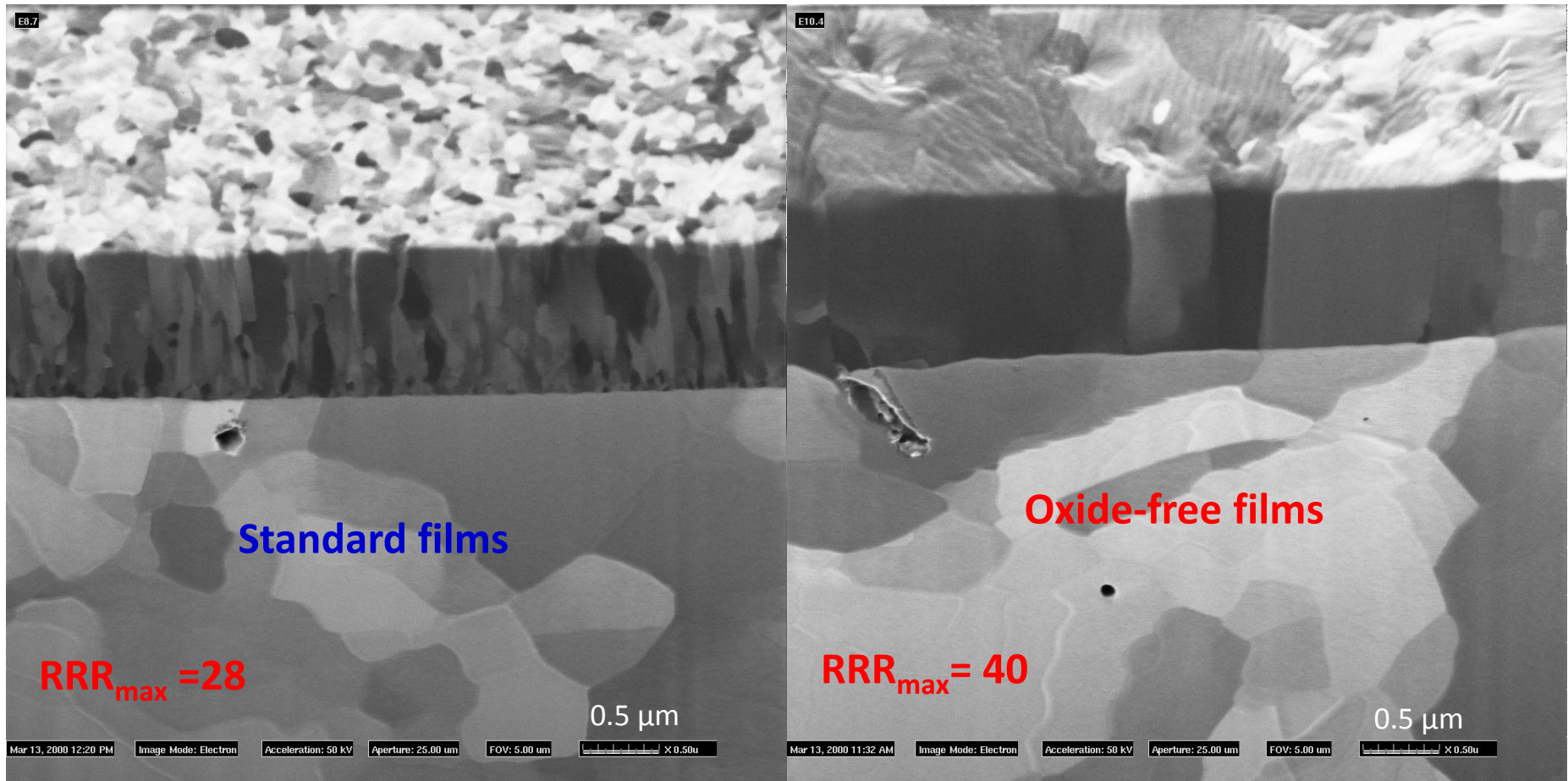
Single cathode system



Double cathode system



Nb/Cu Sputtered Films: Film structure – FIB cross sections



Columnar grains, size ~ 100 nm

In plan diffraction pattern: powder diagram
(110) fiber texture \perp substrate plane

Equi-axed grains, size $\sim 1-5\mu\text{m}$

In plan diffraction pattern: zone axis [110]

Heteroepitaxy

Nb (110) //Cu(010) , Nb (110) //Cu(111), Nb (100) //Cu(110)

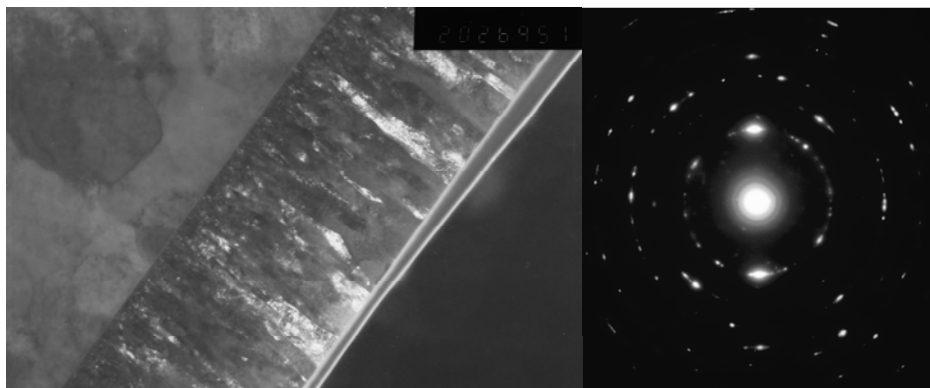
Courtesy: P. Jacob - EMPA

Nb/Cu Sputtered Films: TEM cross-section & in-plane

cross-section

Nb/Cu Oxide

Grain size $\sim 0.1 \mu\text{m}$

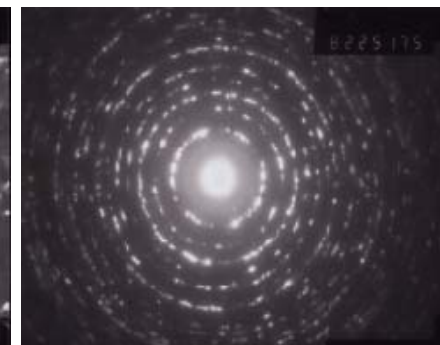
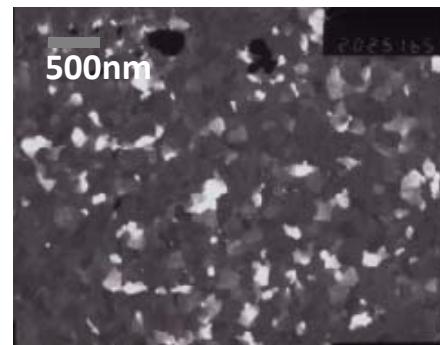


In-plane

Grain size $\sim 100 \text{ nm}$

Fiber texture

Diffraction pattern: powder diagram



Nb/Cu

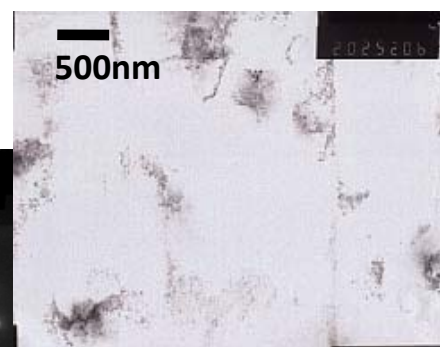
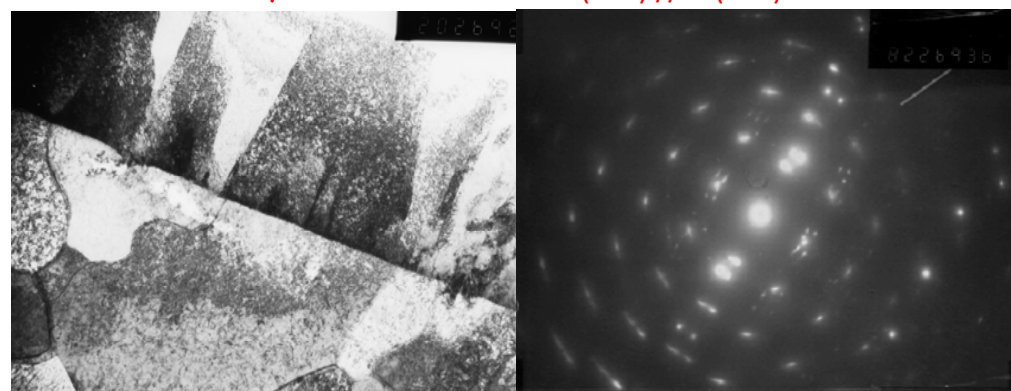
Grain size $\sim 1\text{-}5 \mu\text{m}$

Hetero-epitaxy relationships

Nb (110) // Cu(010)

Nb (110) // Cu(111)

Nb (100) // Cu(110)

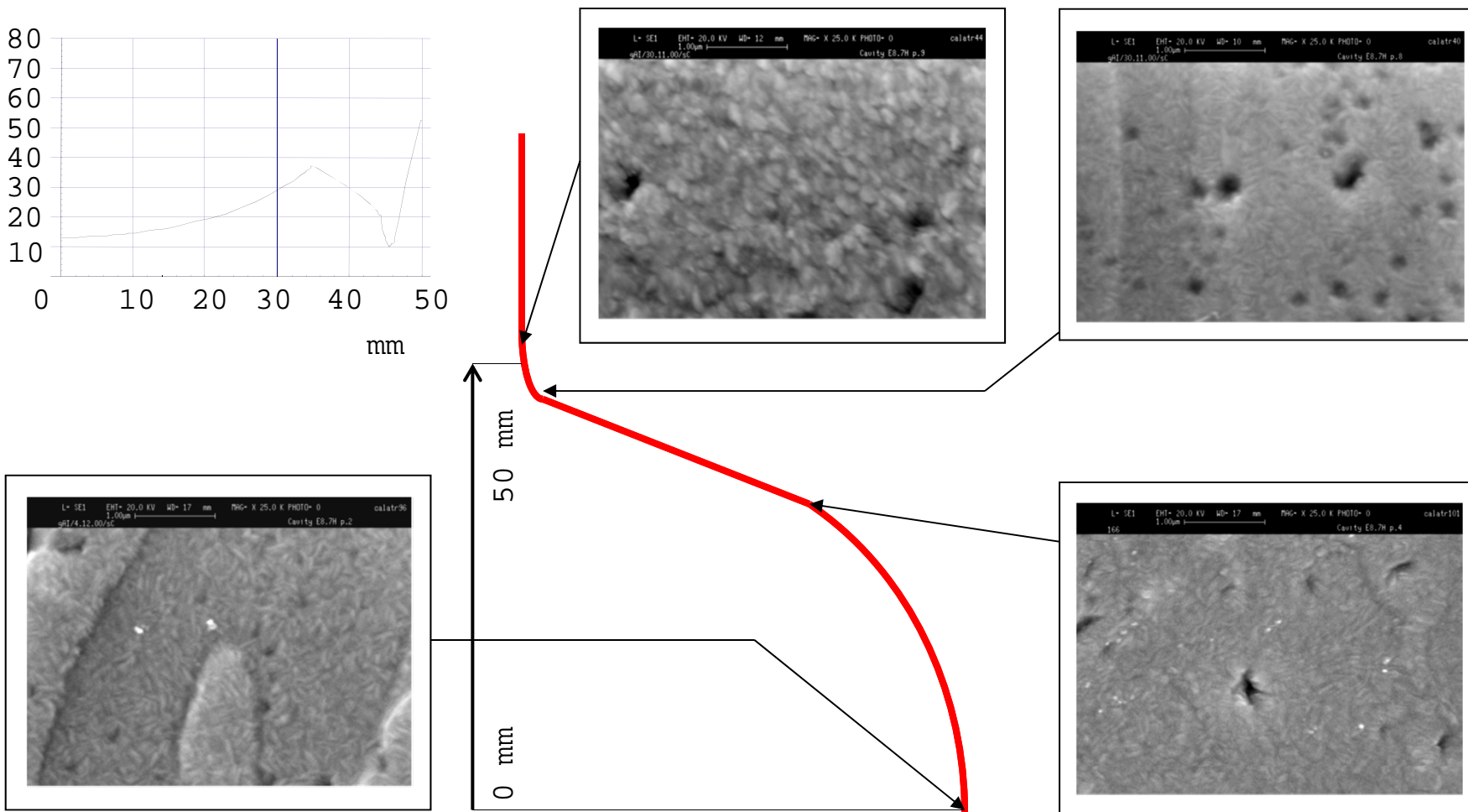
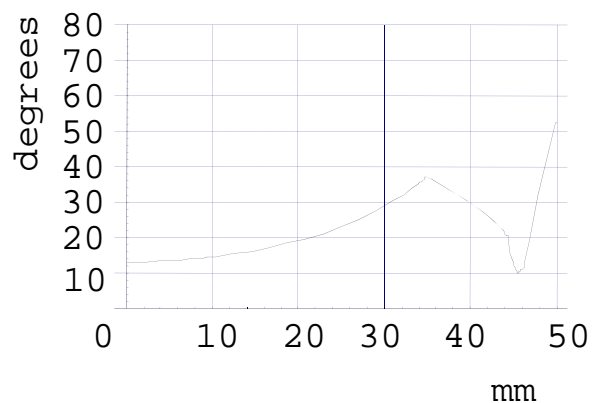


Grain size $\sim 1\text{-}5 \mu\text{m}$

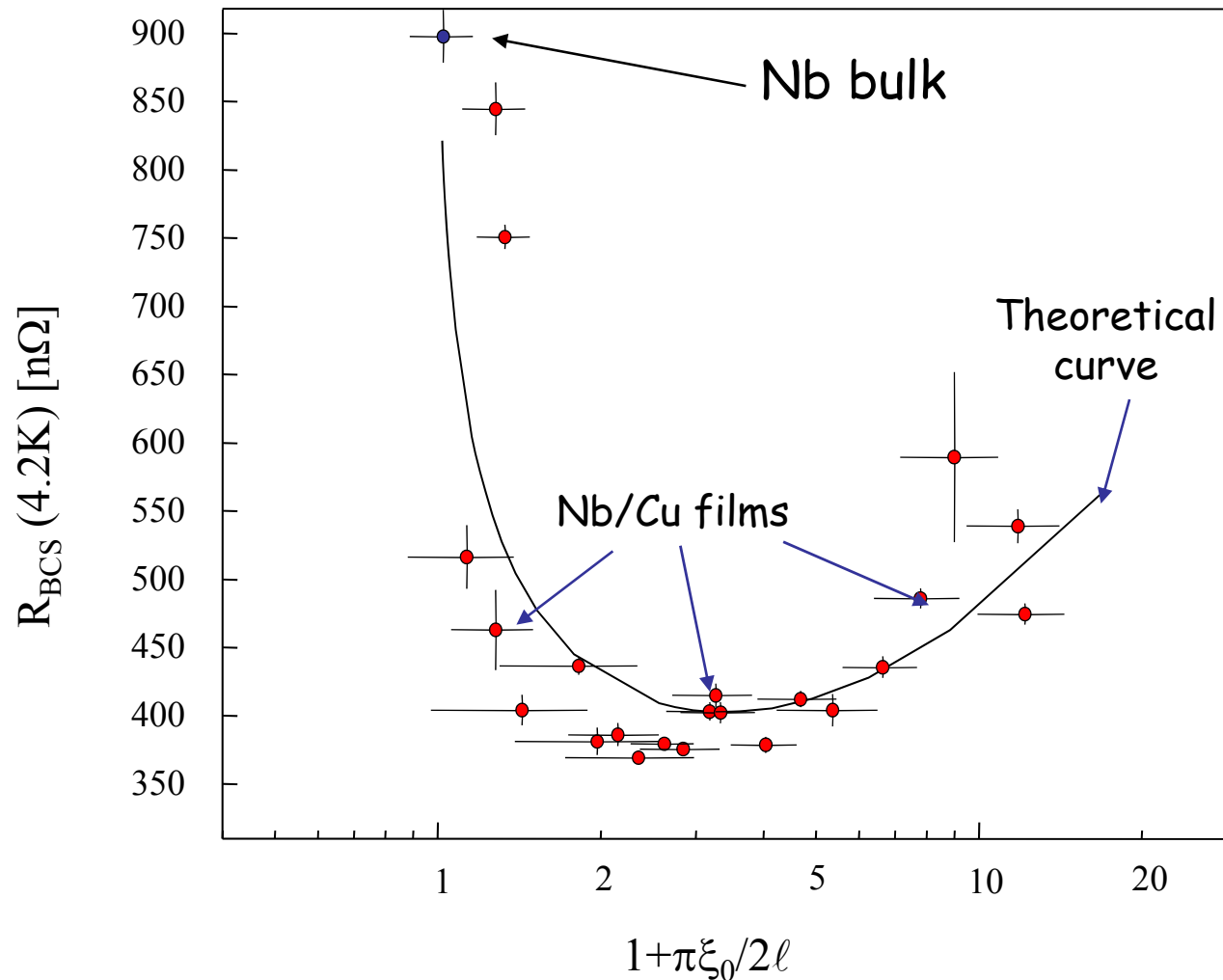
Heteroepitaxy

Diffraction pattern: zone axis [110]

Nb/Cu Sputtered Films: Intrinsic film roughness & incidence angle of the niobium atoms



Theoretical and experimental BCS resistance at zero RF field

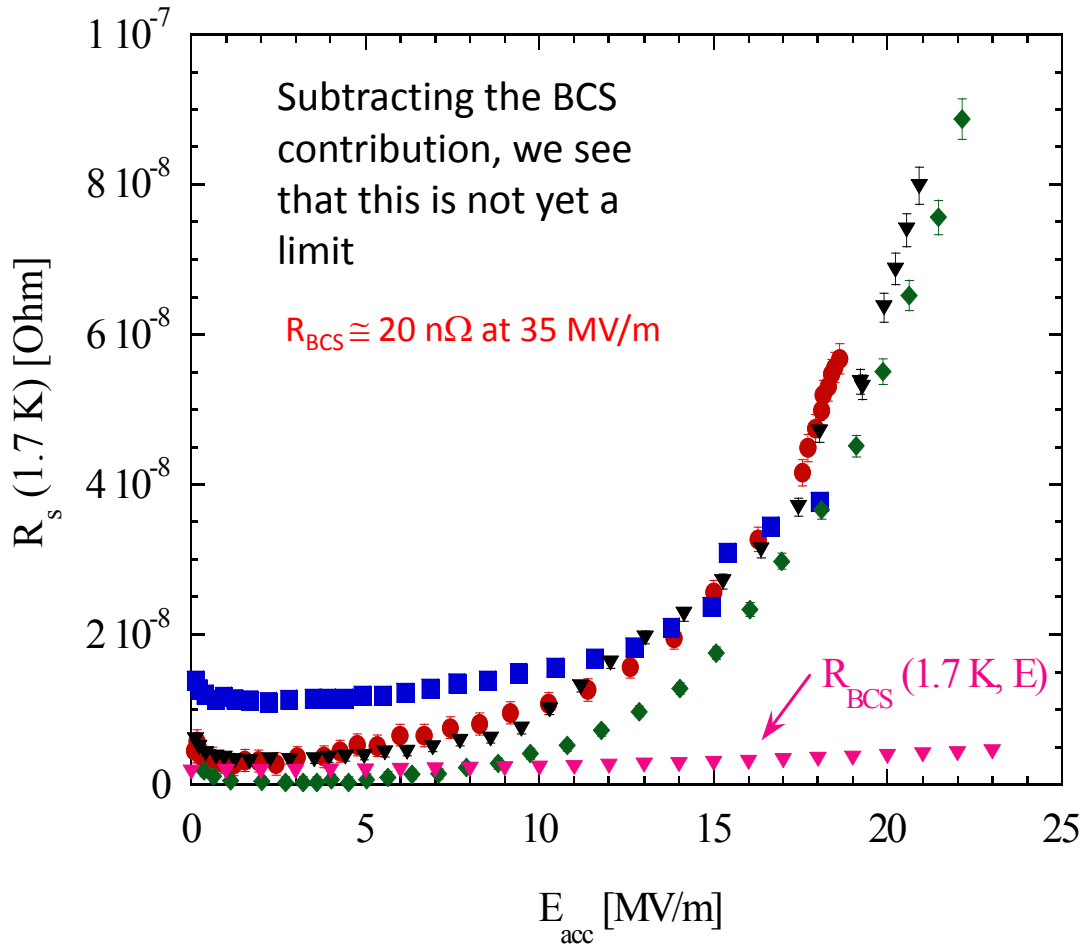


R_{BCS} at 4.2 K
Nb bulk: ~ 900 nΩ
Nb films: ~ 400 nΩ

R_{BCS} at 1.7 K
Nb bulk: ~ 2.5 nΩ
Nb films: ~ 1.5 nΩ

Compilation of results from several Nb/Cu and Nb bulk 1.5 GHz RF cavities

Best RF performance measured on 1.5GHz Nb/Cu sputtered cavities



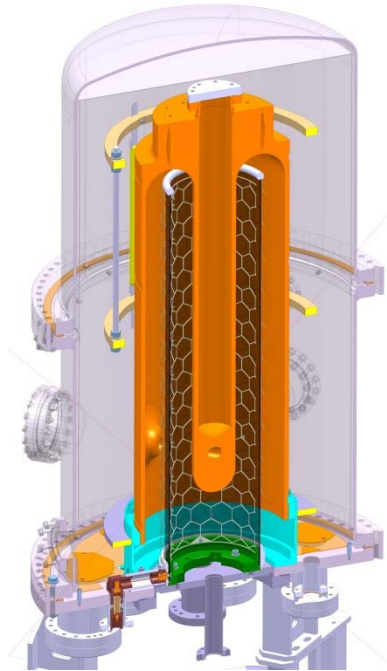
Coatings performed using krypton on electropolished spun copper cavities(Santa Fe 1999)

	Standard	Oxide-free
RRR	11.5 ± 0.1	28.9 ± 0.9
T_c	$9.51 \pm 0.01 \text{ K}$	$9.36 \pm 0.04 \text{ K}$
Ar cont.	$435 \pm 70 \text{ ppm}$	$286 \pm 43 \text{ ppm}$
Texture	(110)	(110), (211), (200)
H_{c1}	$85 \pm 3 \text{ mT}$	$31 \pm 5 \text{ mT}$
H_{c2}	$1.150 \pm 0.1 \text{ T}$	$0.73 \pm 0.05 \text{ T}$
a_0	$3.3240(10) \text{ \AA}$	$3.3184(6) \text{ \AA}$
$\Delta a_{\perp}/a_{\perp}$	$0.636 \pm 0.096 \%$	$0.466 \pm 0.093 \%$
Stress	$-706 \pm 56 \text{ MPa}$	$-565 \pm 78 \text{ MPa}$
Grain size	$110 \pm 20 \text{ nm}$	$> 1 \mu\text{m}$

Attaining even **bulk-like performance** with Nb film would enable major system simplifications.

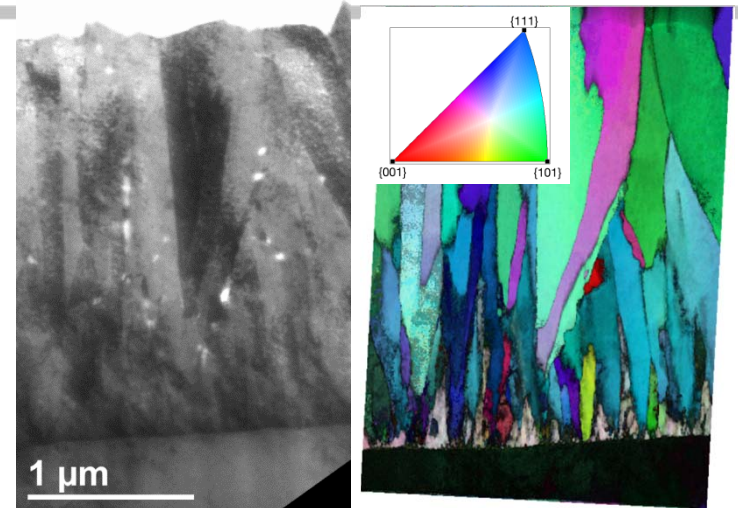
Deposited films, once a production process is developed, may offer the **highest level of quality assurance and thus reliable performance.**

QWR Resonators for HIE-Isolde Upgrade



Nb thin film sputtered
on 3D forged OFE Cu
substrate
By dc-bias diode sputtering

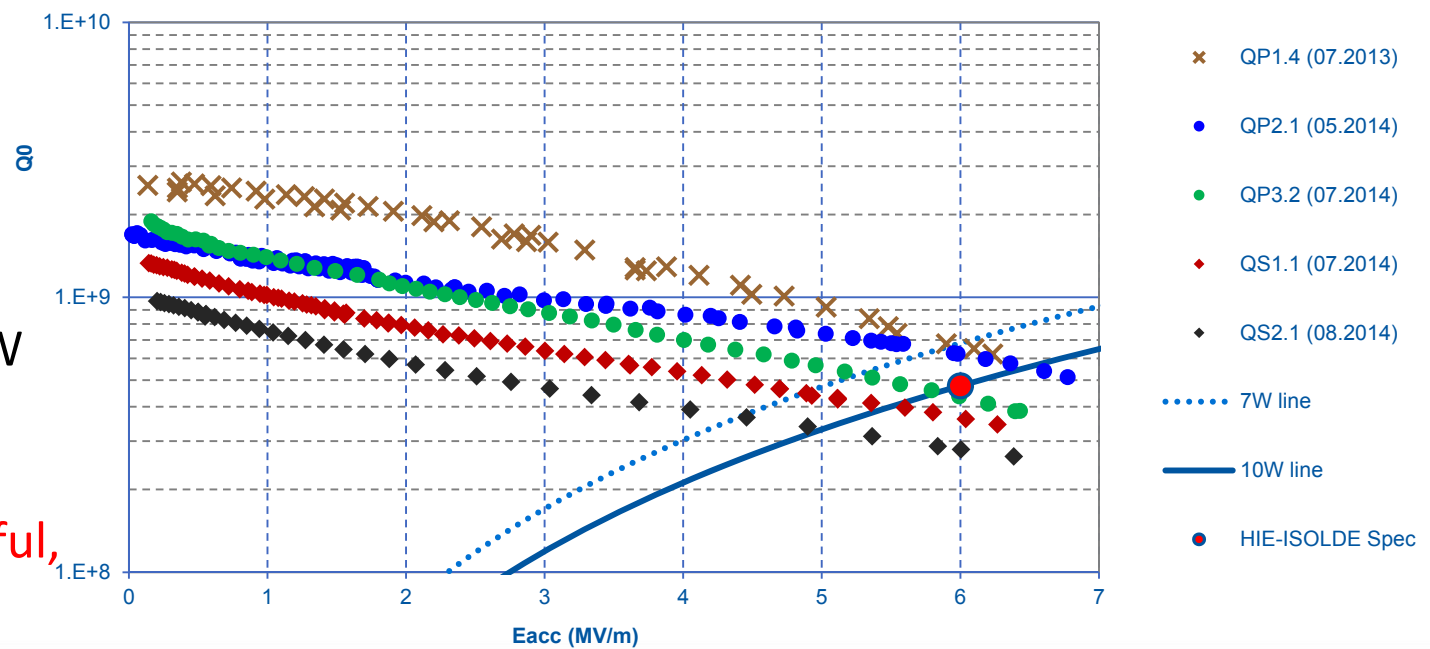
→ average RRR = 35



Specifications:

$E_{acc} = 6\text{MV/m}$
 $P_{cav} @ 6\text{MV/m} = 10\text{W}$
Average $P_{cav} \sim 12\text{W}$

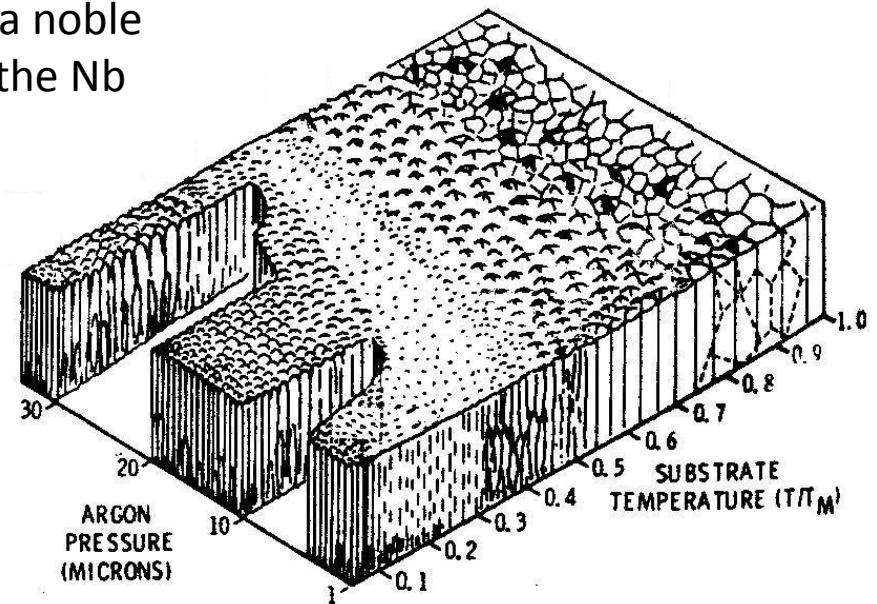
Prototyping successful,
Production started



High-energy deposition techniques

- Crystalline defects, grains connectivity and grain size may be improved with a higher substrate temperature which provides higher surface mobility (important parameter is $T_{\text{substrate}}/T_{\text{melting_of_film}}$)
- However the Cu substrate does not allow heating
- **The missing energy may be supplied by ion bombardment**
 - In bias sputter deposition a third electron accelerates the noble gas ions, removing the most loosely bound atoms from the coating, while providing additional energy for higher surface mobility
 - Other techniques allow working without a noble gas, by ionising and accelerating directly the Nb that is going to make up the coating
 - These techniques allow also to obtain “conformal” coatings that follow the surface profile better filling voids.

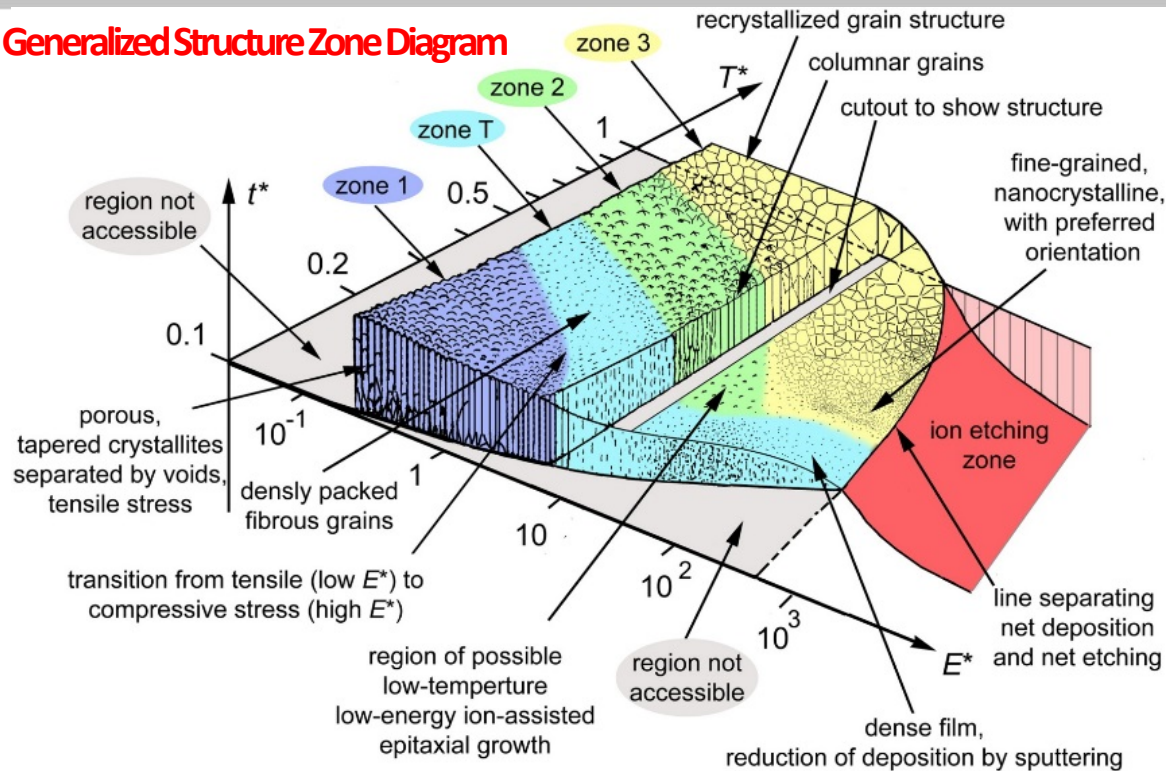
Thornton's “Structure Zone Model”



Energetic Condensation

Condensing (film-forming) species : hyper-thermal & low energies (>10 eV).

Generalized Structure Zone Diagram



derived from Thornton's diagram for sputtering (1974)

A. Anders, Thin Solid Films 518 (2010) 4087

Additional energy provided by fast particles arriving at a surface
 \Rightarrow number of surface & sub-surface processes \Rightarrow changes in the film growth process:

- residual gases desorbed from the substrate surface
- chemical bonds may be broken and defects created thus affecting nucleation processes & film adhesion
- enhanced mobility of surface atoms
- stopping of arriving ions under the surface

\Rightarrow Changes in

- morphology
- microstructure
- stress

As a result of these fundamental changes, energetic condensation allows the possibility of controlling the following film properties:

- Density of the film
- Film composition
- Crystal orientation may be controlled to give the possibility of low-temperature epitaxy

Next generation Nb films

ALL film properties are a direct consequence of the film structure, defect/impurity content... thus the technique, environment, substrate are key factors

Full control of the deposition process & tailored SRF performance

UNDERSTANDING OF

- The **chemistry** of the involved species
 - Reactivity**
 - Stoichiometric sensitivity**
 - Reaction process **temperatures**
- Structure dependence on substrate structure**
- Influence of deposition energy** on resulting structure
- Sensitivity to the presence of contaminating species, defects**
- Protection** of desired film against subsequent **degradation**

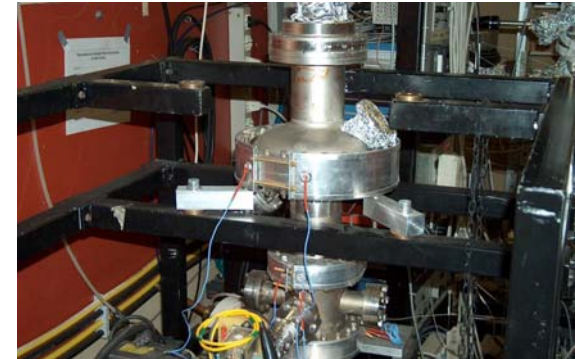
Careful **characterization of the attained composition and microstructure** (RHEED, STM, XRD, EBSD, AFM, optical profilometry, XPS, SIMS, TEM, FIB).



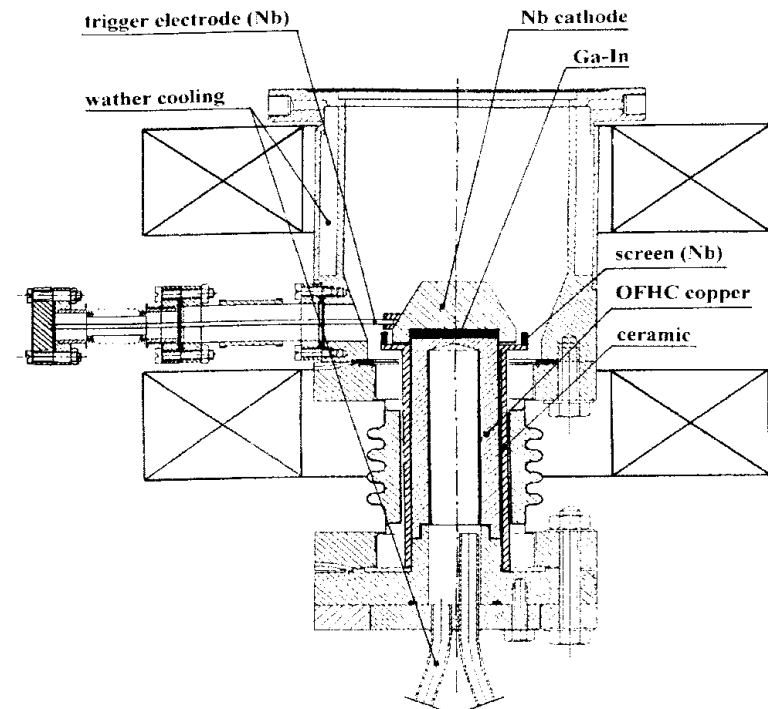
Close association with **resulting RF surface impedance & superconducting properties** (λ , Δ , T_c , H_c , RRR)

Plasma Arc (ARCO-INFN/Soltan Institute)

- In the plasma arc an electric discharge is established directly onto the Nb target, producing a plasma plume from which ions are extracted and guided onto the substrate by a bias and/or magnetic guidance
- Magnetic filtering (and/or arc pulsing) is also necessary to remove droplets

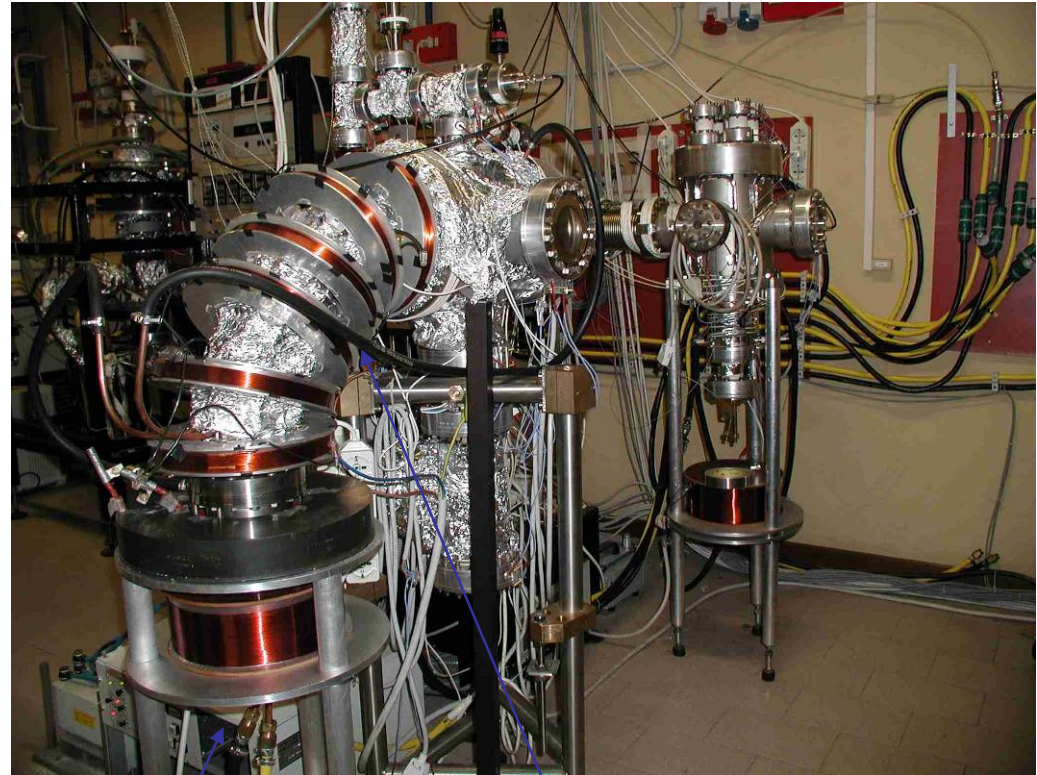
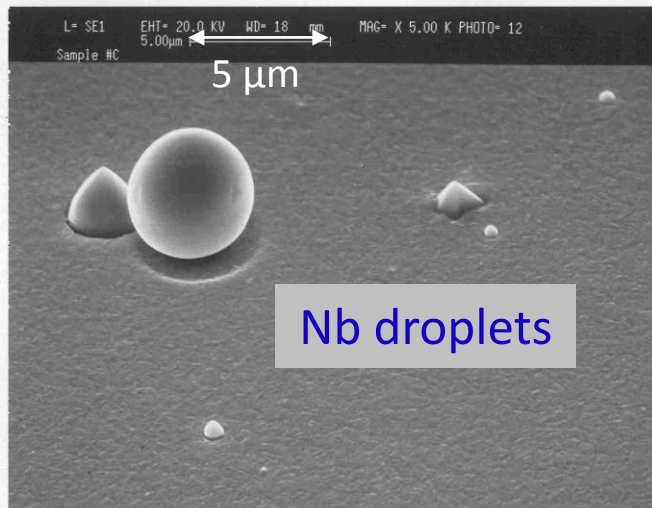
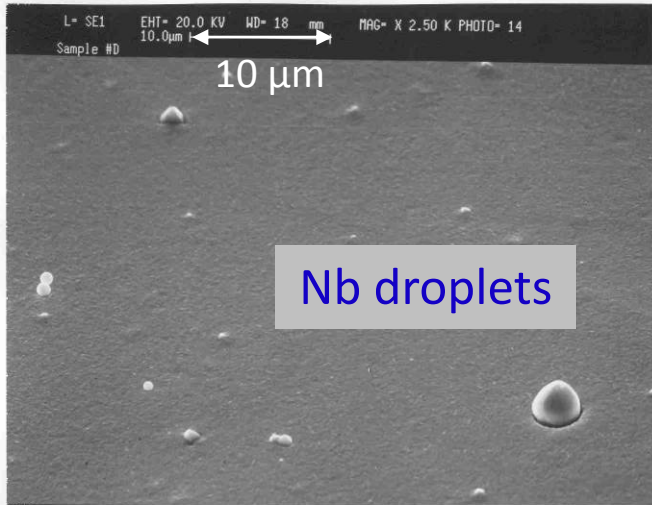


- A trigger for the arc is necessary: either a third electrode, or a laser
- Arc spot moves on the Nb cathode at about 10 m/s
- Arc current is 100-200 A
- Cathode voltage is ~ 35 V
- Ion current is 100-500 mA on the sample holder ($2-10$ mA/cm²)
- Base vacuum $\sim 10^{-10}$ mbar
- Main gas during discharge is Hydrogen ($\sim 10^{-7}$ mbar)
- Voltage bias on samples 20-100 V



From: R. Russo, A. Cianchi, S. Tazzari

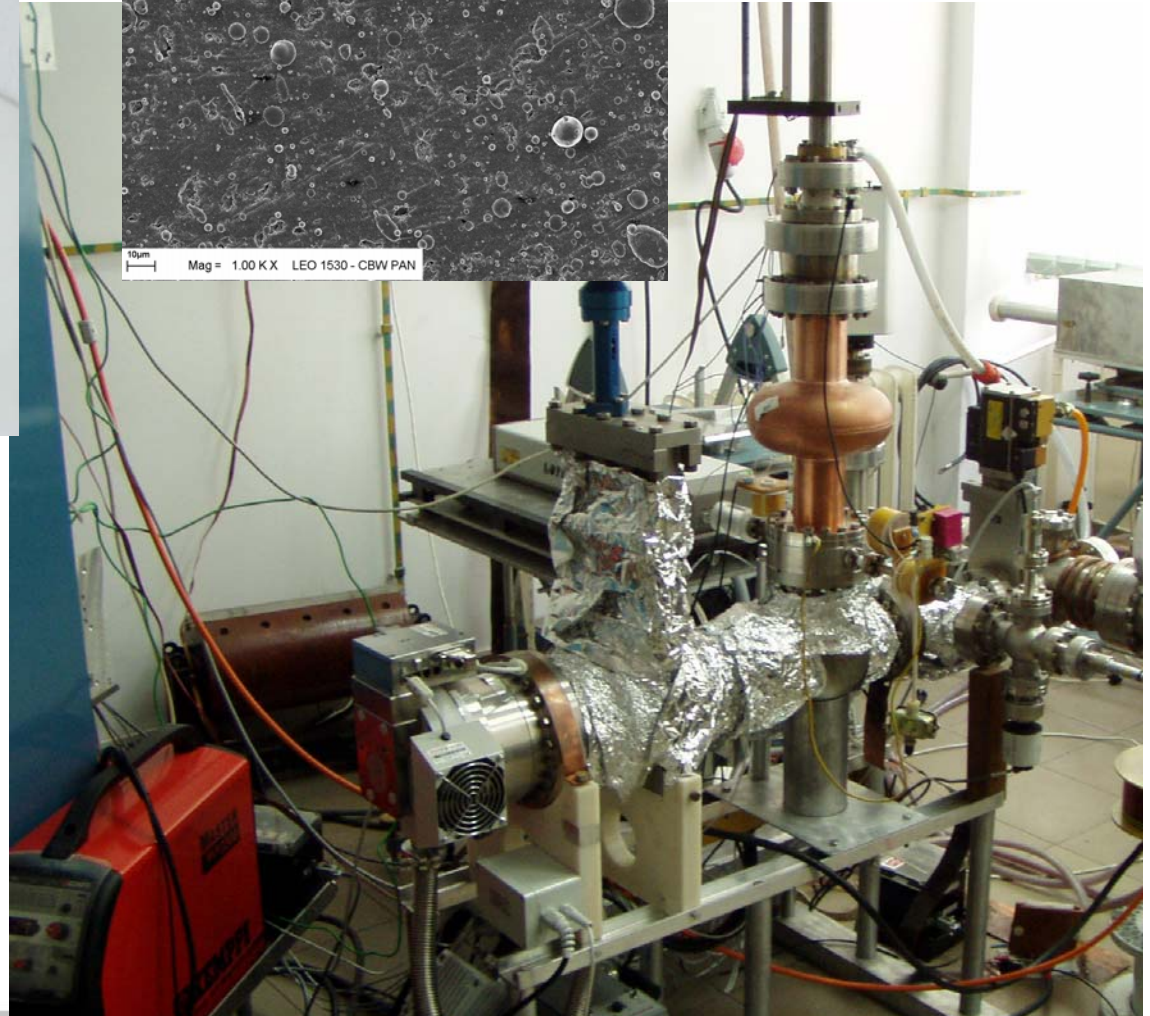
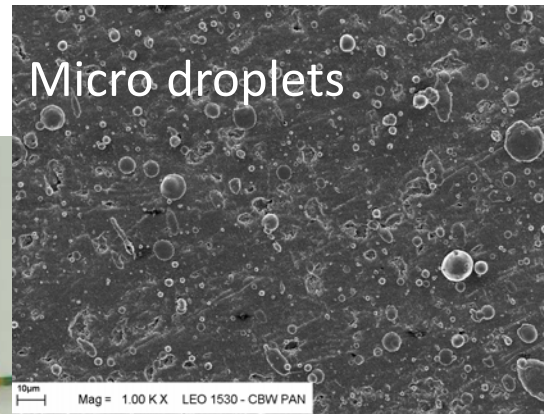
Plasma Arc – Presence of macro-particles



Arc source

Magnetic filter

Thin Film Cavity Production



Linear arc coating at
IPJ Swierk

Planar arc – RF measurements on samples

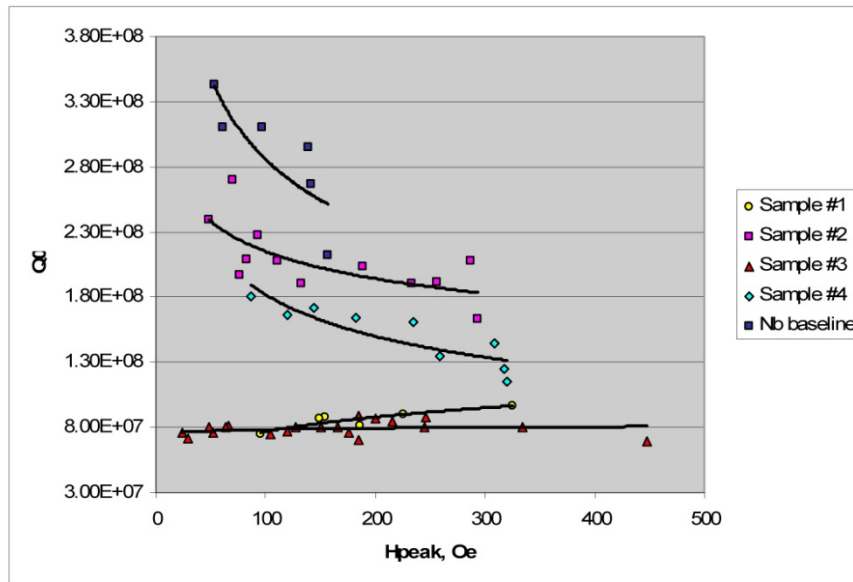


Figure 5: Q_0 versus peak magnetic field for different Nb/Cu end plates and a bulk Nb end plate.



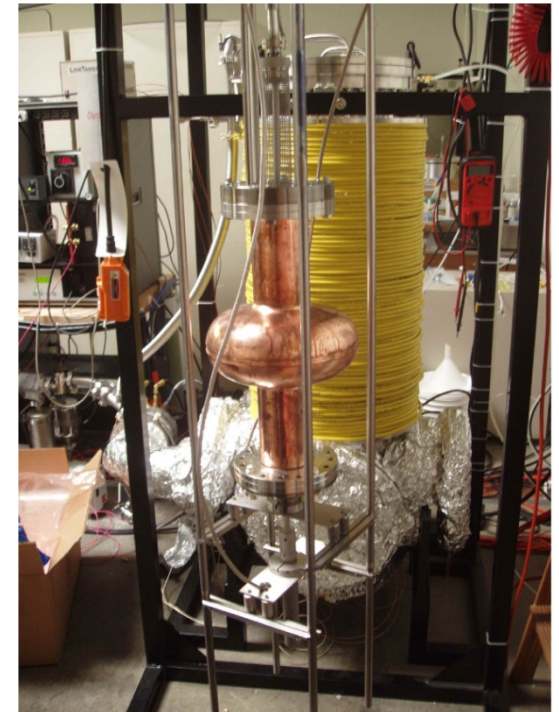
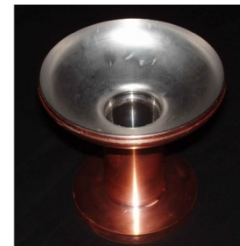
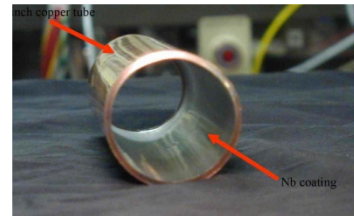
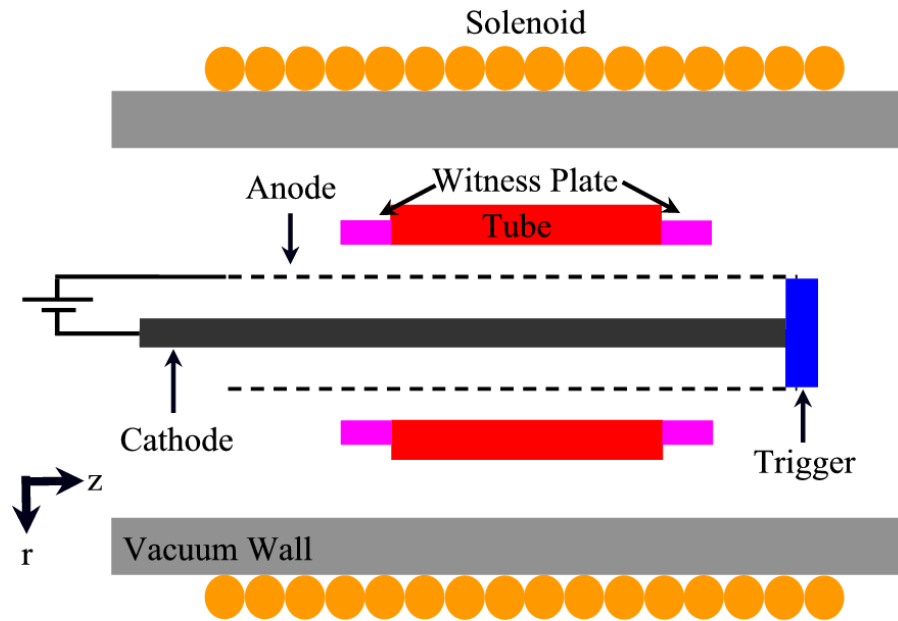
Cu samples with Nb ARC-coating. Used as a base plate of 6 GHz cavity operating in the TE₀₁₁ mode. At low field, the surface resistance is in the range 3-6 $\mu\Omega$ as compared to the BCS R_s of 0.22 $\mu\Omega$ at 2.2 K and small mean free path. The Q remained constant up to a field of 300 Oe.

A baseline of 2.2 $\mu\Omega$ is measured with this cavity with a solid Nb plate

A. Romanenko and H. Padamsee, Proc. SRF2005, Cornell, USA, 2005

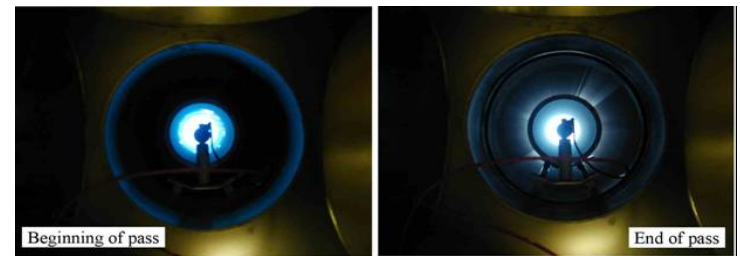
Coaxial Energetic Deposition™ (CED)

Alameda Applied Science Corporation



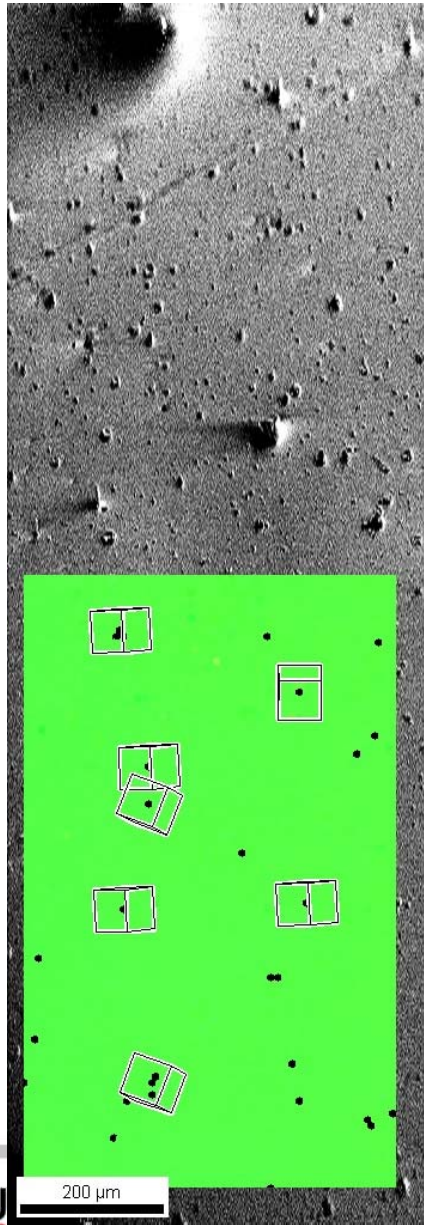
- Cathode: 60 cm conducting rod (1 cm dia).
- Anode: 45 cm Mo mesh tube (4.5 cm ID)
- Substrate: 5 cm ID minimum for this configuration
- Solenoid: **B** ranges from 0-10 mT in z^+ or z^- direction.

Note: Anode does not collect all the arc current. Mesh spacing is much larger than Debye length. Current is measured between the power supply and cathode through known resistor.



Rotating arc moves down the axis of a 4" ϕ tube

CED Films structure



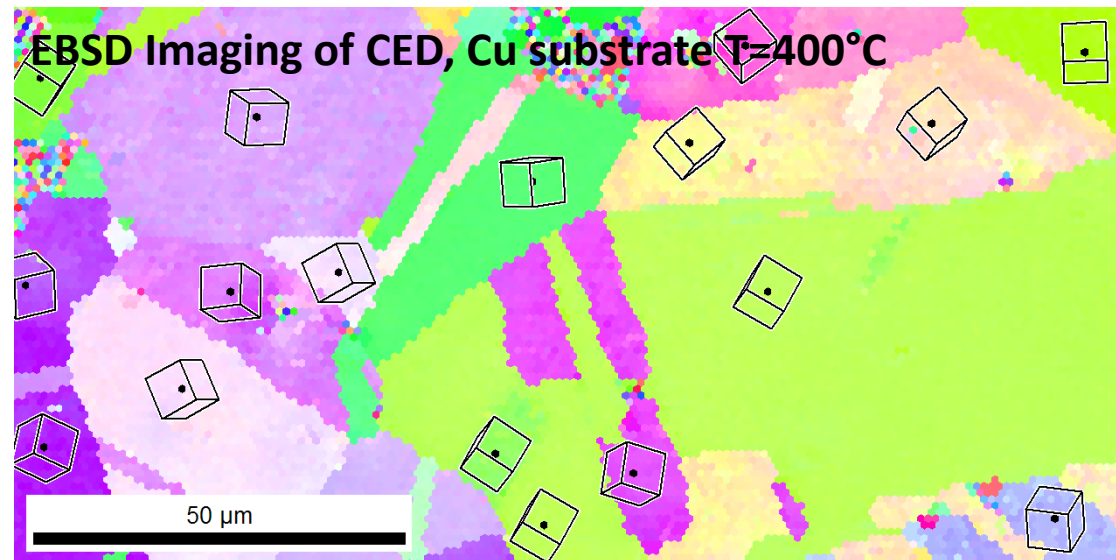
Substrate Al_2O_3 (11-20)

$T = 300^\circ$

$T_c = 9.25\text{K}$

RRR=131

CED Nb Thin Film on a Polycrystalline Cu substrate
Demonstrated Feature of Heteroepitaxy



Working Distance:
15.000000

Number of points:
13094
Number of good
points: 13093

150.00 microns x 74.48
microns
Step: 1.00 microns

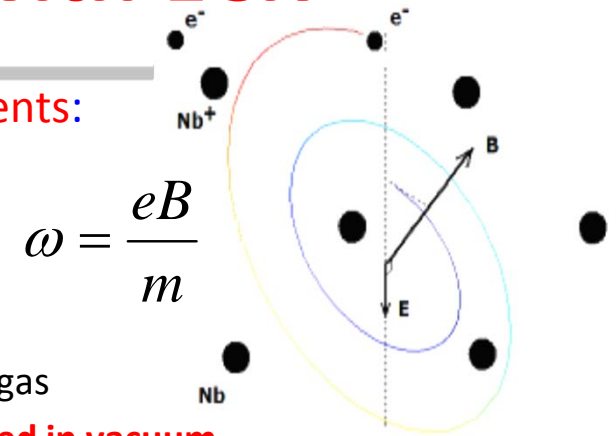
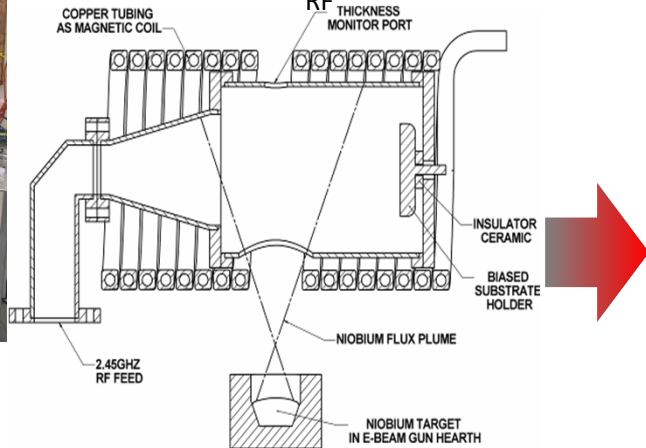
Average Confidence
Index: 0.42
Average Image Quality:
2633.56
Average Fit [degrees]:
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Energetic condensation with ECR

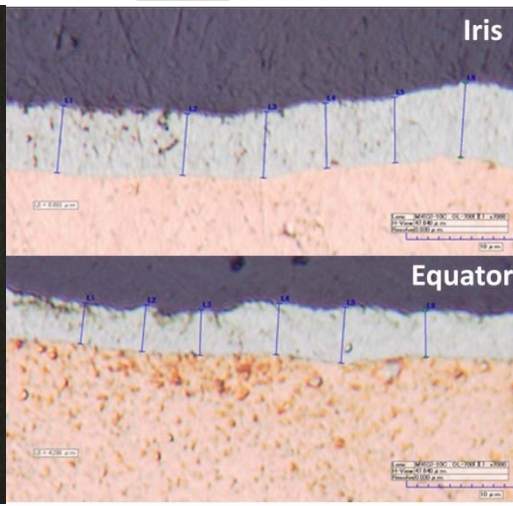


Generation of plasma - 3 essential components:

- Neutral Nb vapor
- RF power (@ 2.45GHz)
- Static B \perp E_{RF} with ECR condition

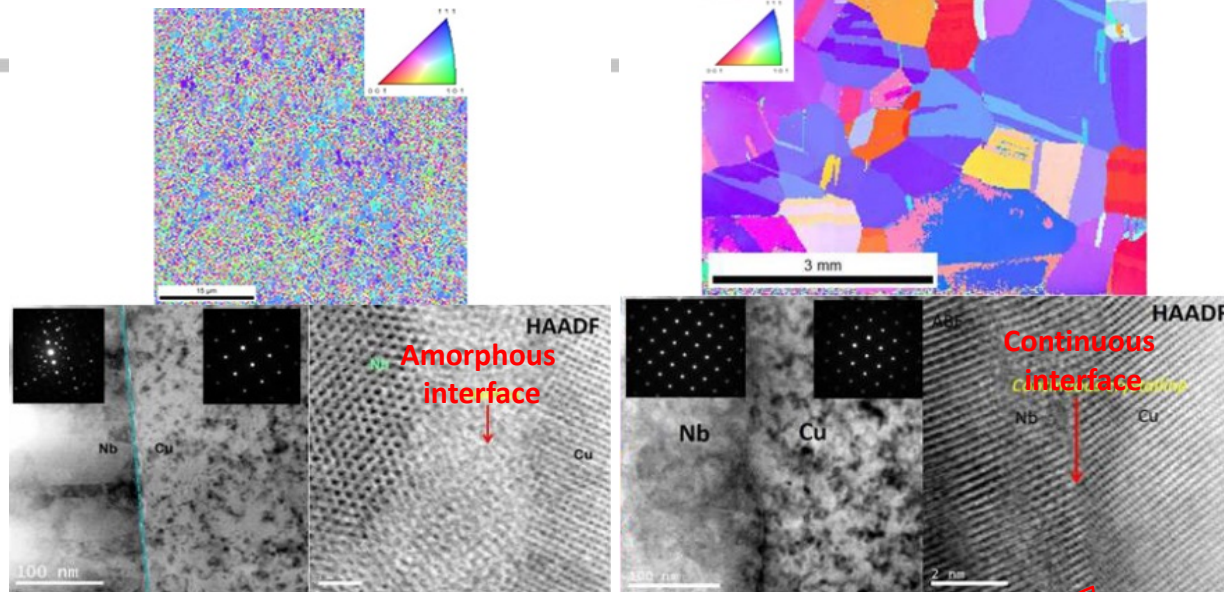


- No working gas
- Ions produced in vacuum**
- Singly charged ions 64eV
- Controllable deposition energy** with Bias voltage
- Excellent bonding
- No macro particles
- Good conformality

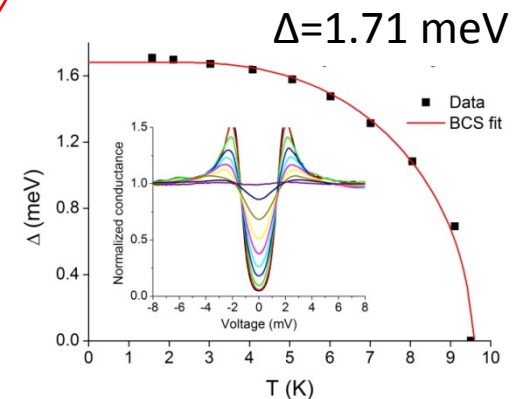
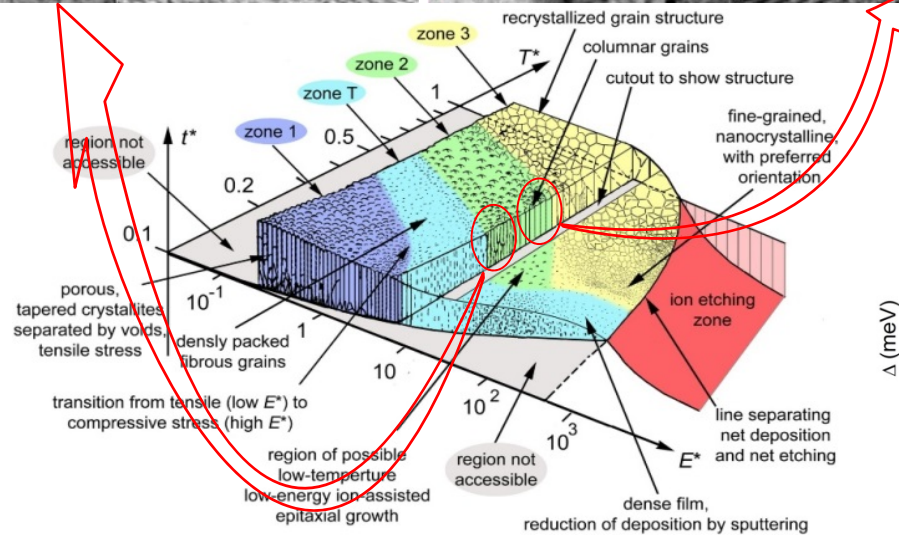
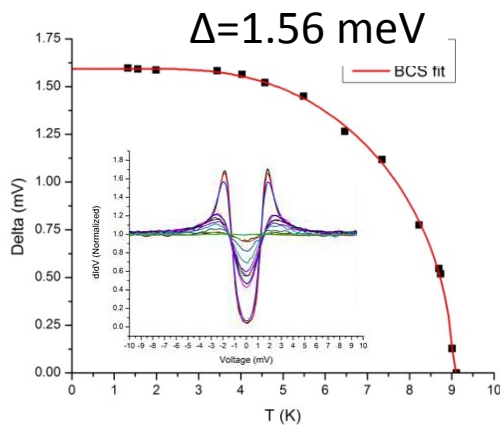


Conformality of the ECR process:
 the film thickness along a 3GHz half-cell profile varies from 4mm (equator) to 6mm (iris)
 Note: the substrate is very rough, was only grossly mechanically polished

Structure, interface and superconducting gap



Gap measurements performed by PCT (point contact tunneling spectroscopy - T. Prosljer)



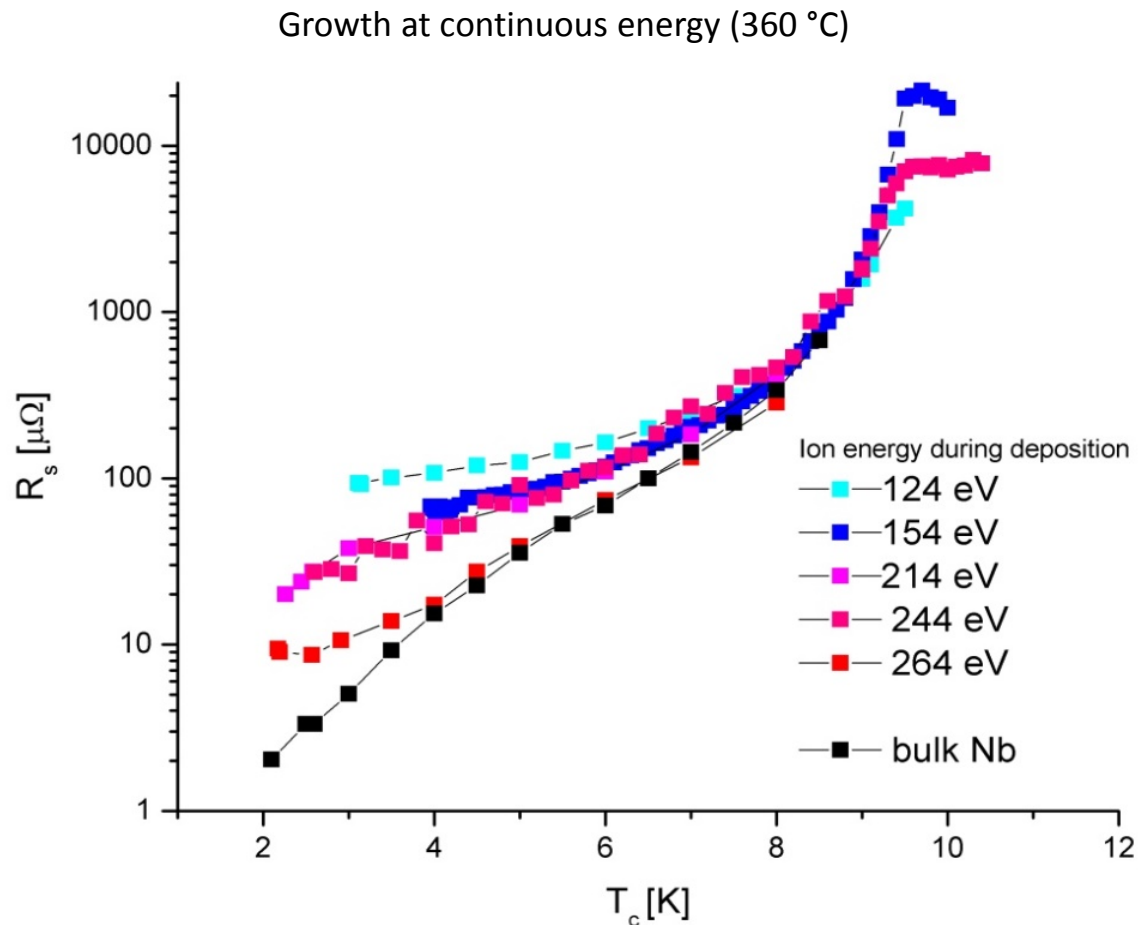
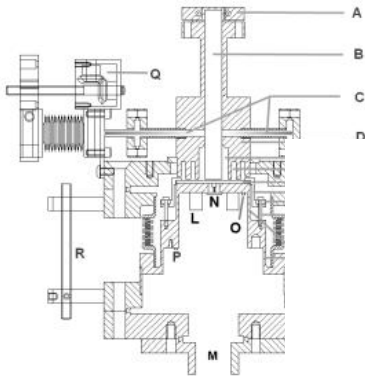
Superconducting gap (1.56-1.62 meV) similar to bulk Nb ($\Delta_{\text{Nb bulk}} = 1.55 \text{ meV}$ measured on the same setup) for hetero-epitaxial ECR Nb films on polycrystalline Cu.

Tailored Nb films via energetic condensation

- ❑ Tune thin film structure and quality with ion energy and substrate temperature on a variety of substrates (amorphous, polycrystalline and single crystal)
- ❑ Achieve film structures and properties only achievable at higher temperature with classic coating methods
- ❑ Tune RRR values from single digits to bulk Nb values → No intrinsic limitations
- ❑ Lower impurity (H) content than bulk Nb
- ❑ Good adhesion to the substrate (delamination threshold determined as function of ion energy and temperature)
- ❑ Grain boundaries not necessarily detrimental (if dense) to R_s
- ❑ Tailoring interface with high energy and subsequent growth at energy minimizing defect creation can contribute to lower R_s

		Substrate	RRR _{max}
Insulating	Single crystal	a-Al ₂ O ₃	488
		r-Al ₂ O ₃	641
		c-Al ₂ O ₃	247
		MgO (100)	188
		MgO (110)	424
		MgO (111)	270
	Polycrystalline amorphous	Al ₂ O ₃ ceramic	135
		AlN ceramic	110
		Fused Silica	84
Metallic	Single crystal	Cu (100)	181
		Cu (110)	275
		Cu (111)	245
	Polycrystalline	Cu fine grains	193
		Cu large grains	305

Influence of ion energy on surface resistance

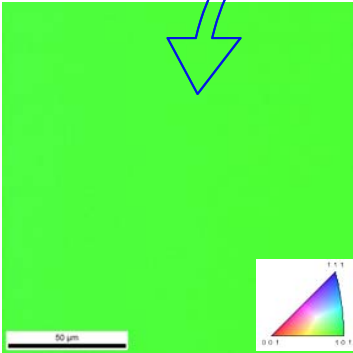
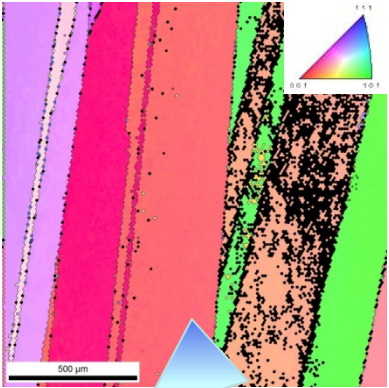
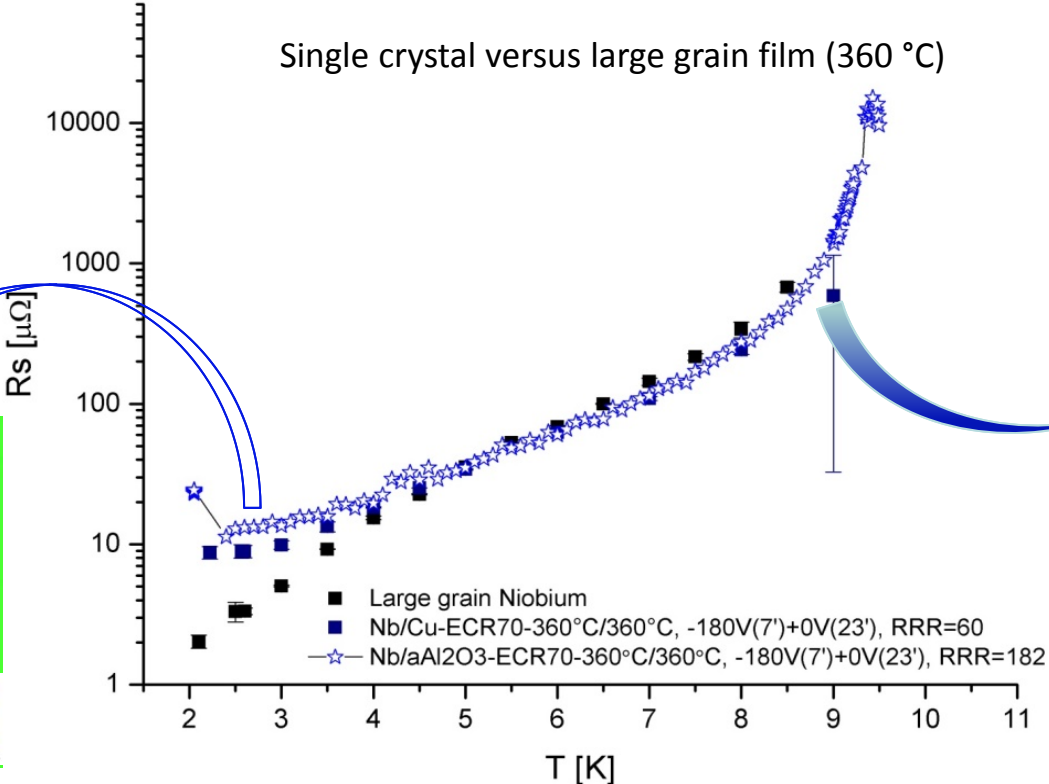


Decrease of R_s with increasing incident ion energy

ECR Nb/Cu – surface resistance

Polycrystalline Nb

Single crystal versus large grain film (360 °C)



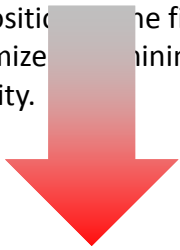
Nb single crystal (110)

If dense, grain boundaries not necessarily detrimental to RF performance

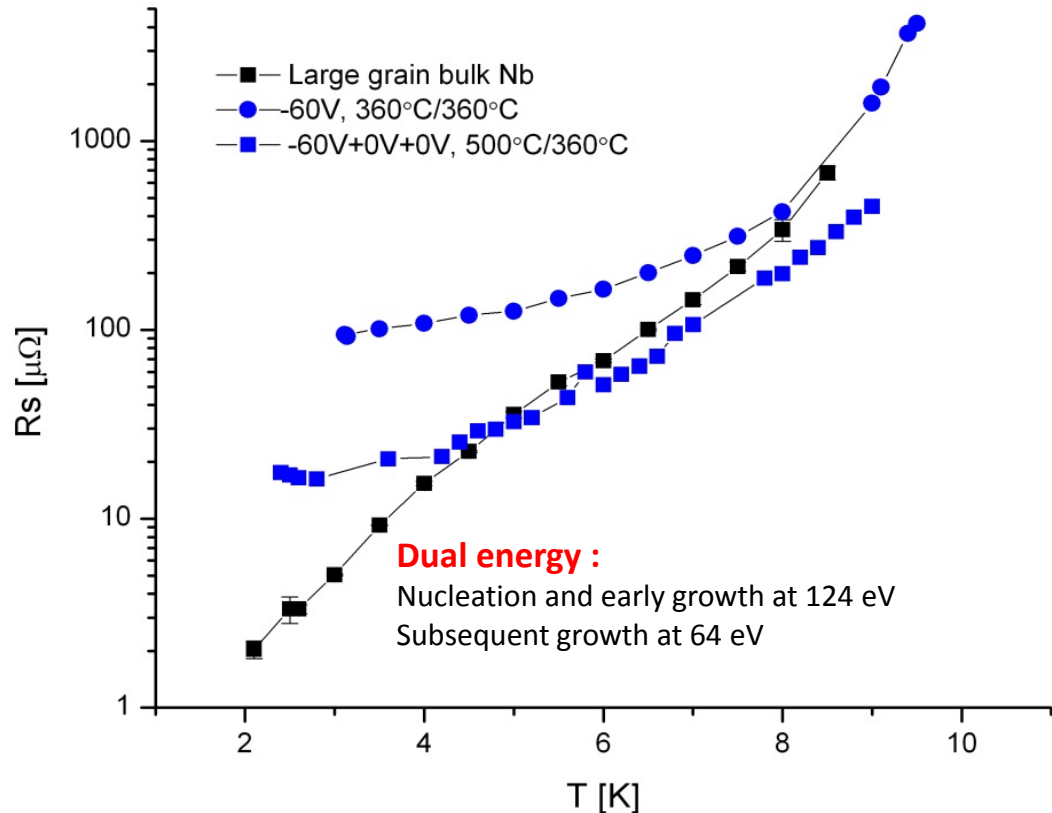
ECR Nb/Cu – Surface resistance

Approach: 3 sequential phases for film growth

- ❑ Film nucleation on the substrate (Nb, Al₂O₃, Cu; single crystal, polycrystalline, amorphous)
- ❑ Growth of an appropriate template for subsequent deposition
- ❑ Deposition of the final surface optimized for minimum defect density.



- ❑ **Film nucleation on the substrate**
- ❑ **Subsequent Growth**



**Addressing the film deposition in 2 phases
(nucleation @ high energy, subsequent growth @ 64 eV)
shows some improvement in R_s**

ECR Nb/Cu– Surface resistance

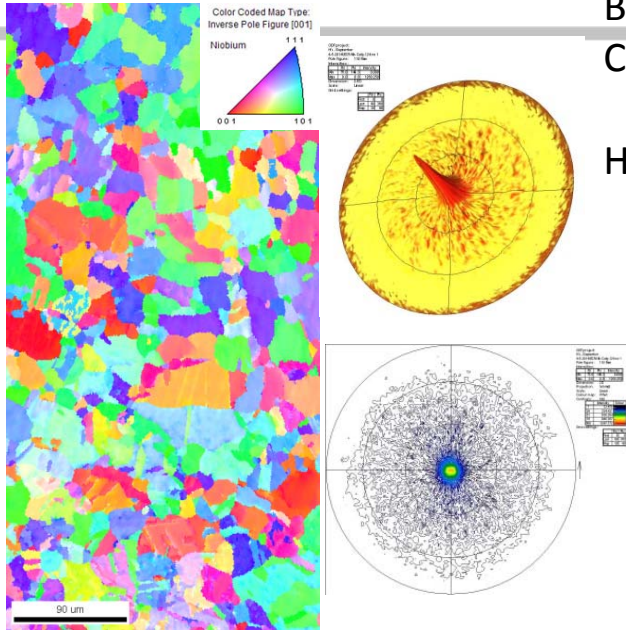
Bake & coating temperature: 360 °C

Coating with dual ion energy: 184 eV for nucleation/early growth
64 eV for subsequent growth

Hetero-epitaxial film Nb on OFHC Cu

$$T_c = 9.36 \pm 0.12 \text{ K}$$

RRR = 99 (witness sample, 179 for Nb/a-Al₂O₃)
EBSD IPF map and XRD pole figure show very good crystallinity and grain sizes in the range of the typical Cu substrate



	R_{res} [nΩ]	$\lambda(0K)$ [nm]
400 MHz	46.6 ± 0.8	40 ± 2
800 MHz	79 ± 2	38 ± 1
1200 MHz	156 ± 11	38 ± 1

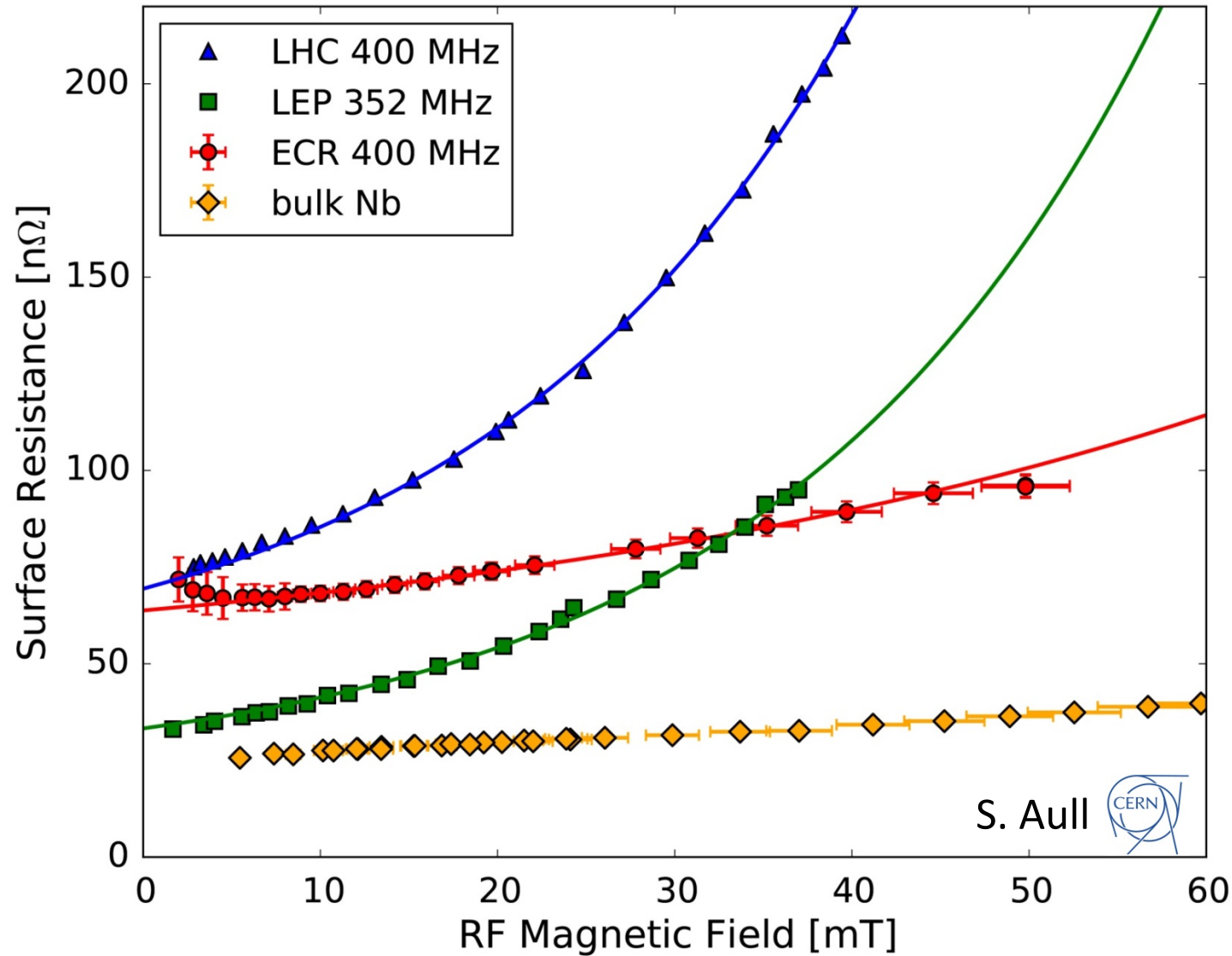
ℓ^* [nm]	RRR
144 ± 20	53 ± 7

* with $\lambda_L = 32$ nm and $\xi_0 = 39$ nm

**Nb film
in the clean limit**



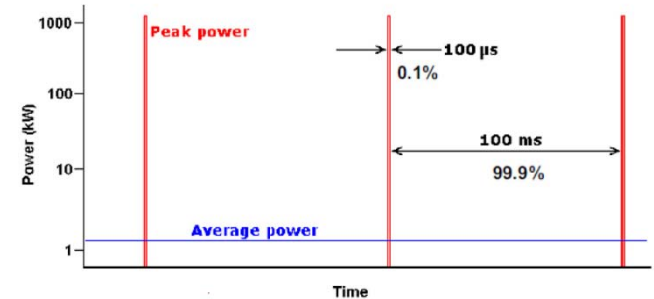
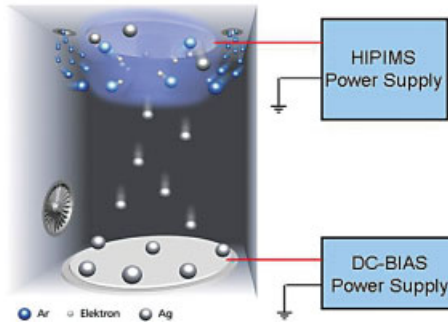
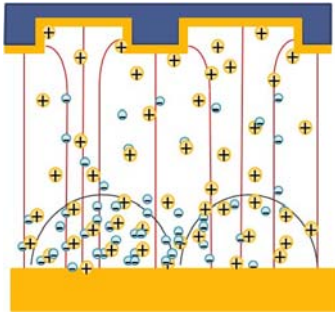
Comparison with Sputtered Nb/Cu and bulk Nb



S. Aull 

HiPIMS

High power impulse magnetron sputtering (HiPIMS) is an emerging technology

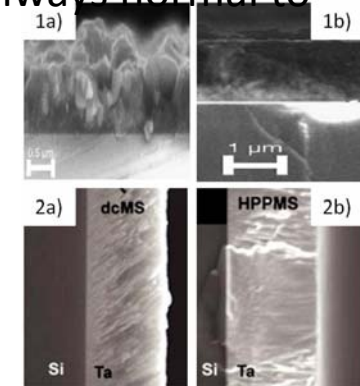


Use of High Power Impulse Magnetron Sputtering (HiPIMS) for Nb coating with enhanced properties

The target material is partly ionized such that the ions will reach the substrate with a higher energy and, in case of an applied bias voltage to the substrate, always normal to the surface.

- very high purity
- better adhesion
- better (normal) conductivity,
- Large crystal grains, low defect density
- dense & smooth films

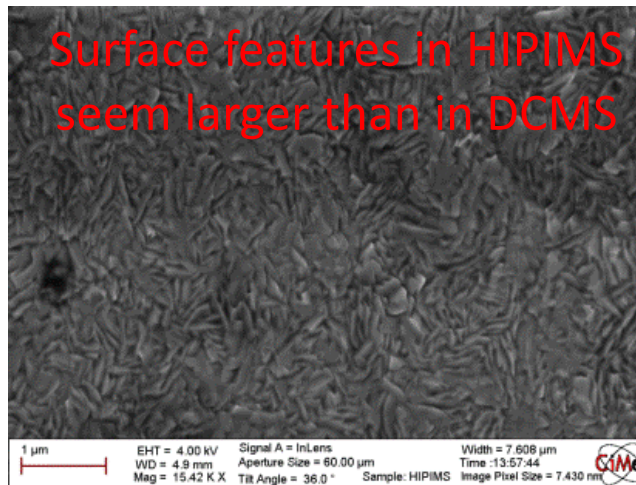
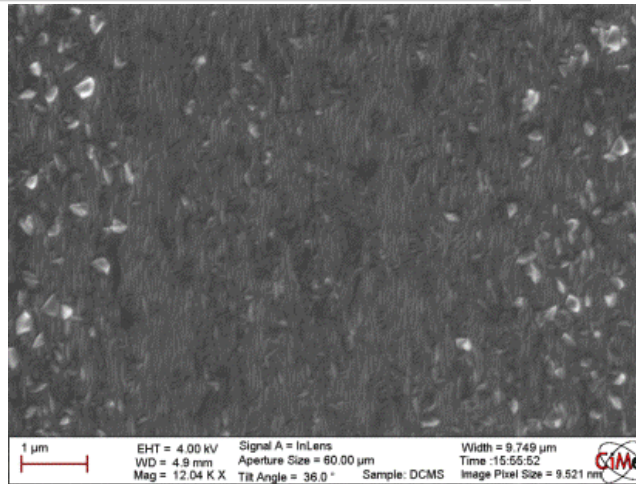
→ Homogeneous coating even on non-flat surfaces



NEW METHODS FOR THIN FILM DEPOSITION AND FIRST INVESTIGATIONS OF THE USE OF HIGH TEMPERATURE SUPERCONDUCTORS FOR THIN FILM CAVITIES

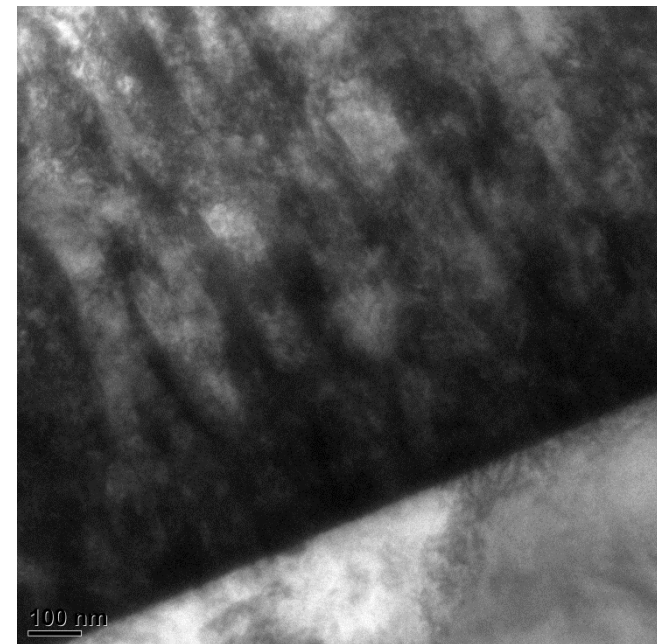
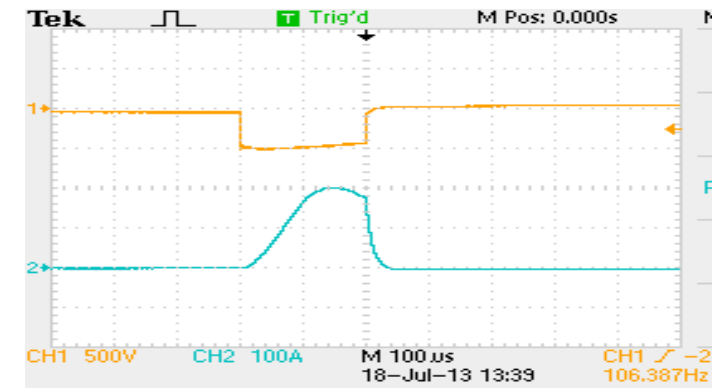
A. Gustafsson *et al.*, CERN (WEPEC047)

Nb films – Energetic condensation with HiPIMS



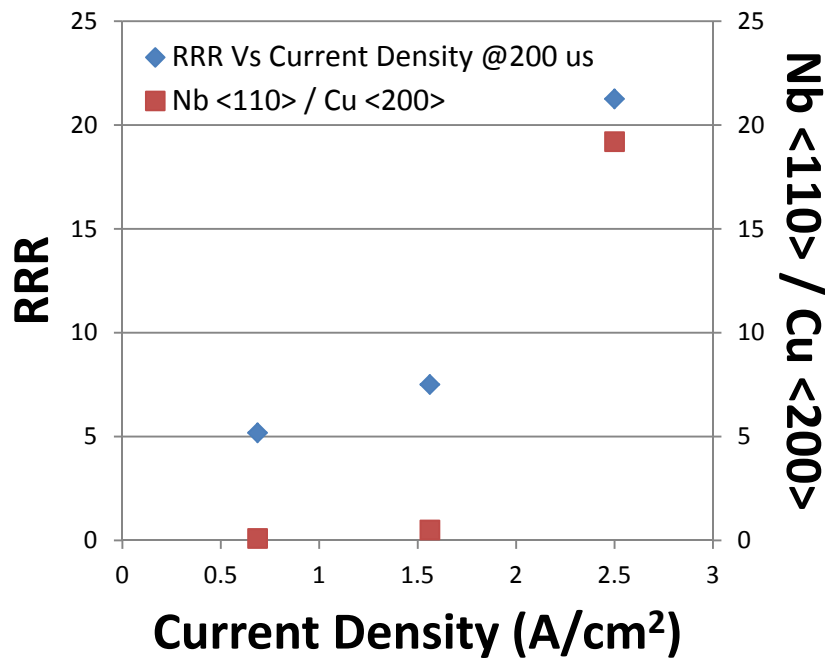
Surface features in HiPIMS seem larger than in DCMS

CERN (G. Terenziani, S. Calatroni, T. Junginger)
Sheffield-Hallam University (A.P. Ehisarian)

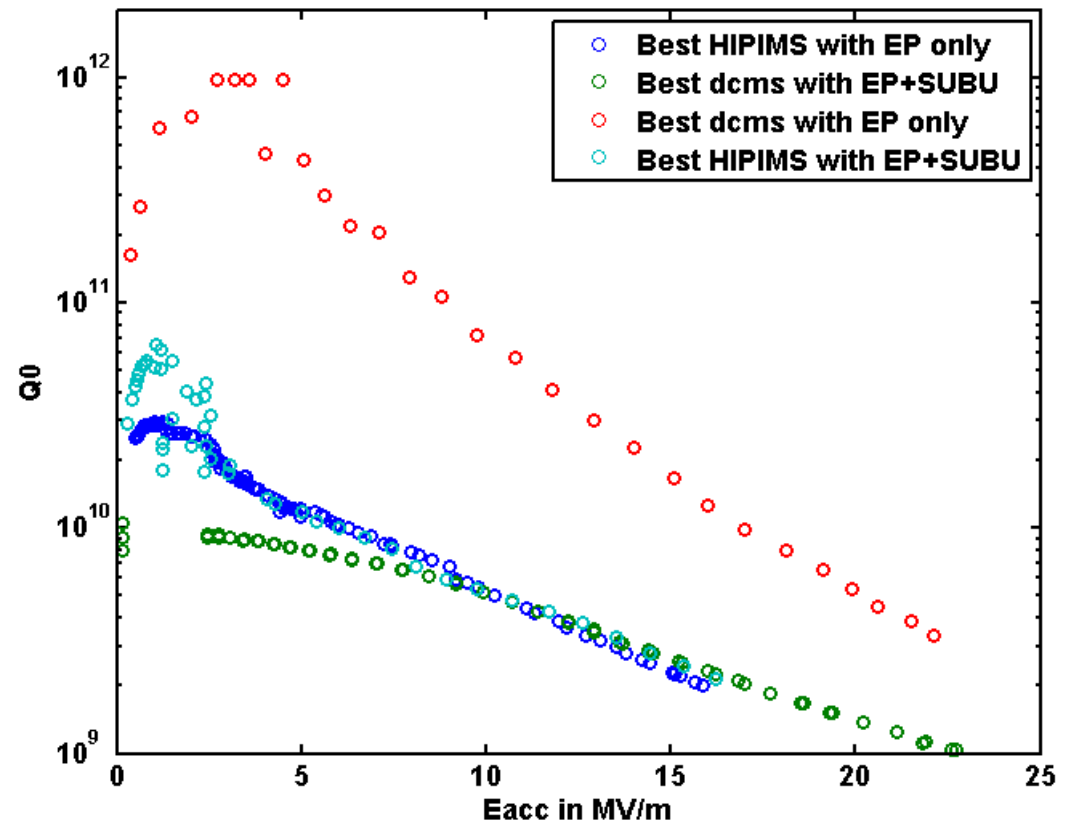


Nb films – Energetic condensation with HiPIMS

Comparison RRR Vs Crystallographic Orientation

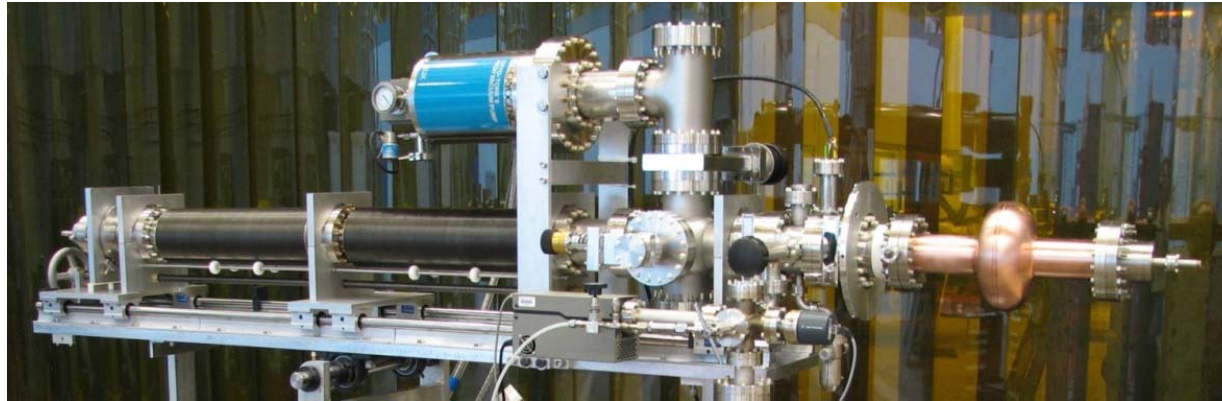


1.3 GHz Cavity



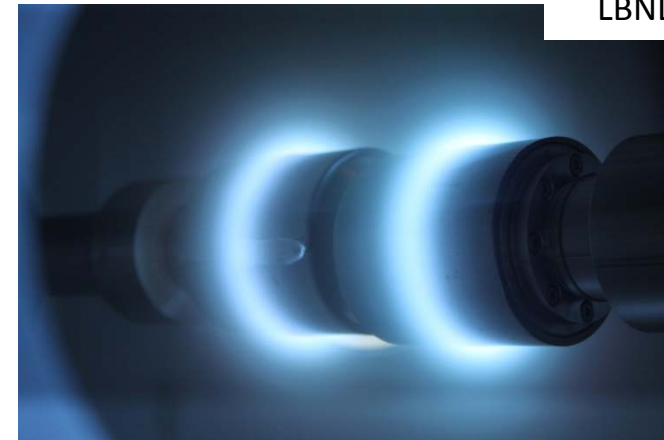
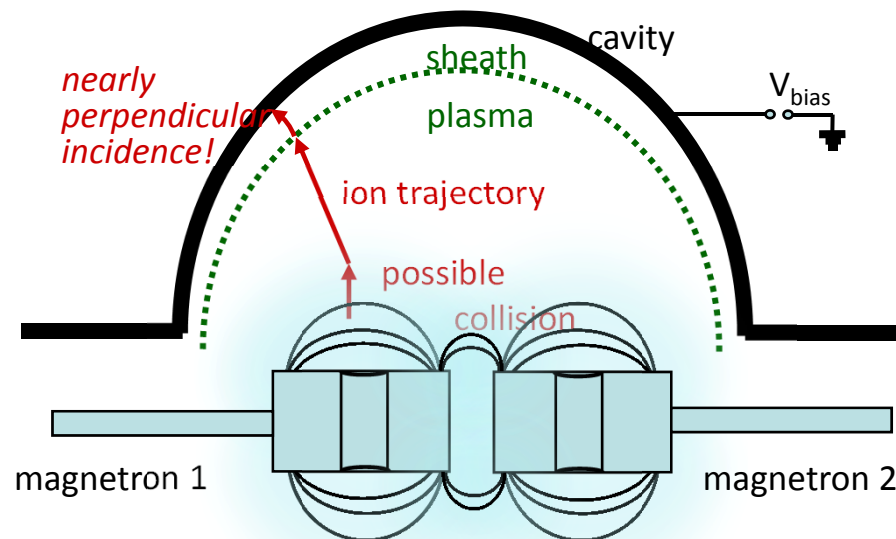
Nb films – Energetic condensation with HiPIMS

HiPIMS cylindrical coating system for single cell under commissioning



HiPIMS dual magnetron system Most effective for Biasing & influencing Ion Energies & Trajectories

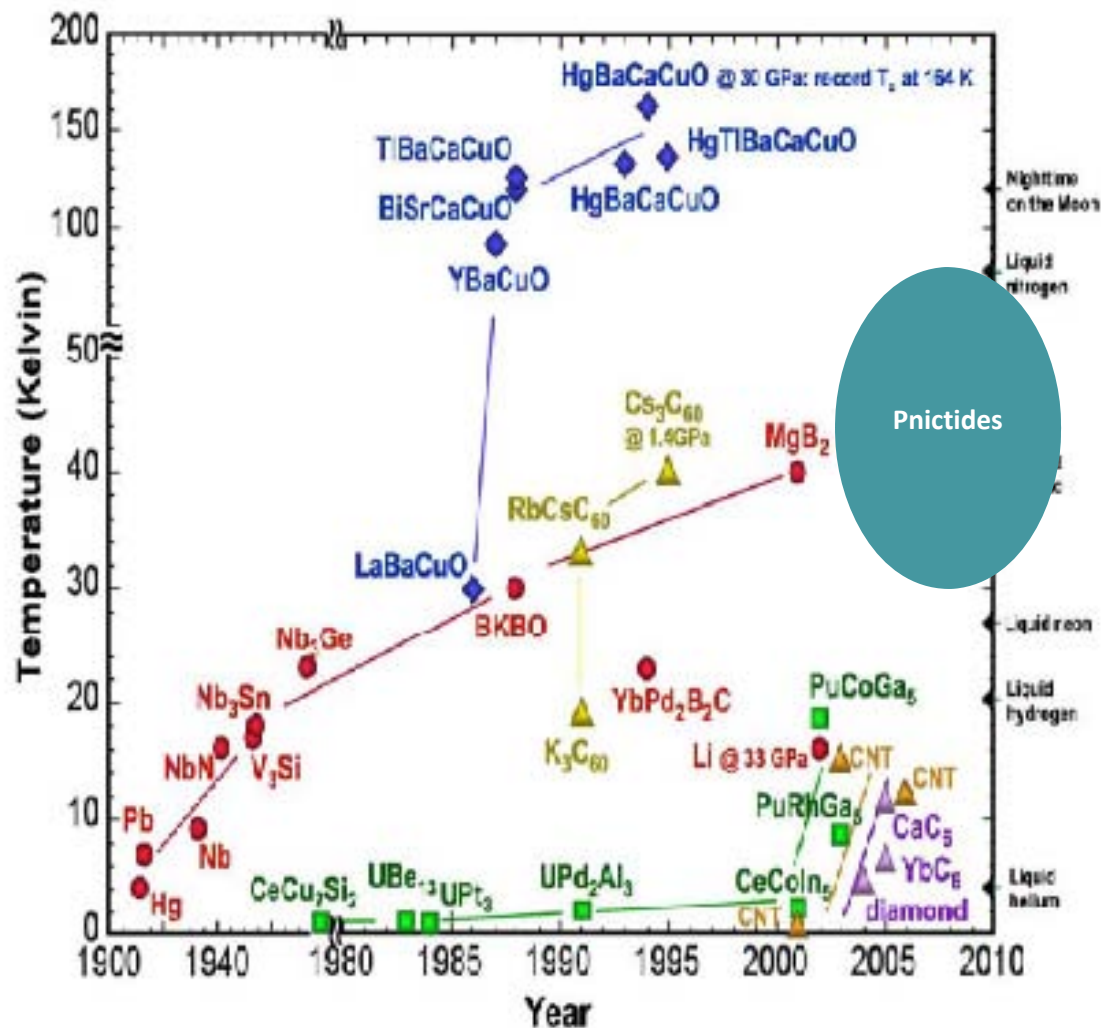
LBNL(A. Anders)



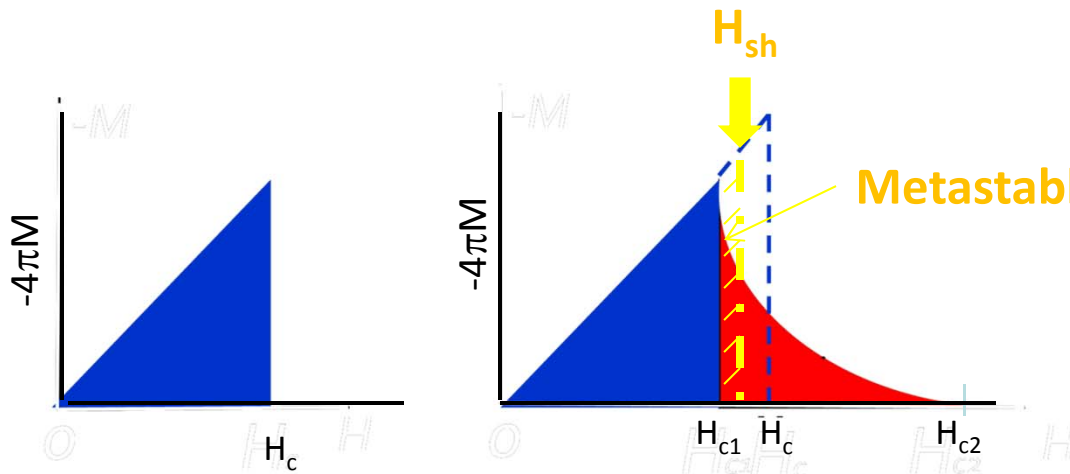
- high power mode (above runaway threshold)
- Dominated by niobium emission

Which materials are suitable for SRF cavities?

Highest $T_c=164$ K (under GPa)



Critical Field



Type-I

Type-II

- Meissner state at $0 < H < H_{c1}$
- Mixed vortex state at $H_{c1} < H < H_{c2}$
- Exponentially small R_s at $H < H_{c1}$ ($Q = 10^{10}-10^{11}$)
- Drastic Q drop due to vortex dissipation at $H > H_{c1}$

Boundary between Type I and Type II determined by the Ginzburg-Landau parameter

$$\kappa = \lambda/\xi$$

H_c is the thermodynamic critical field &

$$H_c = H_{c2}/\sqrt{2}\kappa$$

(for type-II superconductors).

The superheating field H_{sh} is the field up to which the Meissner state metastably persists above H_{c1}

For type-II superconductors, at $T = 0$ in the clean limit with a Ginzburg-Landau parameter $\kappa = \lambda/\xi \gg 1$, H_{sh} has been calculated to be $\sim 0.75H_c$.

$$H_{RFcrit} \approx H_{sh}$$

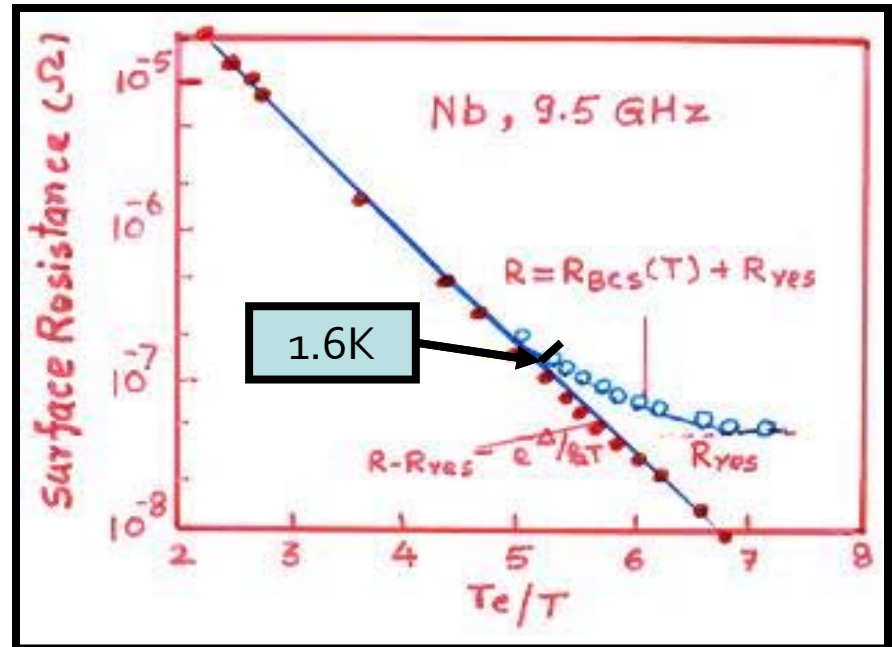
Surface Resistance

The power dissipated per unit area of SC in RF regime
and

$$P = R_s \cdot \frac{H^2}{2}$$

Surface Resistance

$$R_s = R_{BCS}(T) + R_{res}$$



V. Palmieri, 10th Workshop on RF Superconductivity Proceedings, Tsukuba 2001 (Noguchi)
"New materials for superconducting radiofrequency cavities"

BCS Surface Resistance R_{BCS}

If $T < T_c / 2$, for dirty limit superconductors

$$R_{BCS} \cong \frac{R_n}{\sqrt{2}} \left(\frac{\eta\omega}{\pi\Delta} \right)^{\frac{3}{2}} \frac{\sigma_1}{\sigma_n} = A \sqrt{\rho_n} e^{-\frac{\Delta}{K_B T}} (1 + O(\Delta, \omega, T))$$

A constant weakly dependent on material

ω = RF frequency

ρ_n = Normal State conductivity

Δ = Superconducting gap

T_c = Transition Temperature

dependence on ρ_n and T_c represents an immediate criterion for selecting the most favorable candidates for cavities

The higher T_c (= 0.57Δ), the smaller the BCS surface resistance

Material with high normal state conductivity and high T_c (high superconducting gap Δ) should be selected

Residual Resistance R_{res}

Temperature independent

Contributions to residual losses:

Intrinsic:

Inhomogeneties, Metallic Inclusions within I, Grain Boundaries, Oxides

Extrinsic:

Trapped Flux during cooling (can be avoided)

Variety of phenomena involved



Not one formula predicting R_{res}

From literature

Empirically, R_{res} found proportional to at least $\sqrt{\rho_n}$

For two materials with the same R_{BCS} and different T_c and ρ_n ,
the one with the smallest ρ_n should have the smallest R_{res}

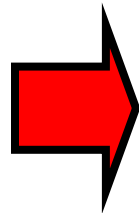
Metallic behaviour is favored

Criteria of choice

THERE IS NO IDEAL SUPERCONDUCTOR FOR CAVITY

CHOICE IS BASED ON COMPROMISE

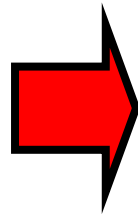
For low RF losses



high T_c

**Metallic behavior in the normal state ,
small r_n**

For high gradients



High H_{sh} , H_c ; small k

Possible Choices among Superconducting Materials

Material	T_c [K]	$\rho_n(\mu\Omega\text{cm})$	$H_c(0)$ [T]	$H_{c1}(0)$ [mT]	$H_{c2}(0)$ [T]	$\lambda(0)$ [nm]	Δ [meV]	Type
Nb	9.22	2	0.2	180	0.28	40	1.55	II
Pb	7.2		0.08	N/A	N/A	48		I
NbN	16.2	70	0.23	20	15	200	2.6	II, B1 comp.
NbTiN	17.5	35		30	15	151		II, B1 comp.
Nb ₃ Sn	18	20	0.54	50	30	85	3.1	II, A15
V ₃ Si	17				24.5	179		II, A15
Mo ₃ Re	15		0.43	30	3.5	140		II, A15
MgB ₂	40	0.1-10	0.43	30	3.5-60	140	2.3; 7.2	II- 2 gaps
Pnictides	30-55		0.5-0.9	30	>100	200	10-20	
YBCO	93		1.4	0.01	>100	150	20	d-wave

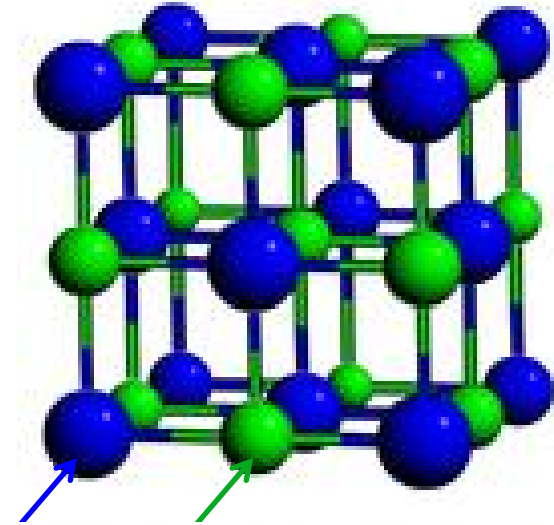
Large gap Δ (good for SRF) is usually accompanied by low H_{c1} (bad for SRF)

$$R_s = \mu_0^2 \sigma_n \lambda^3 \omega^2 \frac{\Delta}{kT} \ln \left(\frac{9kT}{4\hbar\omega} \right) e^{-\Delta/kT}$$

d-wave high- T_c cuprates with nodal gap and $R_s \propto T^s$, $s \simeq 2-3$ are not useful for SRF

B1 compounds – NaCl structure

Metallic atoms A form an fcc lattice
 Non-metallic atoms B occupy all the octahedral interstices.



NB COMPOUNDS

B \ A	Sc	Y	La	Ti	Zr	Hf	V	Nb	Ta	Cr	Mo	W	Re
B					3.4	3.1							
C	<1.38	<1.38		3.42	<0.3	<1.20	0.07 3.2	12	0.35		14.3	10.0	3.4
N	<1.38	<1.4	1.35	5.49	10.7	8.83	8.5	17.3	6.5	<1.28	5.0	<1.38	
P			<1.68										
Sb		<1.02	<1.02										
O				2.0			<0.3	1.39					
S	<0.33	1.9	0.87		3.3								
Se	<0.33	2.5	1.02										
Te		2.05	1.48										

* $T_c = 3.2$ K was registered in vanadium carbide after implantation of C^+ ions ₅

NbN phase diagram

Good SC properties, even if deposited at low temperature
 Low secondary emission coefficient
 Very stable surface properties

The right B1-NbN superconducting phase is the so-called δ -phase

$T_c = 17.2$ K for δ -phase (lattice parameter = 4.388 Å)
 T_c very sensitive to Nitrogen stoichiometry

In sputtered films, the δ -phase can be found mixed to some other low T_c phases

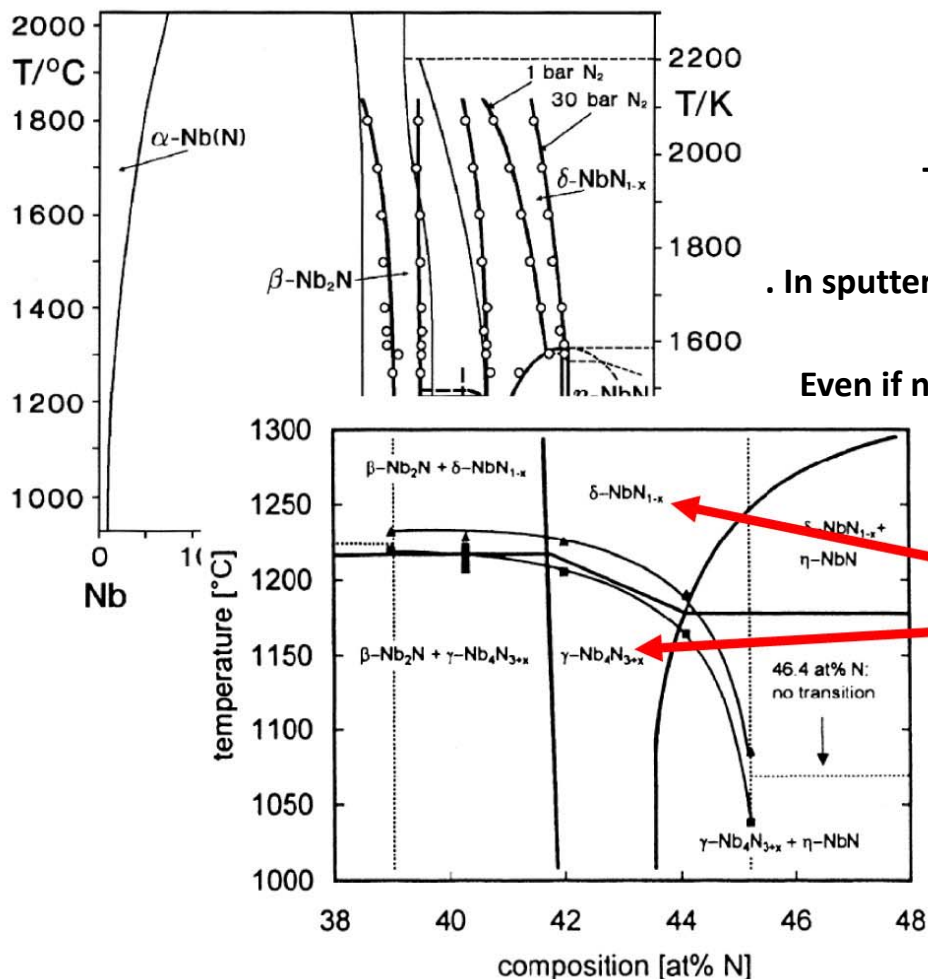
Even if no grain boundaries are present and δ -phase single crystal is considered

the single grain resistivity is not so low.

δ -NbN $\rightarrow T_c \sim 15 - 17.3$ K
 γ -NbN $\rightarrow T_c \sim 12 - 15$ K

Anomalously high resistivity of NbN in the normal state, often higher than $100 \mu\Omega\text{cm}$ due to both metallic and gaseous vacancies randomly distributed in both sublattices

Equiatomic composition is $\text{Nb}_{0.987}\text{N}_{0.987}$ not $\text{Nb}_{1.0}\text{N}_{1.0}$
 Common problem for B1 compounds



Nb Compounds - NbN

The only B1 simple compound that has widely tested for accelerating cavities

Mainly two different techniques have been investigated for this application:

- Thermal diffusion of N into Nb followed by rapid quench cooling
- Reactive Sputtering on metallic or ceramic substrates to Nb cavities

Thermal Diffusion:

Bulk Nb (RRR 300) annealed @ 1550°C for 2h

+

reacted in N₂ vapor(150mbar) @ 1400°C for 4h
Rs=1.3 10⁻⁶ @ 4.2K and 4 10⁻⁹ @1.8K @7.9GHz

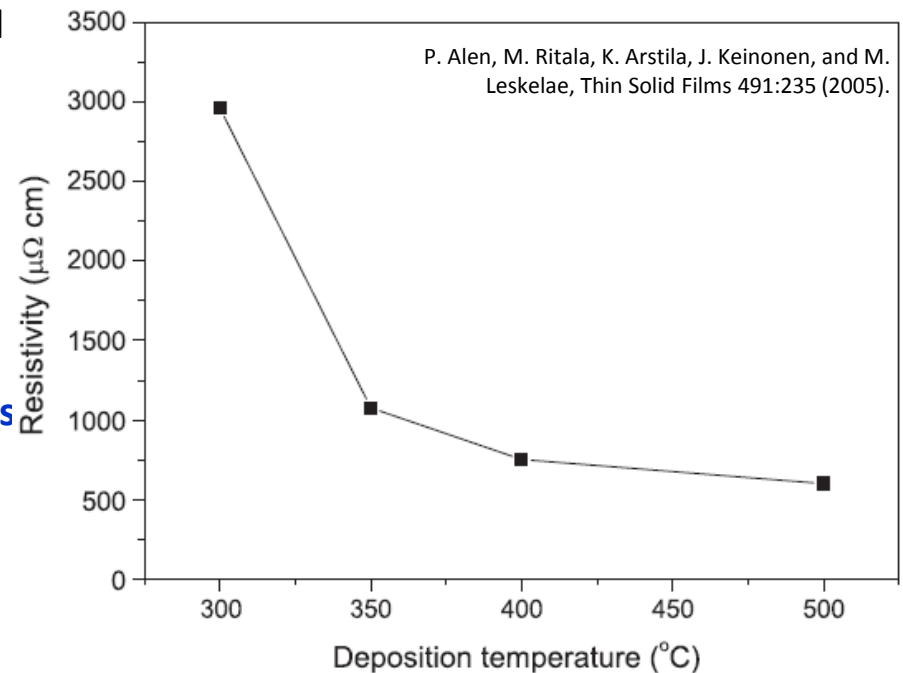
G.Gemme et al., J.Appl.Phys. 77(1), Jan. 1995

Reactive Sputtering:

Sputtering from high purity Nb target in Ar+ N₂ in DC triode magnetron sputtering system

Highest T_c for substrate temp. > 500°C, P_{Ar}=8.10⁻³mbar, P_{N2}=1.10⁻³mbar

A. Nigro et al., Physica Scripta Vol. 38, 483-485, 1988



Nb Compounds - NbTiN

Ternary Nitride $\text{Nb}_{1-x}\text{Ti}_x\text{N}$

Presence of Ti found to reduce significantly the resistivity
And facilitate formation of a pure cubic structure.
The δ -phase remains thermodynamically stable even at RT.
 T_c as high as for good quality NbN, for Nb fraction $(1-x) > 0.5$

extreme hardness, excellent adherence on various substrates, very good
corrosion and erosion resistance, high-sublimation temperature, and relative
inertness

**More metallic nature and better surface properties than NbN
should result in better RF performance**

Nb Compounds - NbTiN

INFN : reactive sputtering with Ar/N₂ in DC Triode Magnetron Sputtering @ 600°C and 200°C

(Nb_{1-x}Ti_x)N films with 1-x < 0.5 present a lower calculated surface impedance, lower critical fields and better surface properties than NbN, especially when deposited at low temperatures.

R. Di Leo et al. J. of Low Temp. Phys, vol 78, n1/2, pp41-50, 1990

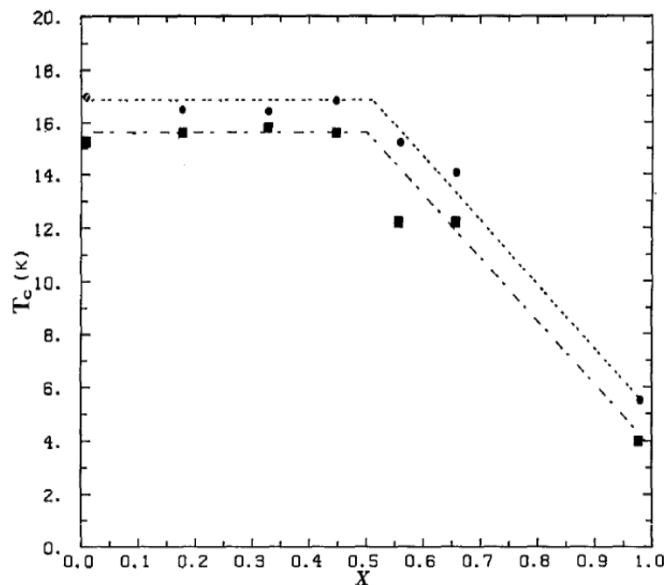


Fig. 1. Superconducting critical temperature T_c as a function of the titanium composition (x) for the $(\text{Nb}_{1-x}\text{Ti}_x)\text{N}$ films deposited at $T_s = 600^\circ\text{C}$ (circles) and at $T_s = 200^\circ\text{C}$ (squares).

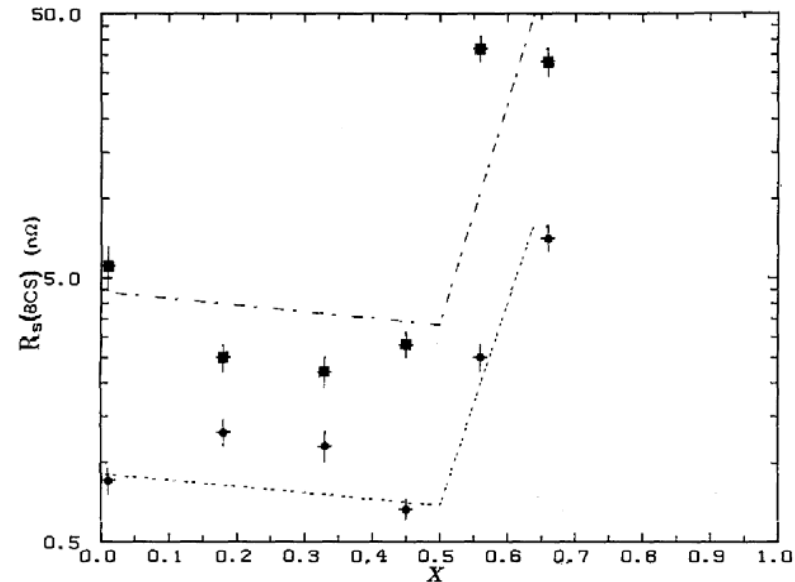


Fig. 3. Calculated BCS surface impedance $R_s(\text{BCS})$ as a function of the titanium composition (x) for the $(\text{Nb}_{1-x}\text{Ti}_x)\text{N}$ films deposited at $T_s = 600^\circ\text{C}$ (circles) and at $T_s = 200^\circ\text{C}$ (squares). The continuous lines correspond to the values of the lines through the data in Figs. 1 and 2.

Nb Compounds - NbTiN

Reactive Magnetron Sputtering:

CEA Saclay :

NbTiN films deposited on 12 cm copper disks by magnetron sputtering and tested in a cylindrical TE_{011} cavity reached RF field levels of 35 mT

low residual surface resistance ($< 100 \text{ n}\Omega$ at 4 GHz) with a very small BCS resistance

4 cavities deposited but no RF measurement due to film blistering on large area of the cavity.

R_s slope significantly decreased when coating with bias ranging from -50V to -100V

P. Bosland et al.

S. Cantacuzène et al.

CERN:

Samples and six 1.5 GHz Cu cavities coated by reactive cylindrical magnetron sputtering

Best cavity result for thicker film ($4.3\mu\text{m}$) and lower deposition temperature (265°C)

$R_s = 330\text{n}\Omega$ @ 4.2K

Q_0 at zero field is higher than the Q-value of Niobium cavities but Eacc limited under 10 MV/m

As for NbN, N stoichiometry critical to obtain the right SC phase

M. Marino, Proceedings of the 8th Workshop on RF Superconductivity, October 1997, Abano Terme (Padua), (Rep) 133/98, vol.IV, p.1076

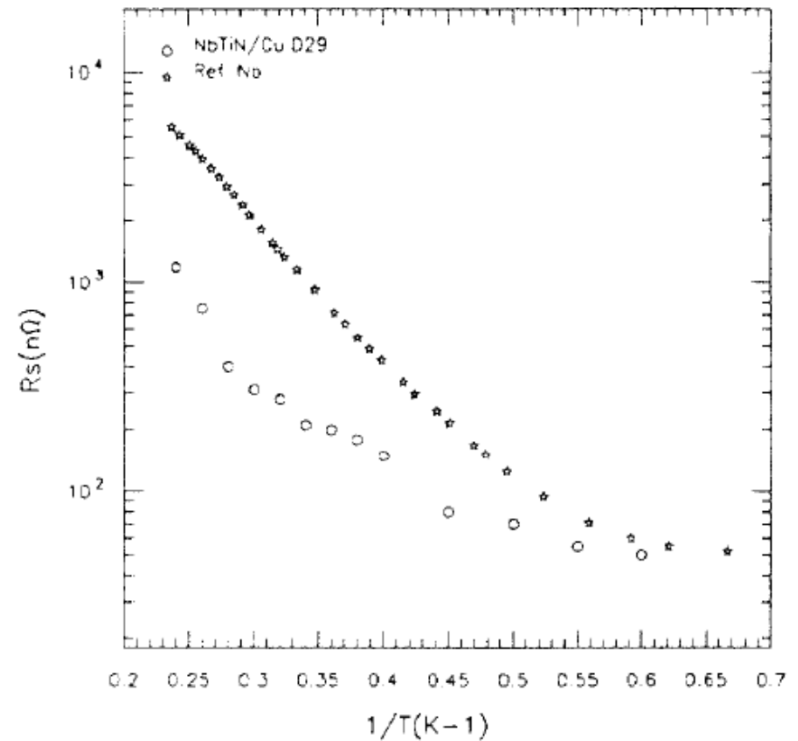


Figure 3 Surface resistance vs temperature for a NbTiN sample, at 4 GHz.

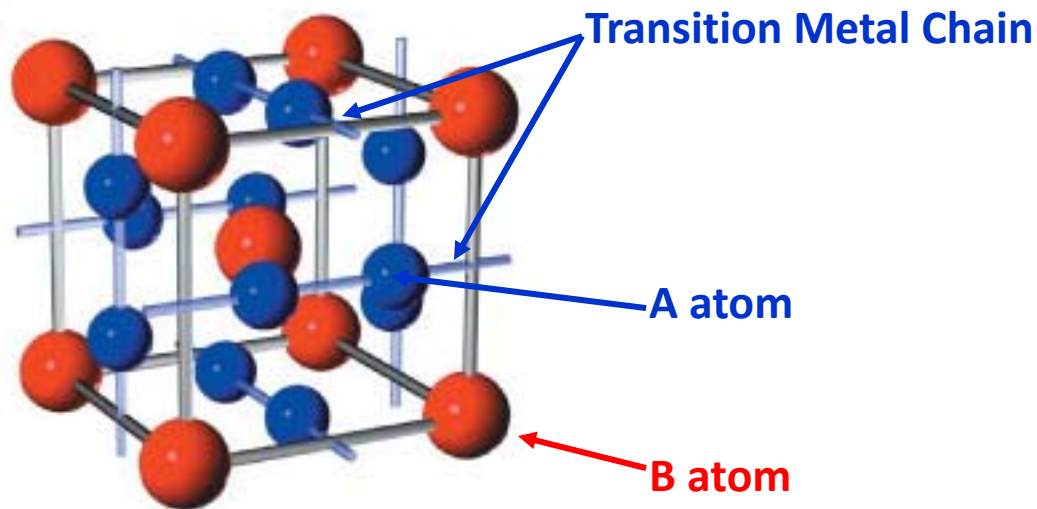
A15 COMPOUNDS

A15 Compounds - Structure

A atoms = Transition elements of group IV, V or VI

B atoms = Non transition or transition elements

A_3B



B atoms occupy corners and centre of BCC structure

A atoms form orthogonal chains bisecting the faces of the BCC unit cell.

Linear Chain Integrity is crucial for Tc (long-range order required)

A15 Compounds – Potential candidates for RF Cavities

Nb_3Sn , Nb_3Al , Nb_3Ge , Nb_3Ga , V_3Si , Mo_3Re

- ❑ Among the Nb and V based high T_c (15 – 20 K)
- ❑ Nb_3Ga and Nb_3Ge do not exist as stable bulk materials at 3:1 stoichiometry
- ❑ Nb_3Al exists only at high temperature causing excessive atomic disorder
- ❑ Production of above materials need non equilibrium processes
- ❑ V_3Ga , V_3Si & Nb_3Sn are stable bulk material and have high T_c
- ❑ Another A-15 compound holding promise is Mo_3Re ($T_c=15K$)

Sharma, International Workshop on Thin Films and New Ideas for Pushing The Limits Of RF Superconductivity, LNL-INFN, Oct.2006

Extreme brittleness so A-15 bulk structure cannot be formed
The A-15 should be produced as thin layer on the interior of the already formed structure
Such a layer need to be only 1 or 2 microns thick
 $\lambda_L (Nb_3Sn) = 65 \text{ nm}$

Thin film route ideal

A15 Compounds – Preparation Methods III

Chemical Vapor Deposition (CVD)

- ❑ MOCVD (*Metal Organic Chemical Vapour Deposition*) is a particular case of CVD in which the precursor is a metallorganic compound
- ❑ Process in which one or more precursors, present in vapor phase, chemically react on an appropriate warm substrate, giving rise to a solid film
- ❑ Deposition rate and structure of the film depend upon temperature and reagent concentration
 - ⇒ Uniformity of temperature and flow of gaseous over entire cavity surface may be difficult with complex geometry

Diffusion Reaction

- ❑ Technique proved successful for magnet conductor application
- ❑ Simple equipment compared to sputtering and CVD

Co-Sputtering

- ❑ Considerable success achieved in synthesizing difficult materials like Nb_3Ge with highest T_c (~23k) or V_3Si
- ❑ Typically two constituents are sputtered simultaneously onto a temperature controlled substrate
- ❑ Stoichiometry dependent on relative positions of target and substrate (can be manipulated to get perfect stoichiometry)
- ❑ Stoichiometry control difficult over large areas like accelerating system and if stoichiometry range for A-15 phase is narrow.

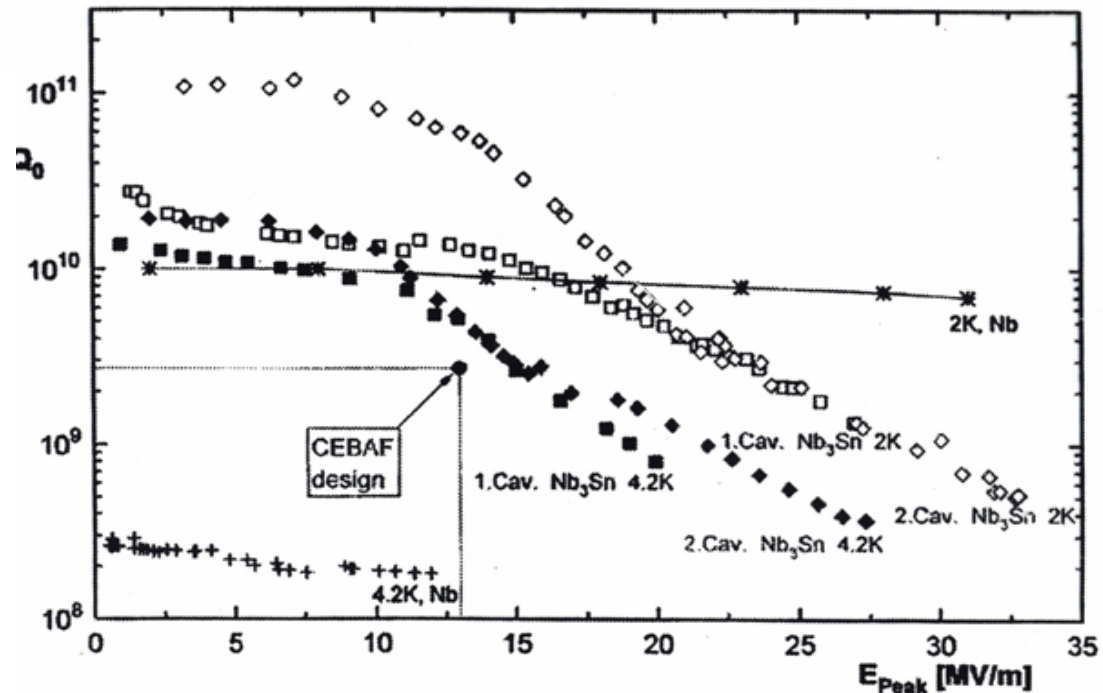
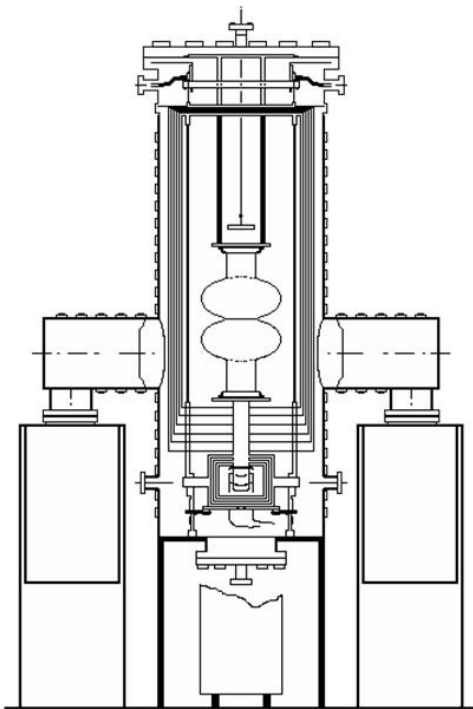
Sputtering

- ❑ To sputter from a single target of correct stoichiometry (prepared by powder sintering)
- ❑ Stoichiometry, Substrate Temperature, Deposition Rate, Deposition Thickness Can be varied independently

Nb₃Sn SRF History

Wuppertal, end '80s :

Nb₃Sn cavity (1.5 GHz) obtained through Sn vapour phase diffusion @ 1200°C



Q vs. E_{peak} of the 1st two Nb₃Sn-coated 1.5GHz single cell cavities in comparison to pure Nb at 4.2K and 2K from CEBAF

5-cell 1.5GHz cavity also coated: $Q_0 \sim 10^9$, $E_{\text{acc}} = 7\text{MV/m}$ with $Q = 8 \cdot 10^8$

G. Müller et al.,

M. Peiniger & H. Piel, *IEEE Trans. On Nucl. Sc.* Vol NS-32, n°5, Oct. 1985

A15 Compounds – Nb₃Sn through liquid diffusion

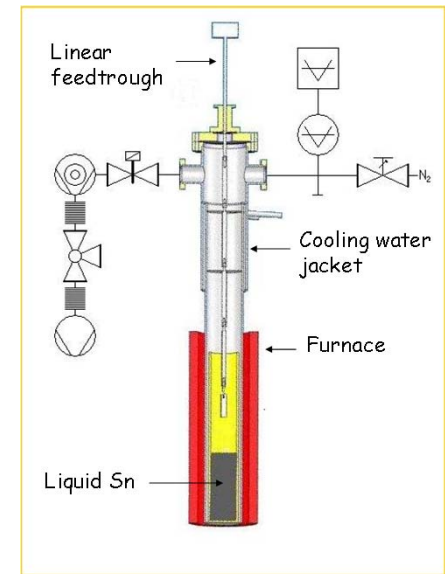
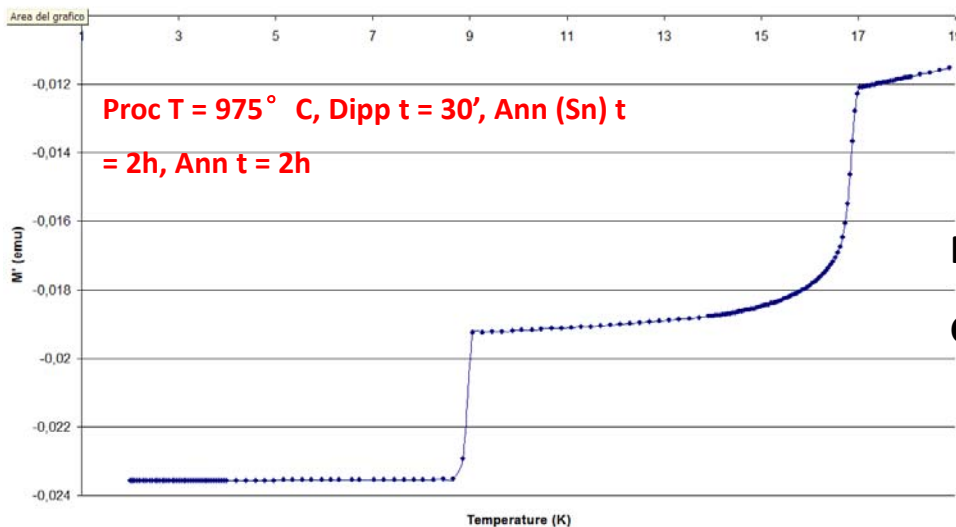
S. Deambrosis, Sharma, *International Workshop on Thin Films and New Ideas for Pushing The Limits Of RF Superconductivity, LNL-INFN , Oct.2006*

S. Deambrosis et al., *Physica C 441 (2006) 108-113*

Nb₃Sn coatings on Nb by liquid diffusion method at INFN Legnaro “Hybrid” Process”

- Substrate thermalization (30 min - 1 h)
- Dipping (few min - 2 h)
- Sample annealing with Sn vapor for a few hours
- Sample annealing without Sn vapor for a few hours

Nb₃Sn 42 1: 975°C x 30' + 2h; 975°C x 2h



No Residual Sn traces on the sample surface

Good Nb₃Sn film superconductive properties

T_c = 16.8 K, DT_c = 0.16 K

No Sn rich Phases

Diffusion temperature to be kept above 930°C to avoid formation of low T_c phases like Nb₆Sn₅ (2.6 K) and NbSn₂ (2.1 K)

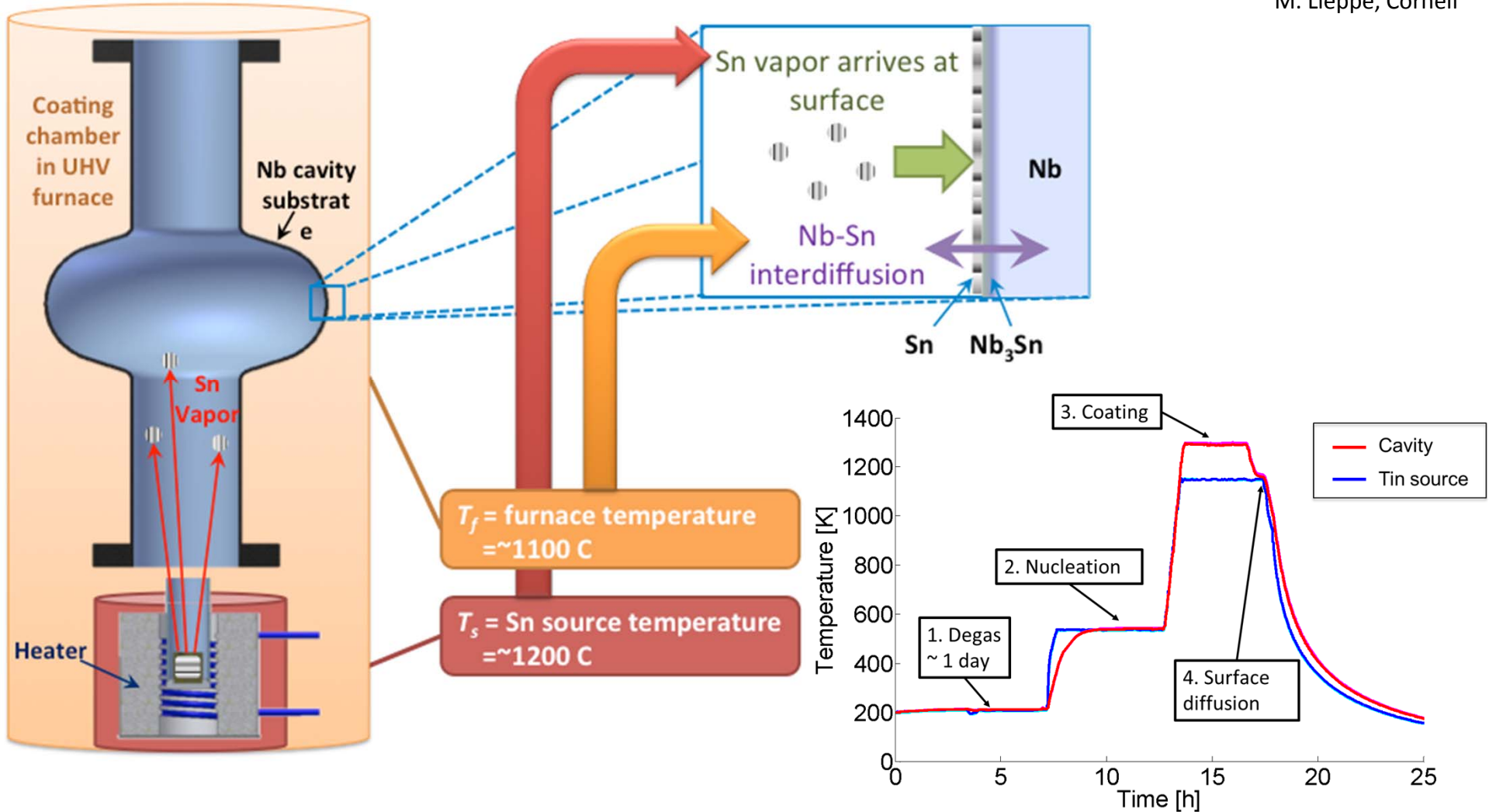
Diffusion time optimize to obtain desired Nb₃Sn thickness

Post diffusion heat reaction important to get rid of the outer Sn layer

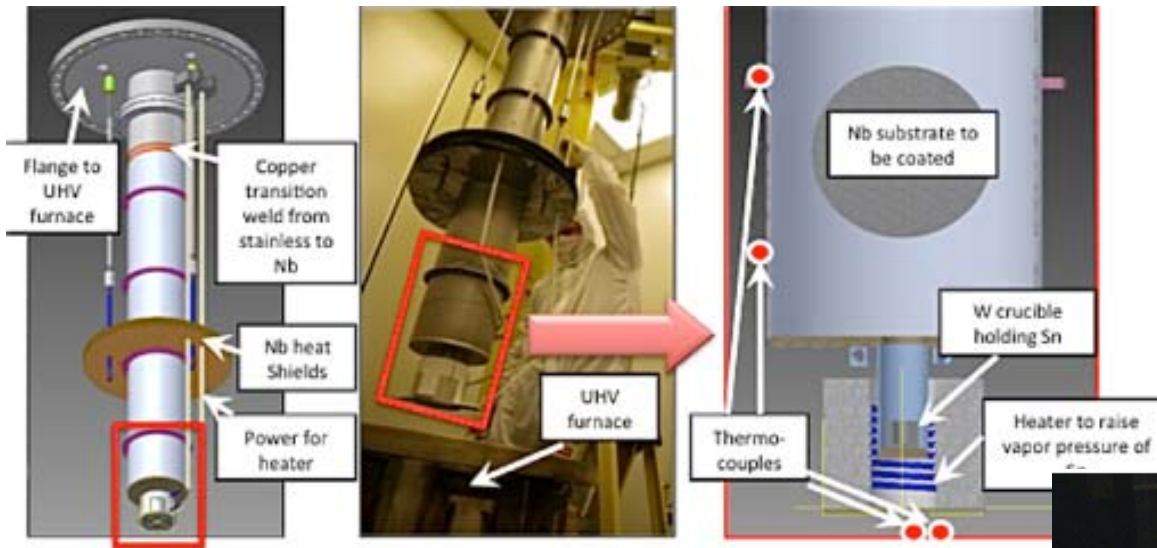
Post diffusion annealing to have enlarged grains and perfect ordering

Nb₃Sn Coating via Vapor Diffusion

M. Lieppe, Cornell

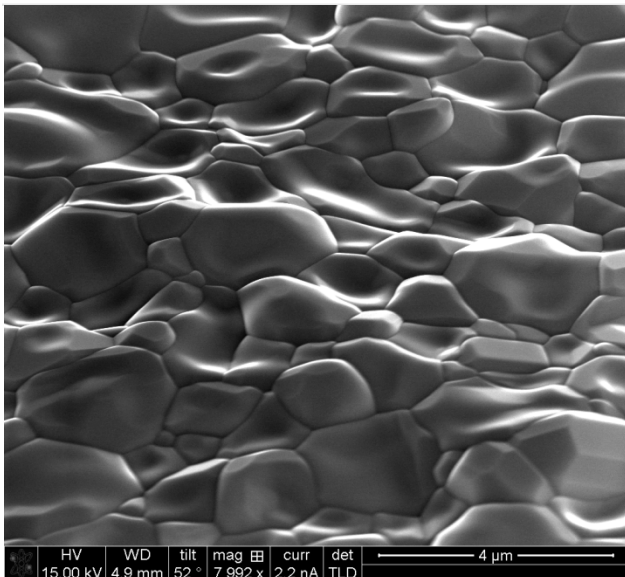


Nb₃Sn



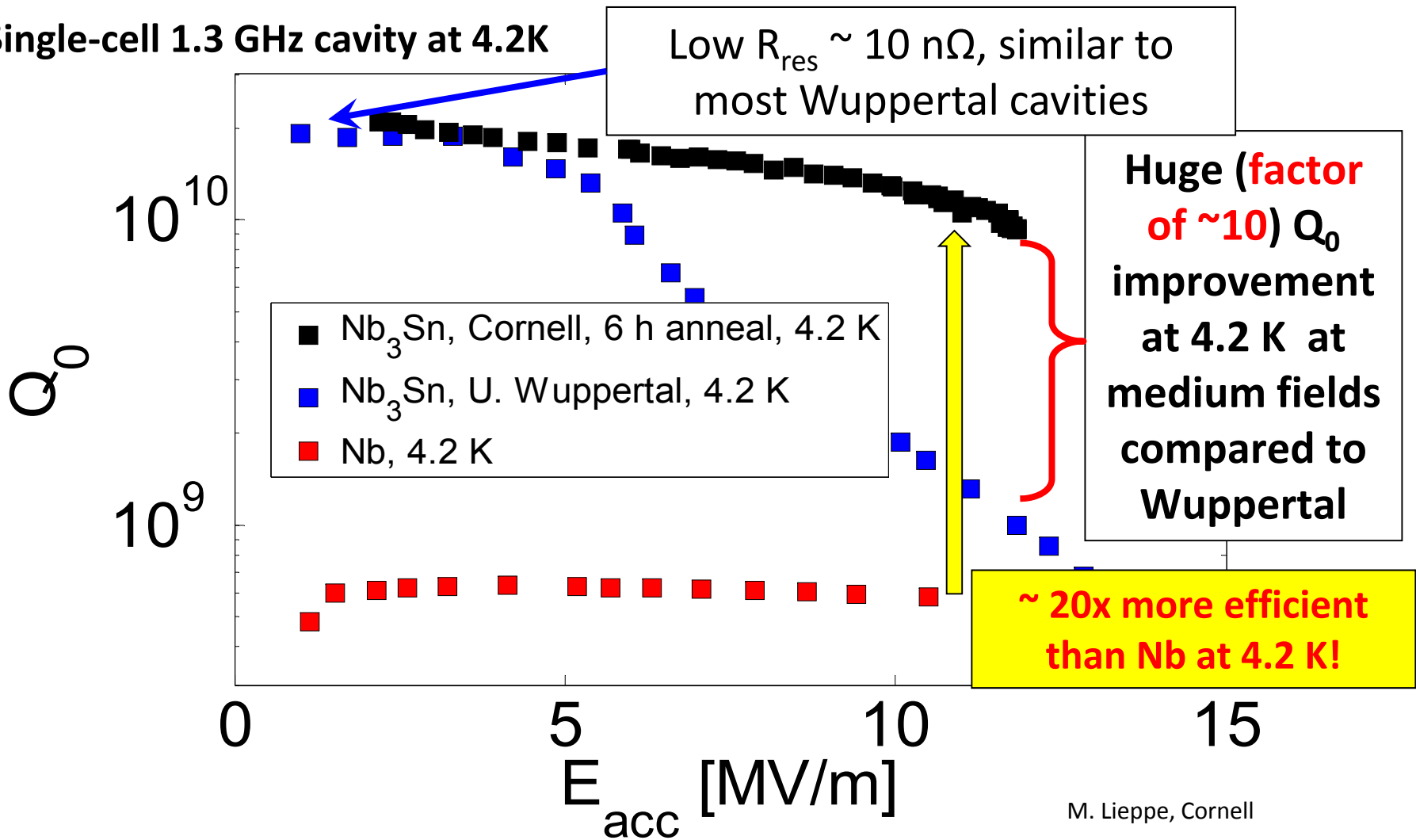
S. Posen, M. Lieppe, Cornell
G. Ereemeev, JLab

Dedicated furnace systems at
Cornell University and JLab for
Nb₃Sn synthetization via Sn
diffusion



Nb₃Sn Cavity Result

Single-cell 1.3 GHz cavity at 4.2K



A15 Compounds – V₃Si

S. Deambrosis et al., Physica C 441 (2006) 108-113

Highly ordered compound, RRR~80 achievable, max T_c (17.1K) when stoichiometric composition (25at.% Si)

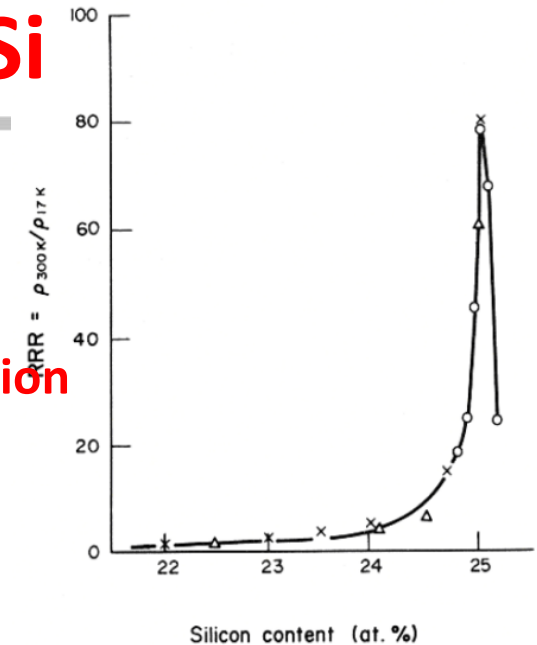
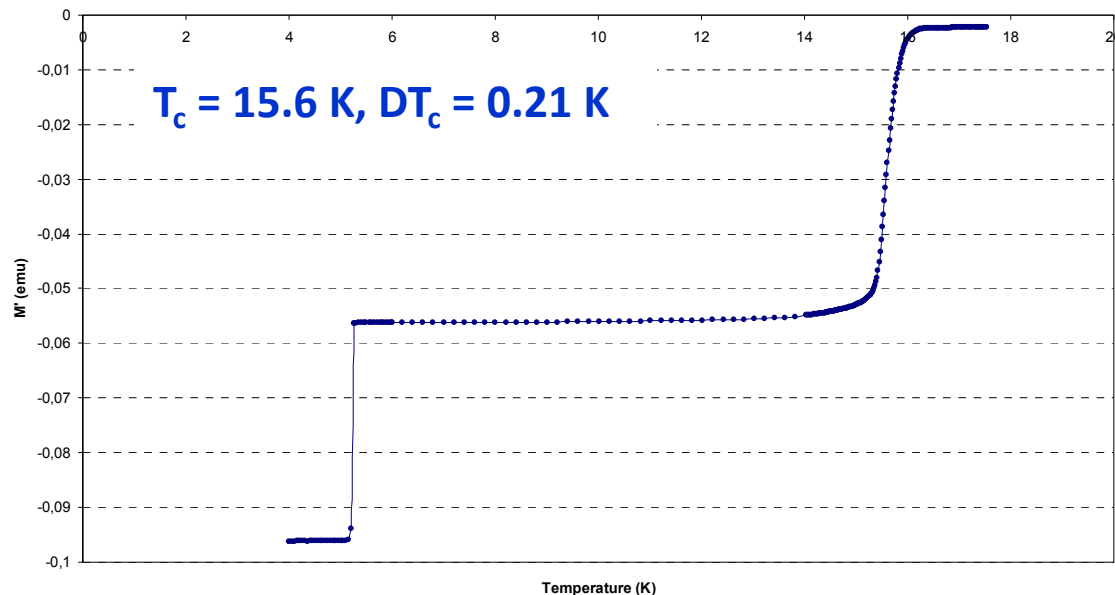
V₃Si layers by silanization of V substrate and Thermal Diffusion

V substrate heated to get SiH₄ decomposition and Silicon diffusion

Film grown by silanization with p (SiH₄) ~ 10⁻³-10⁻⁴ mbar

Annealing in vacuum to get rid of hydrogen

825°C, 4h+8h



* Diffusion parameters and silane flow rate have been optimized

* T_c ~ 16 K is routinely obtained

A15 Compounds – Mo₃Re

Mo₃Re thin films by DC magnetron deposition: Mo₇₅Re₂₅ , Mo₆₀Re₄₀

Solid solution , free of bulk and surface inhomogeneities, low interstitials solubility compared to Nb, low κ , high H_{c1} (500G)

Bulk in σ phase, tetragonal low T_c (6K)

but T_c up to 18K reported in literature with bcc structure

S.M. Deambrosis et al., Physica C 441(2006) 108-113

- * Deposition on Sapphire, Cu and Nb substrates
- * Substrate temperature up to 950° C
- * Post-annealing to increase crystallinity and transition sharpness
- * $T_c = 12K$ obtained for composition Mo₆₀Re₄₀

Higher deposition temperature,
longer annealing time

➔ Higher T_c

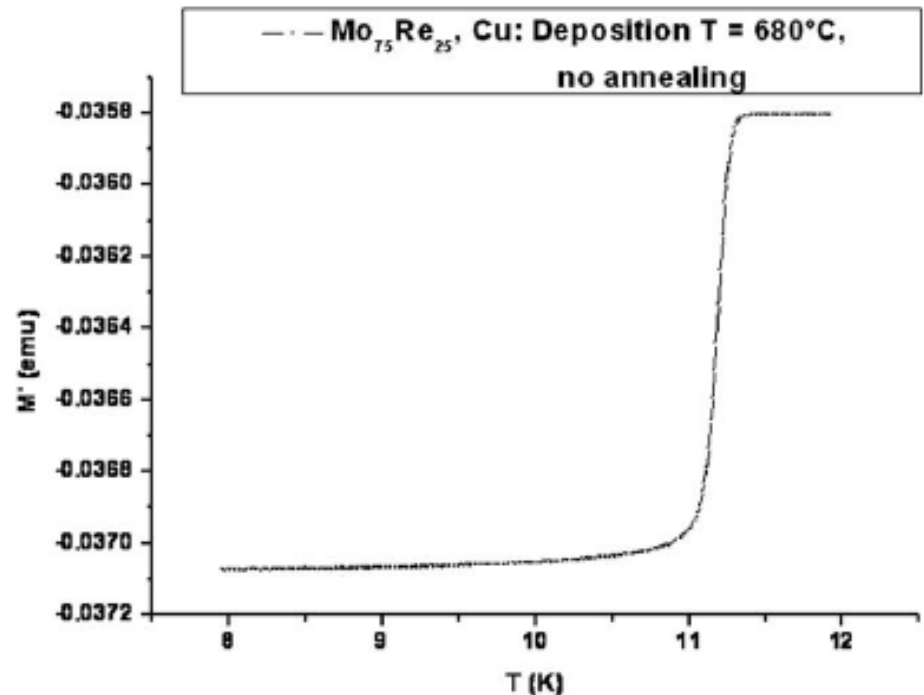


Fig. 4. A Mo₇₅Re₂₅ film deposited on Cu transition curve: deposition $T = 680^\circ\text{C}$, $T_c = 11.18$, $\Delta T_c = 0.08$ K.

MgB₂

Magnesium Diboride (MgB_2)

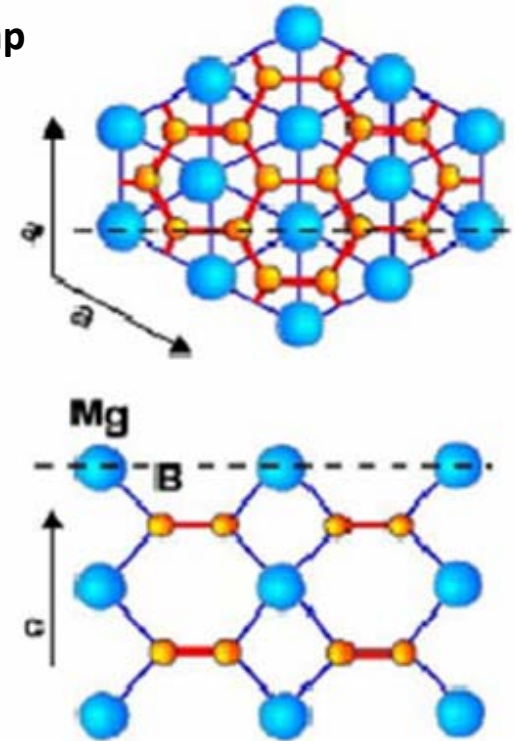
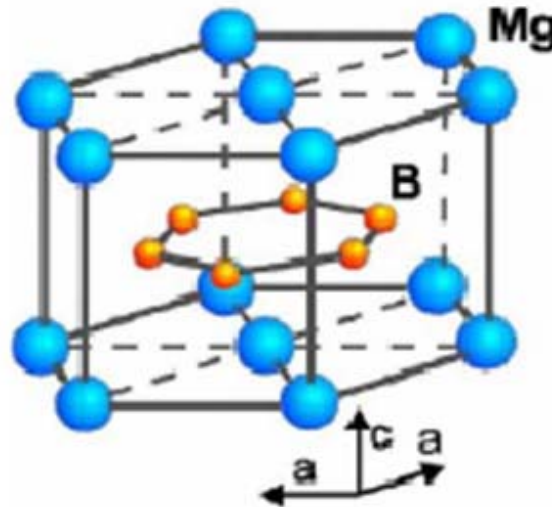
Graphite-type boron layers separated by hexagonal close-packed layers of magnesium

Superconductivity comes from the phonon-mediated Cooper pair production similar to the low-temperature superconductors except for the **two-gap nature**.

$$T_c \sim 40 \text{ K}$$

Compared to cuprates:

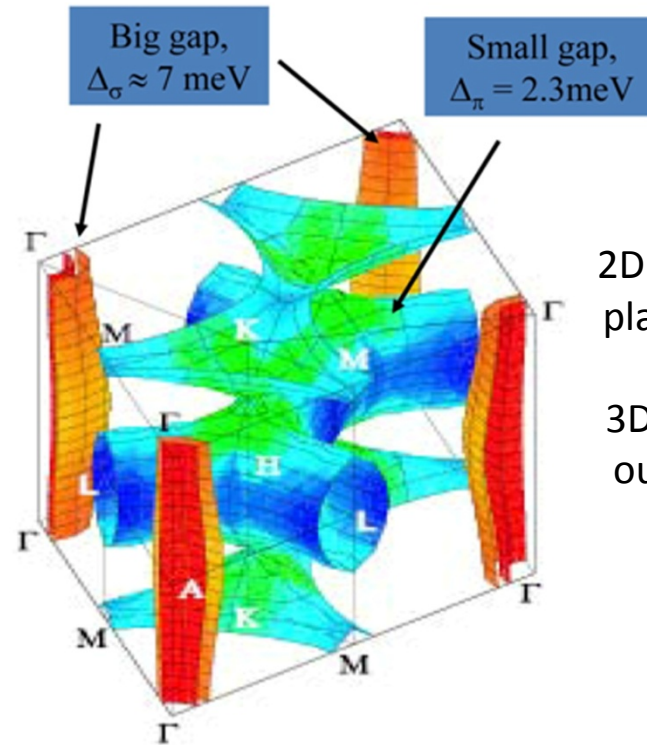
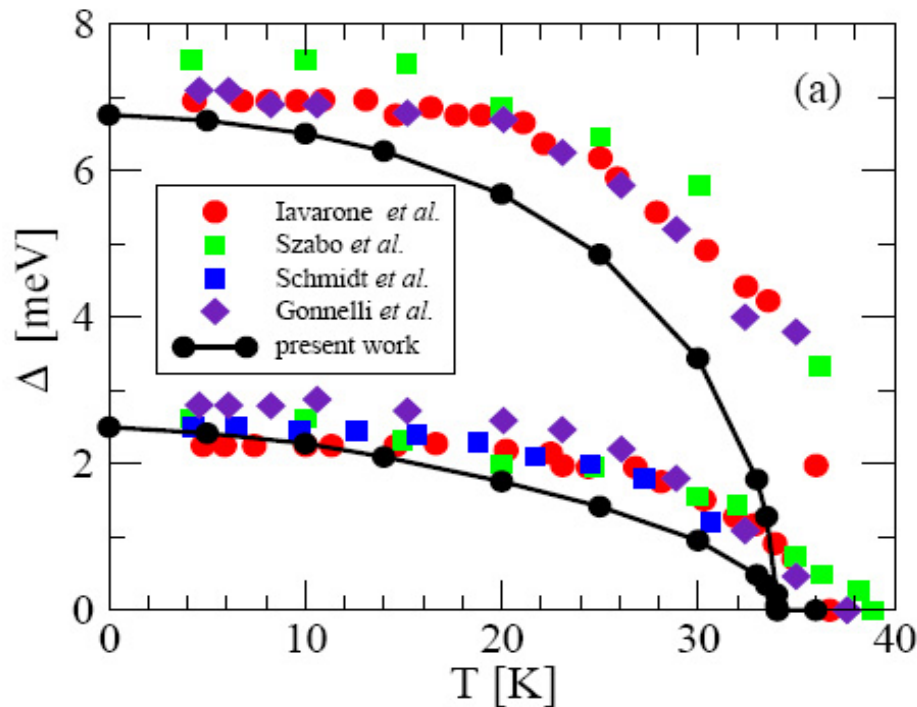
- Cheaper
- Lower anisotropy
- Larger coherence length
- Transparency of grain boundaries to current flows



C. Buzea and T. Yamashita, Superconductor Sci. Technol. 14 (2001) R115.

MgB₂: Two Energy Gaps

A. Floris et al., cond-mat/0408688v1 31 Aug 2004



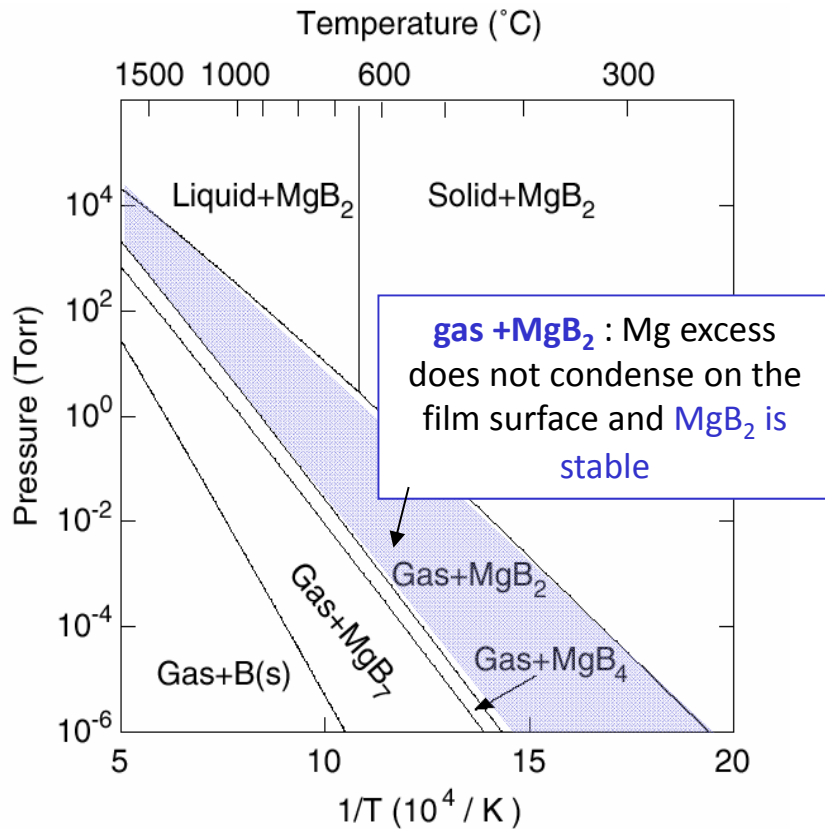
2D big gap for in-plane σ -orbitals & 3D small gap for out-of-plane π -orbitals

Liu, Mazin and Kortus (2002);
Choi et al, (2002)

R_s is dominated by the smaller gap, so MgB₂ may not be better than Nb₃Sn because $\Delta_{\pi}^{\text{MgB}_2} = 2.3 \text{ meV} < \Delta^{\text{Nb}_3\text{Sn}} = 3.1 \text{ meV}$, but better than Nb ($\Delta^{\text{Nb}} = 1.5 \text{ meV}$)

RF response has shown lower energy gap behavior. This must be compared to $\Delta \approx 1.5 \text{ meV}$ for Nb. There is room for better performance than Nb, since the resistivity can also be made quite low (best values are $\leq 1 \mu\Omega\text{cm}$).

MgB₂-Thin films growth



Z.-K. Liu et al., APL 78(2001) 3678.

evaporation Mg pressure from MgB₂ < decomposition curve of MgB₂ < Mg vapor pressure

At P=10⁻⁶ Torr and T > 250°C no accumulation of Mg will take place on the substrate and the growth of the superconducting phase is very slow due to a large kinetic energy barrier.

At low Mg pressure only extremely low deposition temperatures can be used

optimal T for epitaxial growth $\sim T_{\text{melt}}/2$

For MgB₂, 540° C → it requires P_{Mg} ~11 Torr

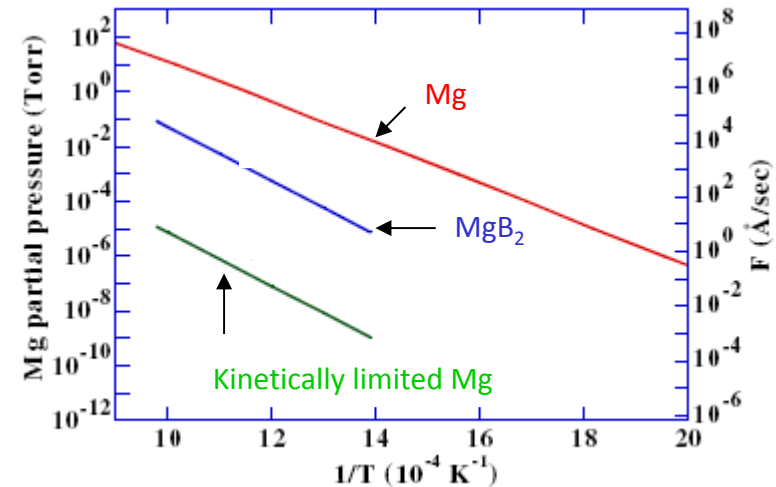
Too high for UHV deposition techniques (PLD, MBE...)

At P_{Mg} = 10⁻⁴-10⁻⁶ Torr, compatible with MBE, T_{sub} ~ 400° C

MgB₂ is stable, but no MgB₂ formation:

Mg atoms re-evaporate before reacting with B

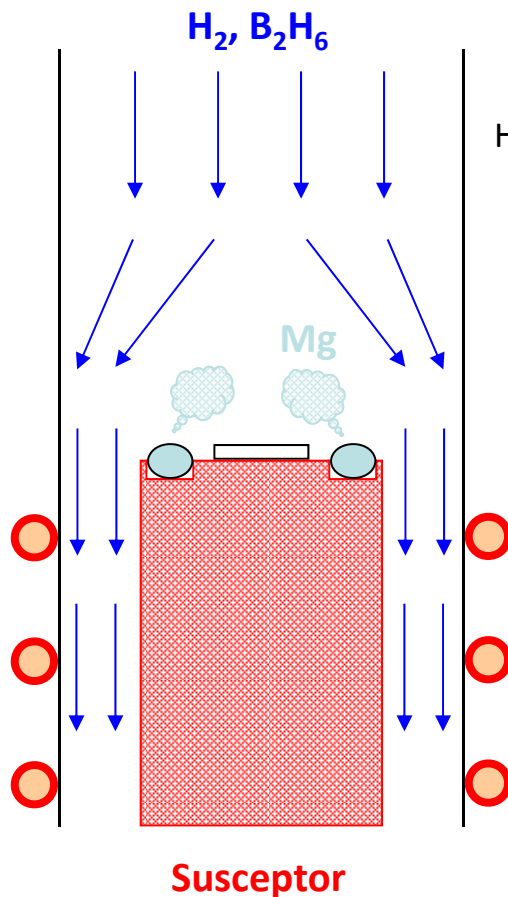
Kinetic of Mg is also important



M. Naito and K. Ueda, SUST 17 (2004) R1

MgB₂ – HPCVD on metal substrates

X. Xi- TempleUniversity



Hybrid Physical Chemical Vapor Deposition

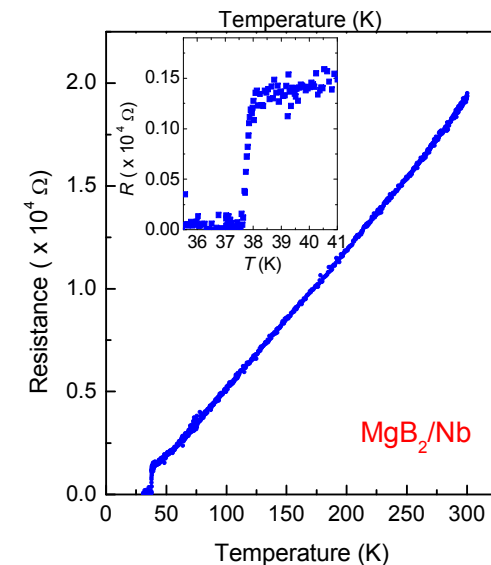
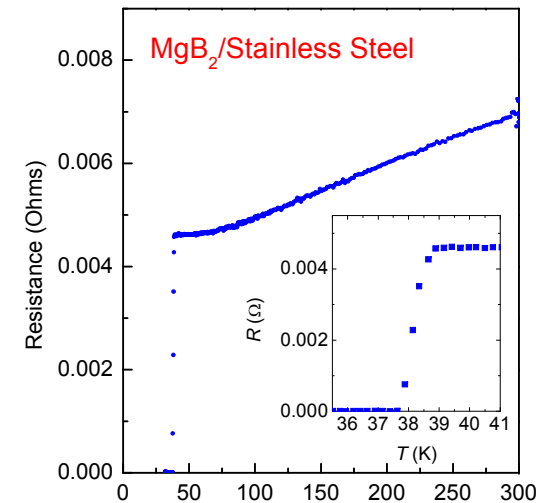
High T_c has been obtained in polycrystalline MgB₂ films on stainless steel, Nb, TiN, and other substrates.

Clean HPCVD MgB₂ thin films with excellent properties:

- RRR > 80
- low resistivity (< 0.1 $\mu\Omega$) and long mean free path
- high $T_c \sim 42$ K (due to tensile strain), high J_c (10% depairing current)
- low surface resistance, short penetration depth
- smooth surface (RMS roughness < 10 Å with N₂ addition)
- good thermal conductivity (free from dendritic magnetic instability)

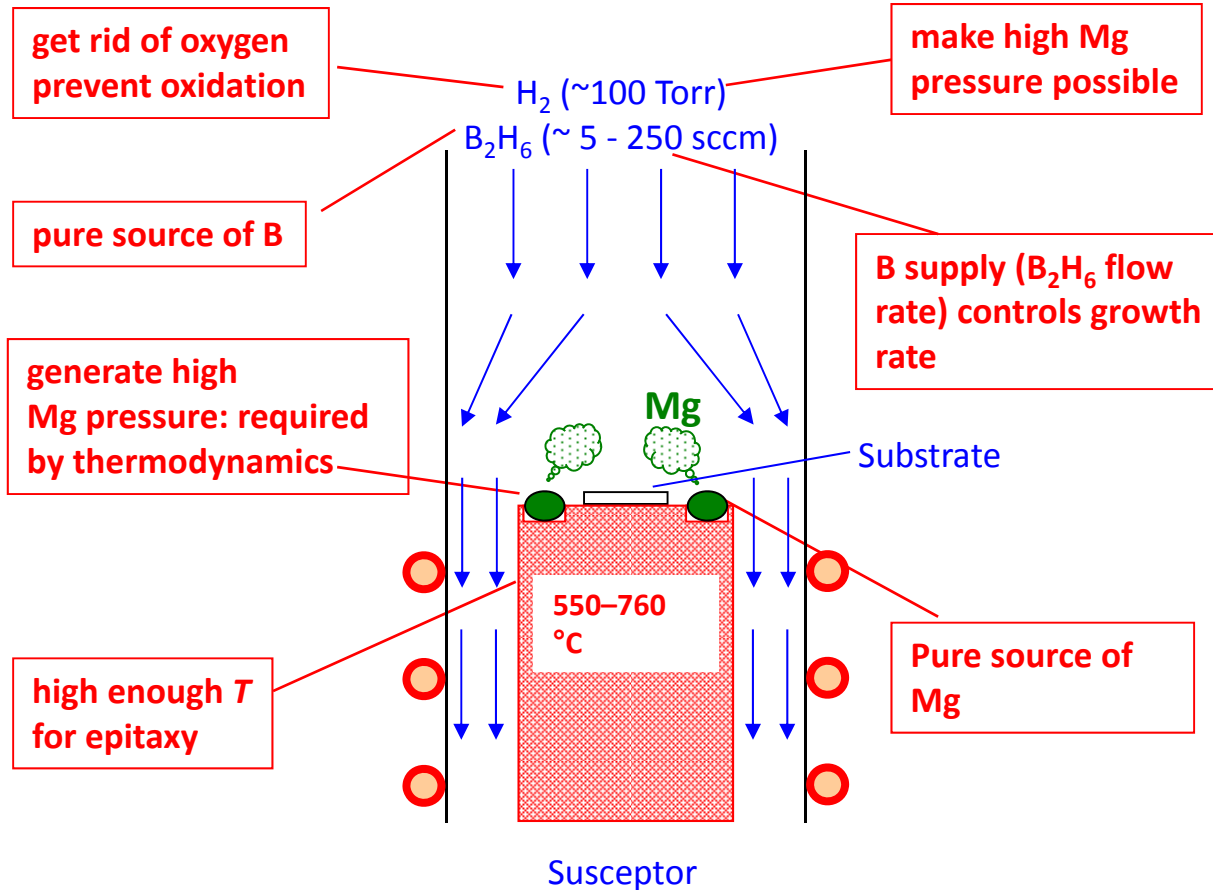
Critical engineering considerations:

- generate high Mg pressure at substrate (cold surface is Mg trap)
- deliver di-borane to the substrate (the first hot surface di-borane sees should be the substrate)



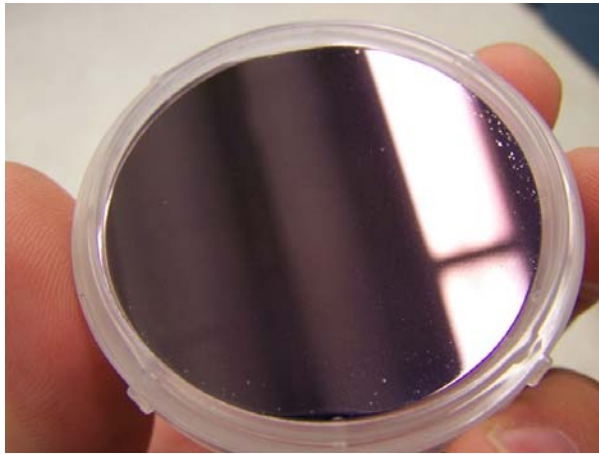
Hybrid Physical Chemical Vapor Deposition (HPCVD) at Temple University

Schematic View

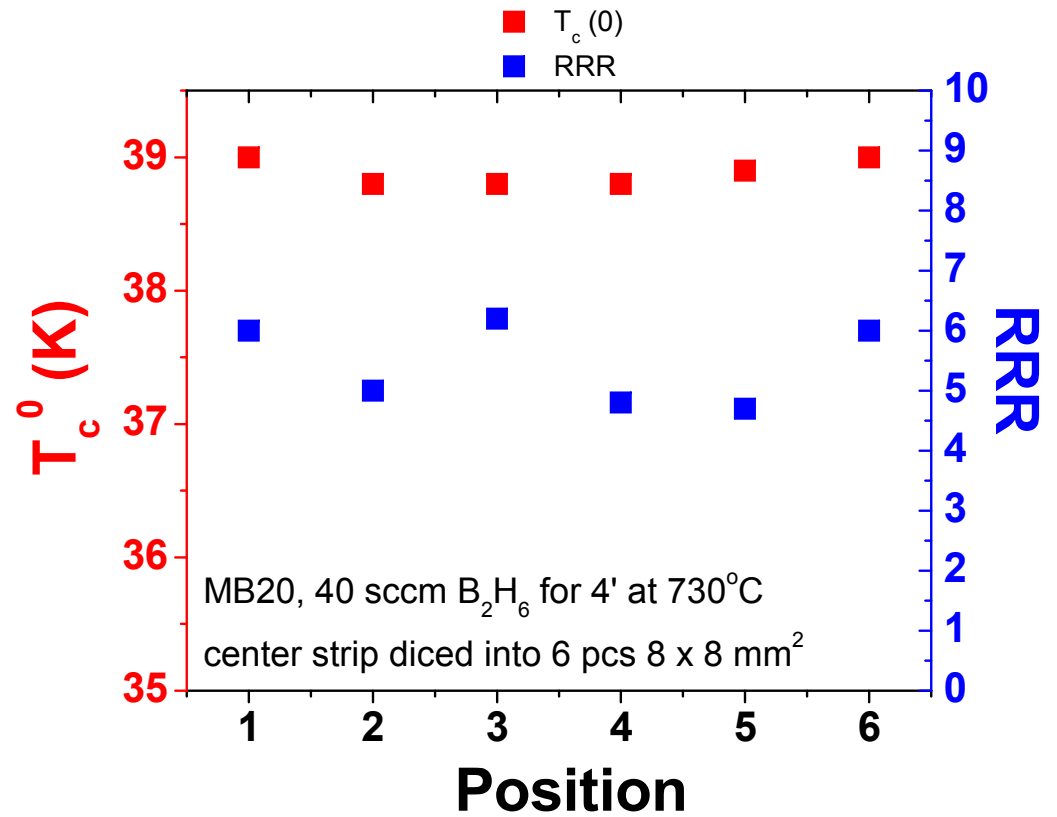
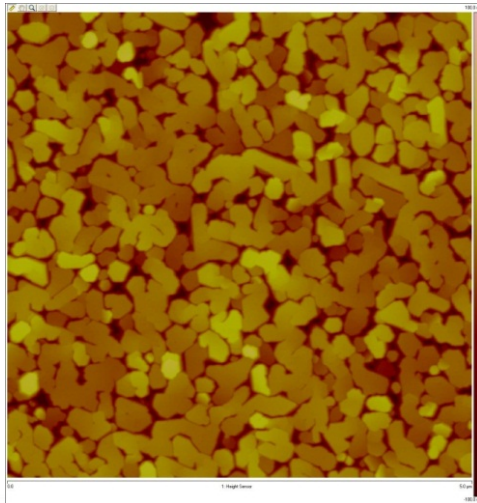


2" MgB₂ Films Grown by HPCVD

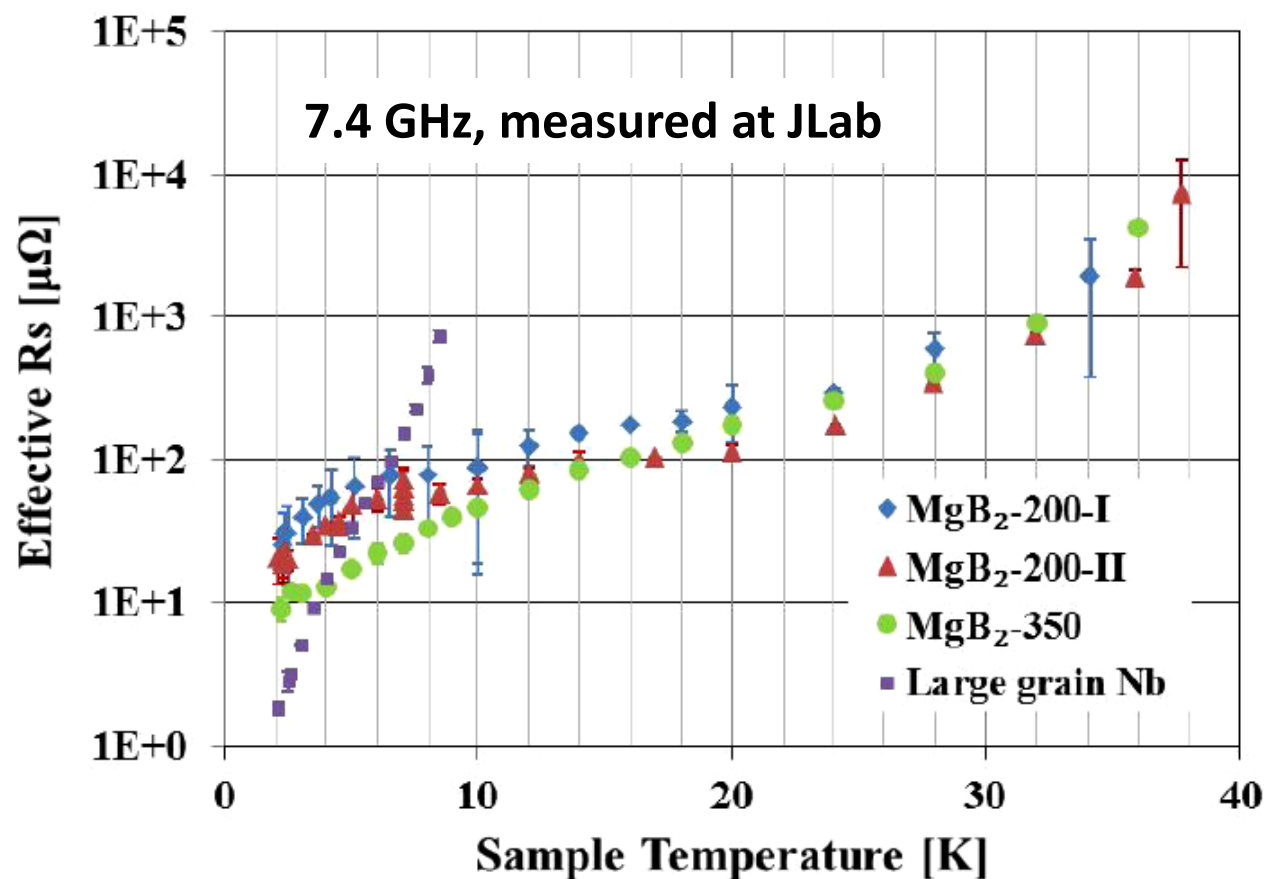
200 nm 2" MgB₂ film on sapphire



AFM image of film surface



Surface Resistance Compared to Large Grain Nb



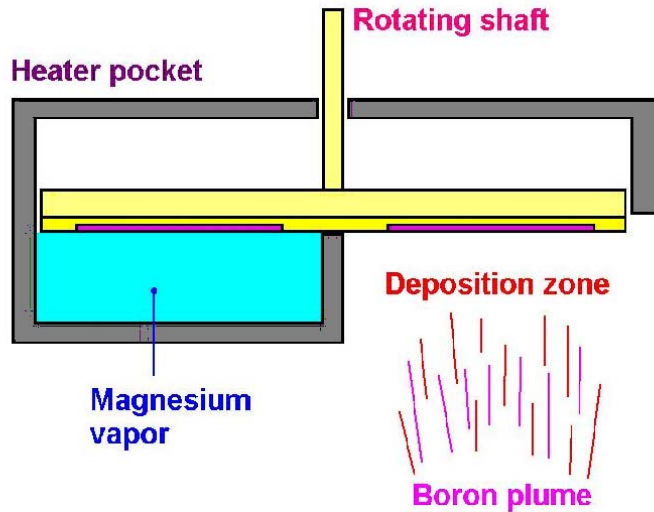
Surface resistance of 2" dia. 350 nm MgB₂ film on sapphire comparable to the best large grain Nb at 4 K.

MgB₂ - Reactive Evaporation

T. Tajima, LANL

RF measurement @ MIT/Lincoln Lab

Superconducting Technologies Inc.



In-situ reactive evaporation @ 550°C

Compared to Nb:

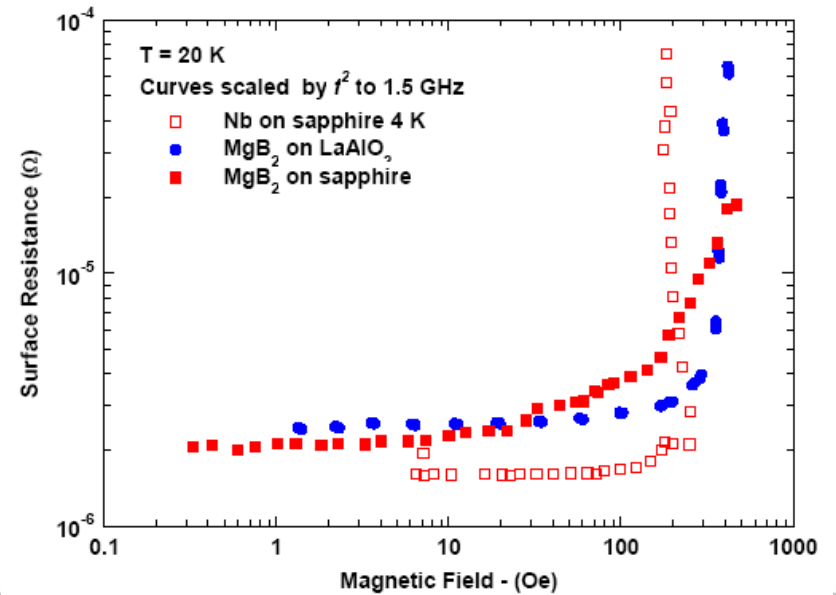
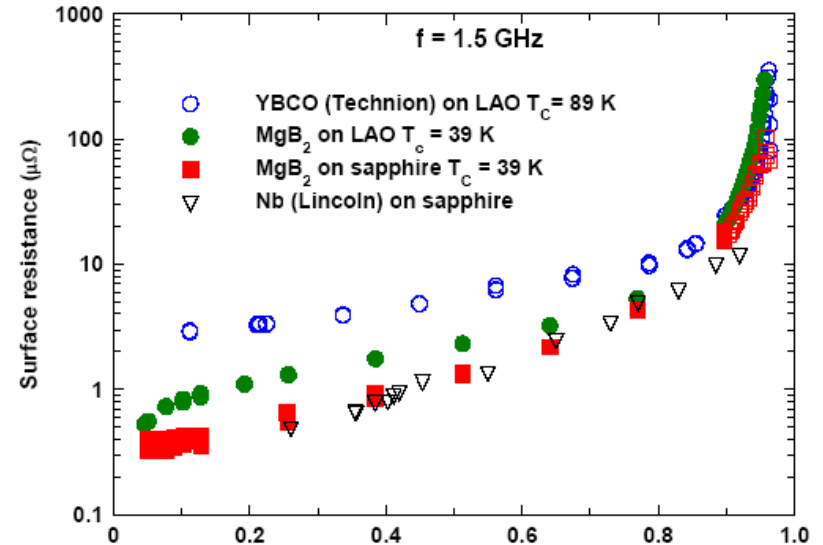
- Higher T_c
- low resistivity
- larger gap
- higher critical field

B.H. Moeckly et al., *IEEE Trans. Appl. Supercond.* 15 (2005) 3308.

T. Tajima et al, *Proc. PAC05*.

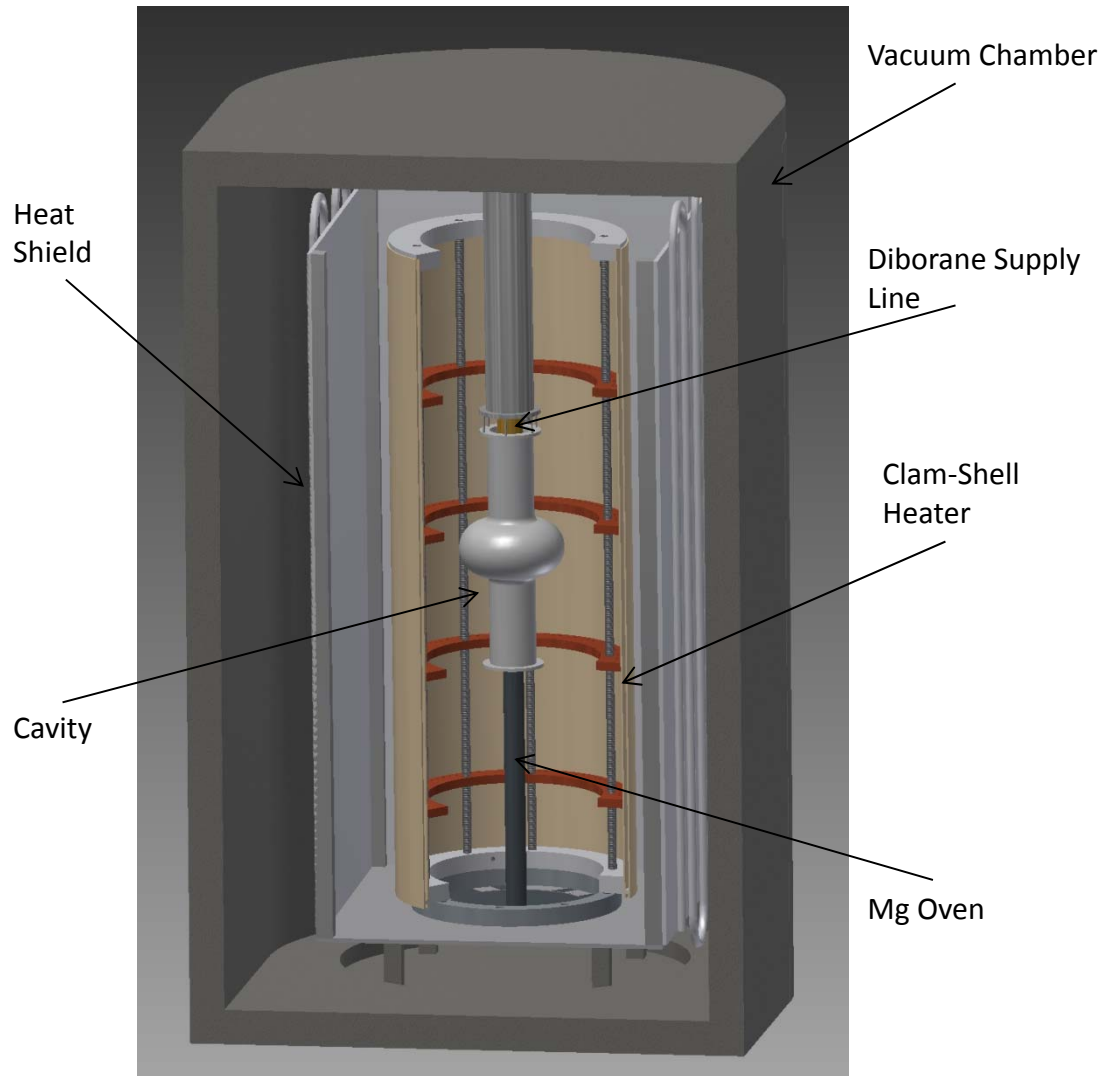
B.H. Moeckly, *ONR Superconducting Electronics Program Review*
Red Bank, NJ, February 8, 2005

Oates, Agassi, and Moeckly, *ASC 2006 Proceeding*



MgB₂ – An Idea for SRF Cavity Coating

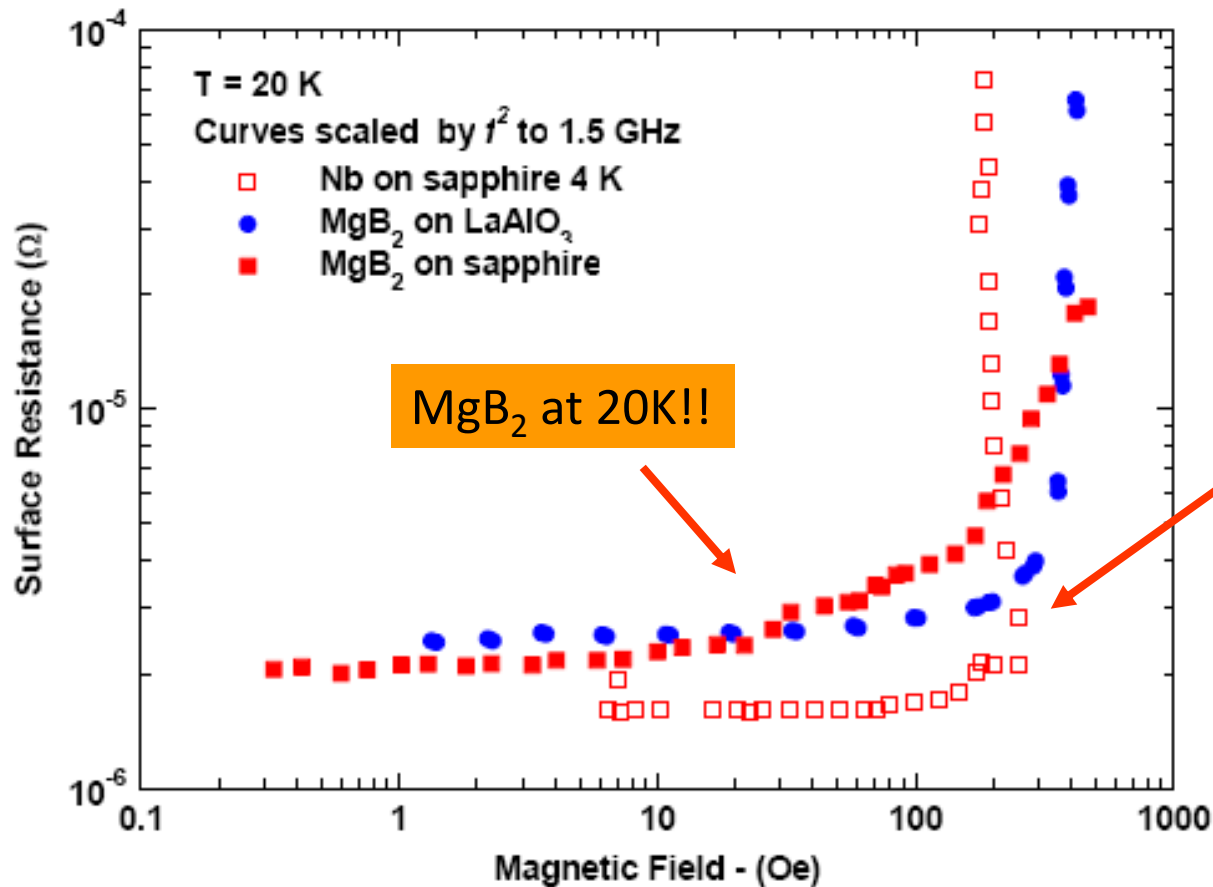
In-Situ Coating of MgB₂: 3 GHz Cavity



- Need enough space between diborane supply line/Mg oven and cavity tubes
- Better control of gas flow inside the cavity
- Largest size that can be accommodated by the existing vacuum chamber

R_s for MgB_2

Results at MIT in 2006 showed R_s comparable to Nb even at 20 K! and the field at which the R_s start to increase rapidly was higher than Nb!!



Coated at STI with reactive evaporation
Thickness: 500 nm

Oates et al., IEEE Trans. Appl. Supercond. 17 (2007) 2871.

MgB₂ – Challenges

Keys to high quality MgB₂ thin films:

- High Mg pressure for thermodynamic stability of MgB₂
- oxygen-free or reducing environment
- clean Mg and B sources

Challenges

Film properties degrade with exposure to moisture: resistance goes up, T_c goes down

Clean cavity surface leads to degradation in water and moisture

... need of a cap layer?

Diffusion of Mg into substrate (Cu...), need of a buffer layer

Safety ... procedures for use of diborane

OXYPNICTIDES

A new superconducting family

Oxypnictide base

ReOMP_n

- M = Fe, Co, Ni
- Pn = As or P
- Re = La, Nd, Sm, Pr

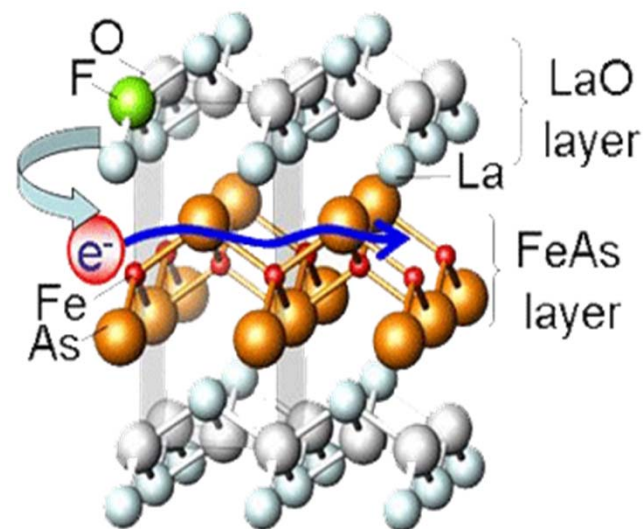
Layered as HTS –superconducting
AsFe layers and T_c from <10K to 55K

Superconductivity occurs on the FeAs layer
with magnetic pair-breaking Fe²⁺ions

Another families: Ba_{1-x}K_xFe₂As₂

Up to a few thousand compounds

Materials usually brittle



Iron-Based Layered Superconductor La[O_{1-x}F_x]FeAs (x = 0.05–0.12)
with T_c = 26 K

Yoichi Kamihara,*† Takumi Watanabe,‡ Masahiro Hirano,^{1,§} and Hideo Hosono^{1,¶§}

ERATO-SORST, JST, Frontier Research Center, Tokyo Institute of Technology, Mail Box S2-13, Materials and Structures Laboratory, Tokyo Institute of Technology, Mail Box R3-1, and Frontier Research Center, Tokyo Institute of Technology, Mail Box S2-13, 4259 Nagatsuta, Midori-ku, Yokohama 226-8503, Japan

Received January 9, 2008; E-mail: hcsono@mssl.titech.ac.jp

Tests in magnetic fields up to 45 T suggest the H_{c2} of
LaFeAsO_{0.89}F_{0.11} may be ~64 T.
A different La-based material (La_{0.8}K_{0.2}FeAsO_{0.8}F_{0.2})
tested at 6 K predicts H_{c2} ~122 T.

Can pnictides be useful for TFSRF?

For s-wave members of the pnictide family, one can expect a much lower R_s at 2K because

$$\Delta_{\text{oxy}} = 5\text{-}10 \text{ meV} > \Delta_{\text{Nb}_3\text{Sn}} = 3 \text{ meV}$$

Normal skin effect ($l \ll \lambda$): multiple impurity scattering in the λ -belt:

$$R_s \sim (\mu_0^2 \omega^2 \lambda^3 \sigma_n \Delta / T) \exp(-\Delta / T)$$

Anomalous skin effect ($l \gg \lambda$) in the clean limit:

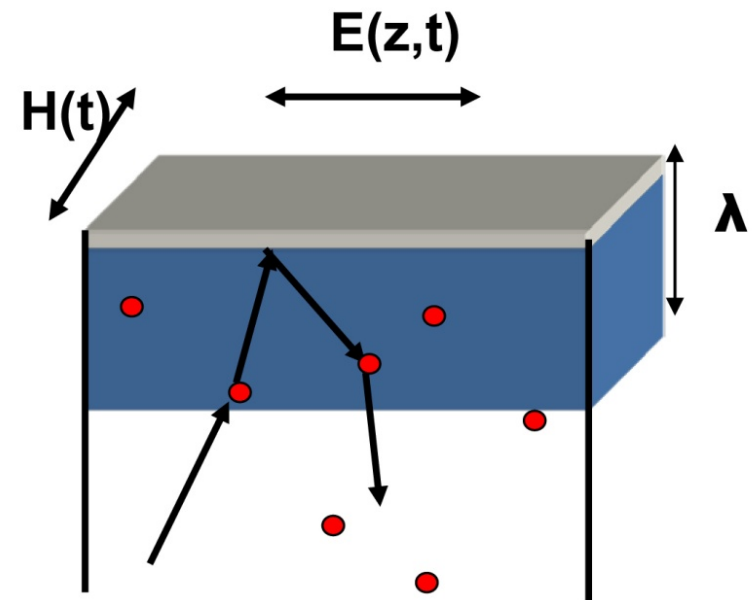
$$\text{Effective } \sigma_{\text{eff}} \sim e^2 n \lambda / p_F$$

High $\rho_n \sim 1 \text{ m}\Omega\text{cm} \sim 10 \rho_n \text{ MgB}_2$, **big $\lambda = 180\text{-}250 \text{ nm}$** ,

low $H_{c1} \sim 10 \text{ mT}$

R_s can be much lower than R_s of Nb_3Sn , but TF multilayer coating is necessary

First high quality epitaxial films have been grown by C.-B. Eom (UW)

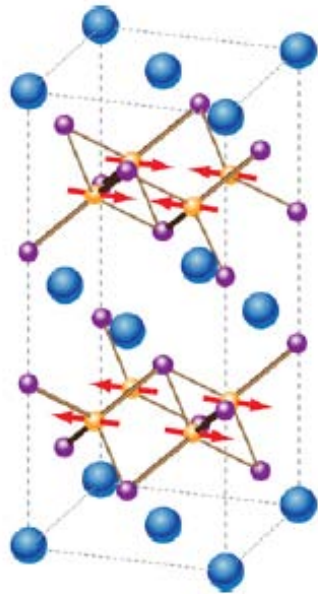


Pnictides Films Deposition

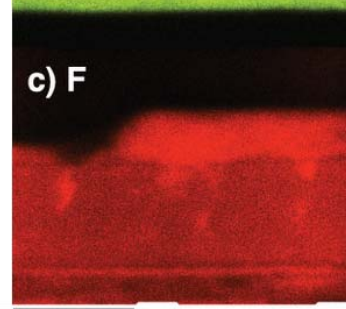
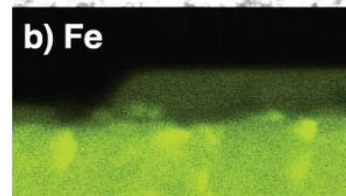
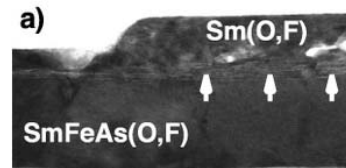
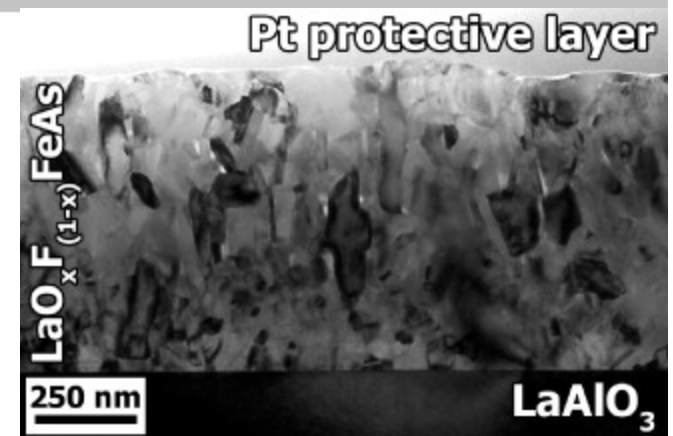
Coating Methods:

Pulse Laser Deposition from stoichiometric target

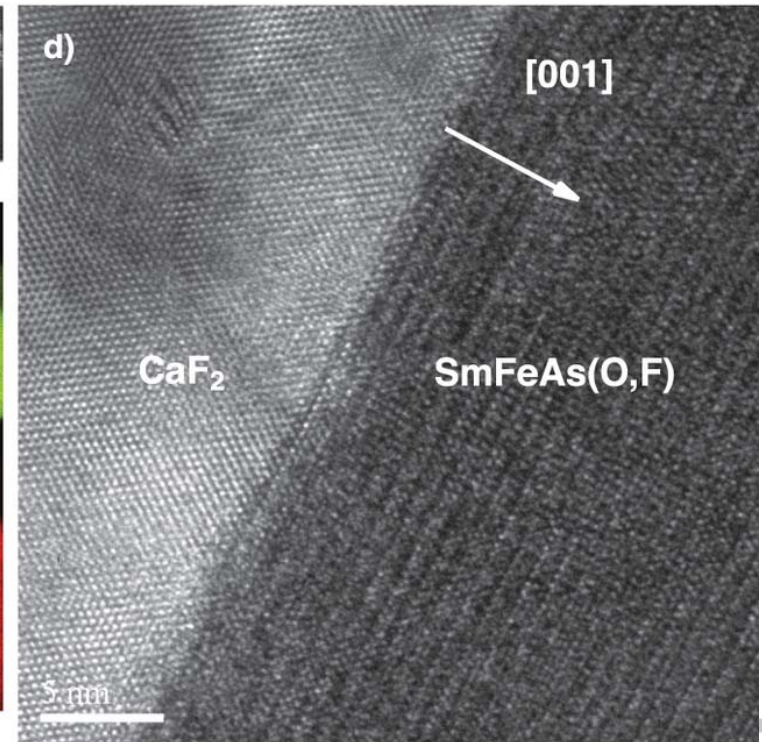
Ion Beam Assisted - MBE



First high quality epitaxial films have been grown by C.-B. Eom (UW)



100 μm

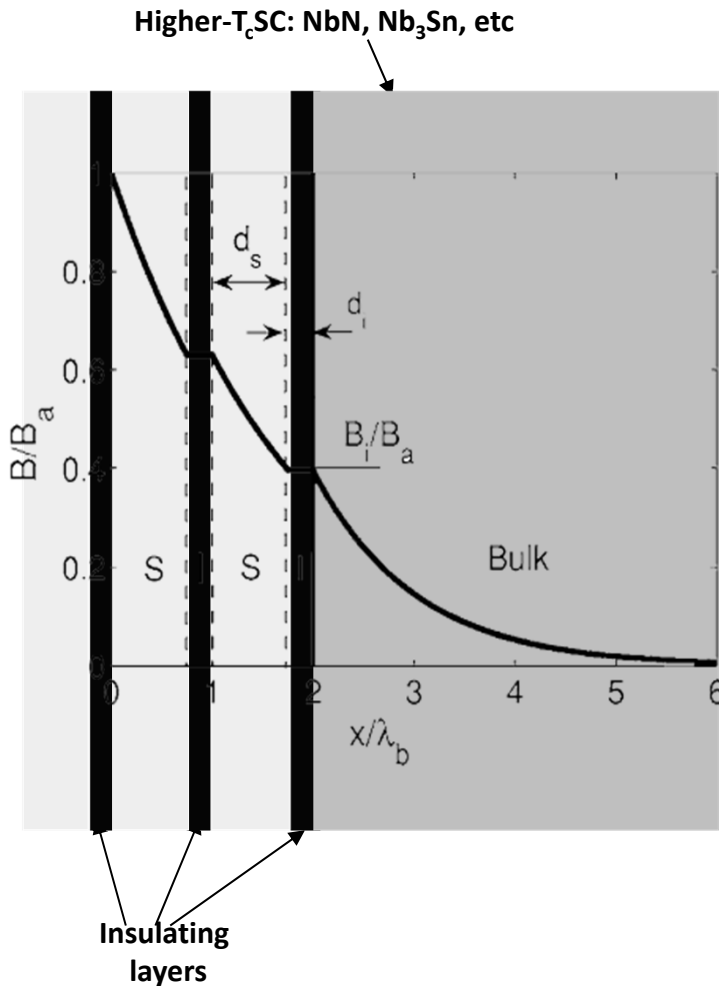


SIS MULTILAYERS

Beyond Nb: SIS Multilayers

Taking advantage of the high T_c superconductors with much higher H_c without being penalized by their lower H_{c1} ...

Alex Gurevich, *Appl. Phys. Lett.* 88, 012511 (2006)



Multilayer coating of SC cavities:

alternating SC and insulating layers with $d < \lambda$

Higher T_c thin layers provide magnetic screening of the Nb SC cavity (bulk or thick film) without vortex penetration

No thermodynamically stable parallel vortices due to the enhancement of H_{c1} in thin films with $d < \lambda$ (Abrikosov, 1964)

$$H_{c1} = \frac{2\phi_0}{\pi d^2} \left(\ln \frac{d}{\xi} - 0.07 \right)$$

The breakdown field could be increased up to the superheating field H_s of the coating: 450 mT for Nb₃Sn

- Strong reduction of BCS resistance because of using SC layers with higher Δ (Nb₃Sn, NbN, etc)

Pushes the accelerating gradient above 100 mV/m

- Possibility to move operation from 2K to 4.2K

Screening field in a multilayer

- Solutions of London equation for a layer with the penetration depth λ on a substrate with the penetration depth λ_0

$$h_1(x) = H[(1 - c)e^{-x/\lambda} + ce^{x/\lambda}], \quad 0 < x < d$$

$$h_2(x) = Hbe^{(d-x)/\lambda_0}, \quad x > d$$

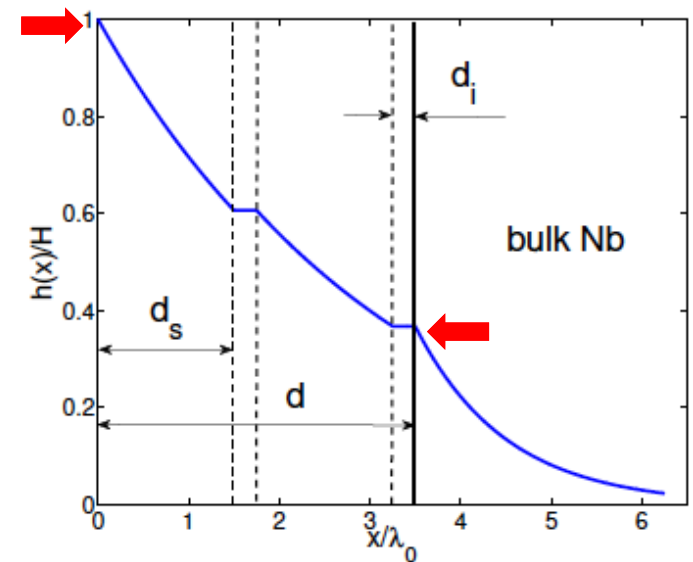
where c and b are given by:

$$c = \frac{k}{k + g^2}, \quad b = \frac{(1 + k)g}{k + g^2}$$

- Important parameters

$$k = \frac{\lambda - \lambda_0}{\lambda + \lambda_0}, \quad g = \exp(d/\lambda)$$

T. Kubo, Y. Iwashita, and T. Saeki, *APL* 104, 032603 (2014);
A. Gurevich, *AIP Advances*, 5, 017112 (2015)



Meissner state breaks down at the surface of either ML or Nb where the current densities $J(0) = h'(0)$ and $J(d) = h'(d)$ are maximum
for the SC substrate (Nb) with $\lambda_0 < \lambda$, both c and k are positive

Current counterflow induced by the Nb substrate

- Current density in the layer $J(x) = -h'(x)$:

$$J(x) = [(1 - c)e^{-x/\lambda} - ce^{x/\lambda}]H/\lambda$$

- Current density at the surface $J(0)$ is reduced by the substrate with $\lambda_0 < \lambda$:

$$J(0) = \left[1 - \frac{2k}{k + \exp(2d/\lambda)} \right] \frac{H}{\lambda} < J_d = \frac{H_s}{\lambda}$$

Counterflow induced by the substrate reduces the current density at the ML surface, allowing the Meissner state in the ML to survive up to fields exceeding the bulk superheating field H_s

For a thick ML with $d \gg \lambda$, the maximum field H_m is limited by H_s :
at optimum thickness d_m the field H_m exceeds both H_s and H_{s0}

Optimum thickness

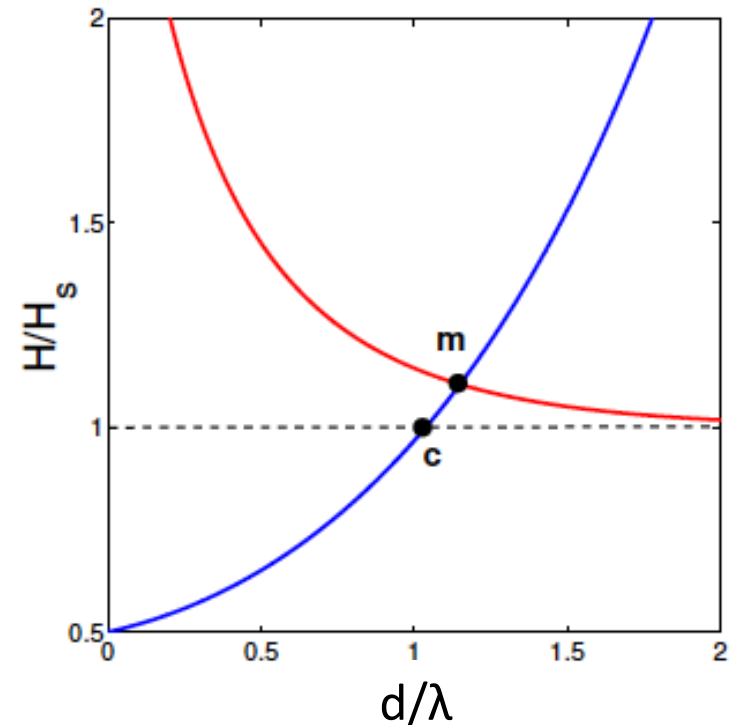
- The Meissner state is stable if the screening current density at the surface of both the ML and the substrate is smaller than the depairing limit:

$$J(0) < J_d = H_s/\lambda \text{ and } J(d) < J_{d0} = H_{s0}/\lambda_0$$

$$\frac{(e^{2d/\lambda} - k)H}{e^{2d/\lambda} + k} \leq H_s, \quad \frac{H(1+k)e^{d/\lambda}}{e^{2d/\lambda} + k} \leq H_{s0}$$

The Meissner state is below both blue and red lines.

The crossing point defines the optimum thickness d_m for **maximum H_m which exceeds the superheating fields of both the layers and the substrate**



$$d_m = \lambda \ln(\mu + \sqrt{\mu^2 + k}), \quad \mu = \frac{H_s \lambda}{(\lambda + \lambda_0) H_{s0}}$$

for $H_s = 2H_{s0}$ and $k = 1/2$,
 $d_c = \ln[\mu + (\mu^2 - k)^{1/2}]$

Maximum screening field

- The maximum screening field H_m corresponds to $d = d_m$ for which

$$H_m = \left[H_s^2 + \left(1 - \frac{\lambda_0^2}{\lambda^2} \right) H_{s0}^2 \right]^{1/2}$$

H_m at the optimum thickness **exceeds the bulk superheating fields of both Nb and the layer material**. For $\lambda \gg \lambda_0$, practically for $\lambda > 160$ nm for a SC layer on the Nb cavity with $\lambda_0 = 40$ nm, H_m approaches the limit

$$H_m \rightarrow \sqrt{H_s^2 + H_{s0}^2}$$

Let us evaluate H_m for a ML on clean Nb with $\lambda_0 = 40$ nm and $H_{s0} = 1.2H_c = 240$ mT (the GL result for clean Nb) and different layer materials, such as Nb₃Sn, NbN, pnictides, **and also dirty Nb**

A. Gurevich, *AIP Advances*, 5, 017112 (2015)

Estimates of H_m and d_m

- **Nb₃Sn:** $H_s = 0.84H_c = 454$ mT and $\lambda = 120$ nm (moderately dirty):

$$H_m = 507 \text{ mT}, \quad d_m = 1.1\lambda = 132 \text{ nm}$$

doubles the superheating field of clean Nb

- **Ba_{0.6}K_{0.4}Fe₂As₂,** $T_c = 38$ K, $H_c = 0.9$ T, $H_s = 756$ mT, $\lambda = 200$ nm

$$H_m = 930 \text{ mT}, \quad d_m = 1.78\lambda = 356 \text{ nm.}$$

almost quadruples the superheating field of clean Nb

- **dirty Nb layer:** $H_c = 200$ mT, $H_s = 170$ mT, $l = 2$ nm, and $\lambda = \lambda(\xi_0 / l)^{1/2} = 180$ nm

$$H_m = 288 \text{ mT}, \quad d_m = 0.44\lambda = 79 \text{ nm.}$$

20% gain as compared to $H_s = 240$ mT of clean Nb

SIS Multilayers – The Benefits

- Multilayer S-I-S-I-S coating could make it possible to take advantage of superconductors with much higher H_c , than those for Nb without the penalty of lower H_{c1}
- Strong increase of H_{c1} in films allows using rf fields $> H_c$ of Nb, but lower than those at which flux penetration in grain boundaries may become a problem
- Strong reduction of BCS resistance because of using SC layers with higher Δ (Nb₃Sn, NbN, ...)

The significant performance gain may justify the extra cost.

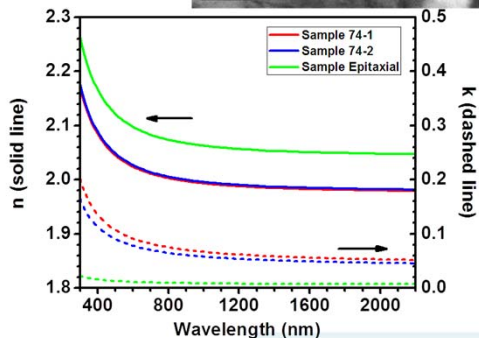
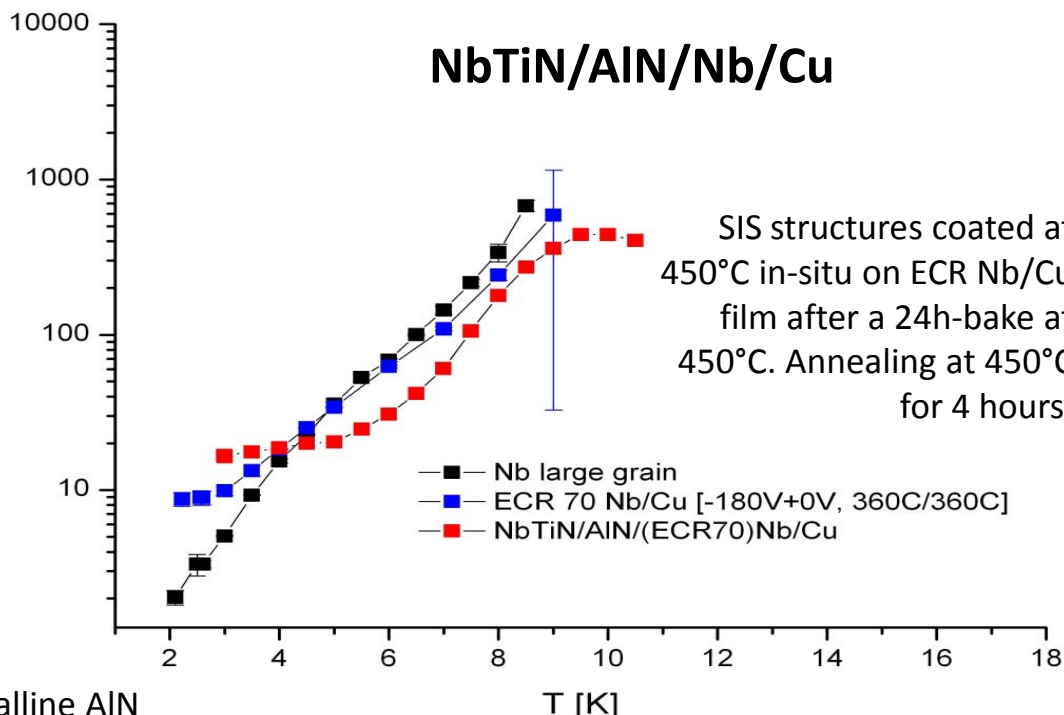
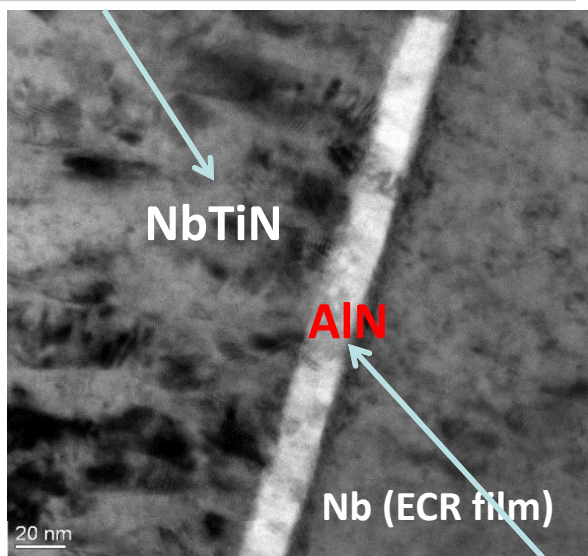
... but ...

Technical challenges, influence of composition on H_{c1} and H_c , influence of the morphology and composition at grain boundaries,

...

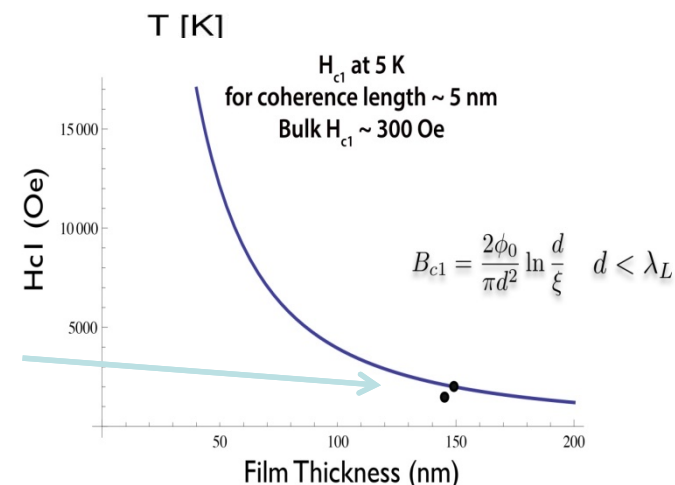
SIS NbTiN/AlN structures on Nb surfaces

150 nm polycrystalline NbTiN
T_c=16.9K



20 nm polycrystalline AlN
n=2.02
(range in Lit. n = 1.9-2.1)

	Thickness [nm]	H _{c1} [mT]	T _c [K]
NbTiN/MgO	2000	30	17.25
NbTiN/AlN/AlN ceramic	145	135	14.84
NbTiN/AlN/MgO	148	200	16.66



Lower BCS resistance beyond 4 K for SIS coated Nb/Cu film compared to standard TiM & bulk Nb. Similar effect observed for NbTiN/AlN/bulk Nb.

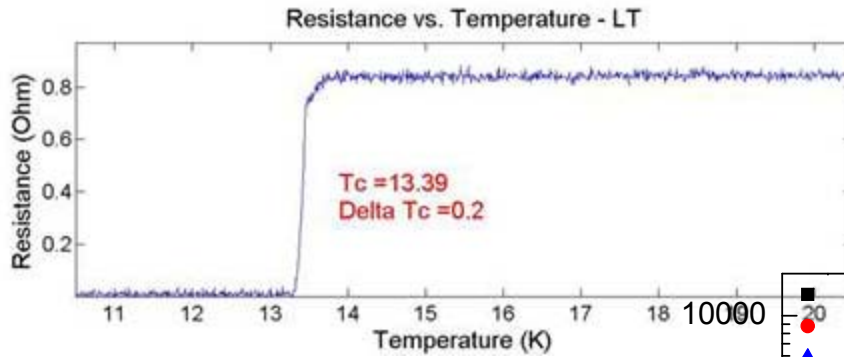
50 nm NbN

15 nm MgO

250 nm Nb

MgO (100)

NbN based SIS structures



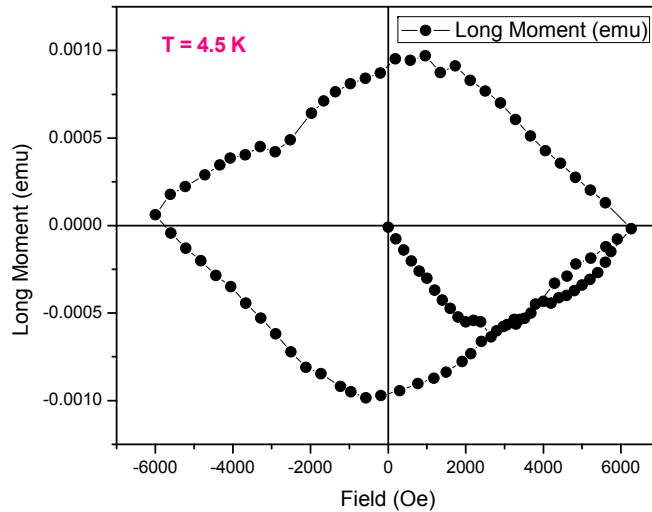
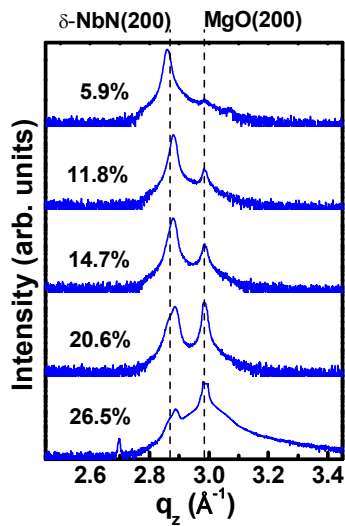
College William & Mary (Prof. Lukaszew)

CEA Saclay (C. Antoine)

Standalone NbN films $T_c \sim 15$ K

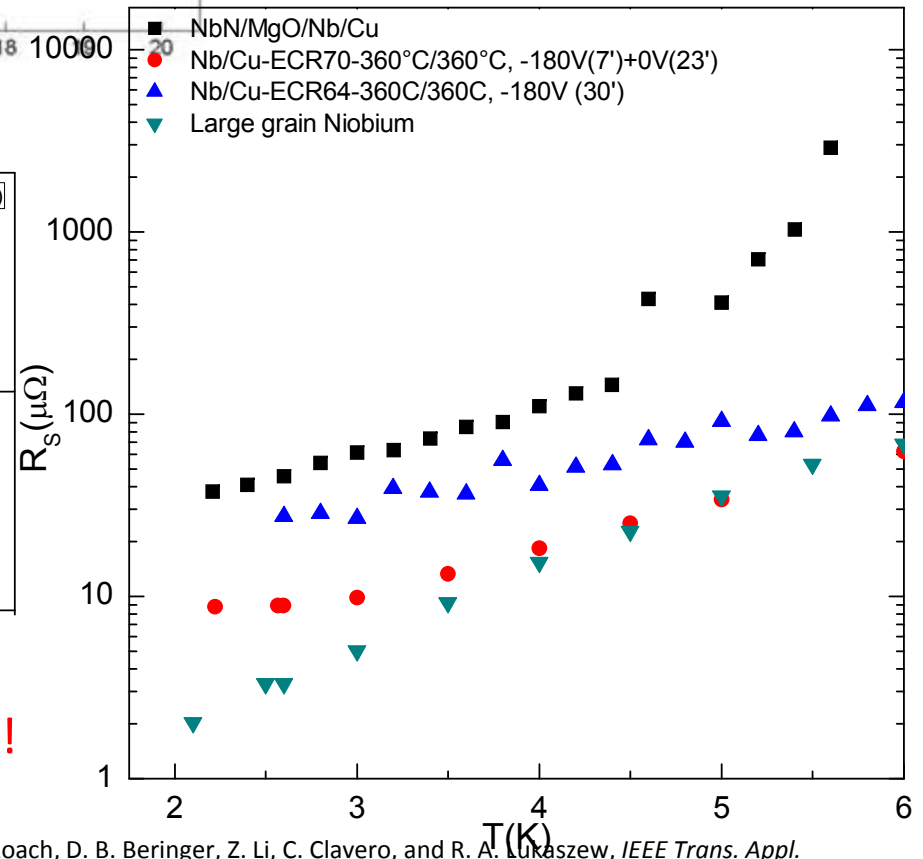
NbN tri-layer $T_c \sim 13.4$ K

Single crystal NbN films



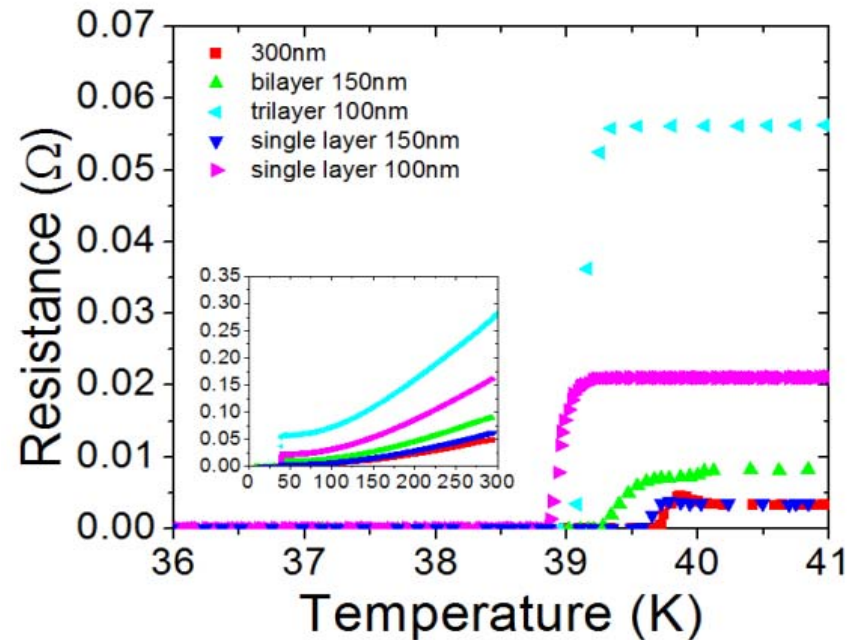
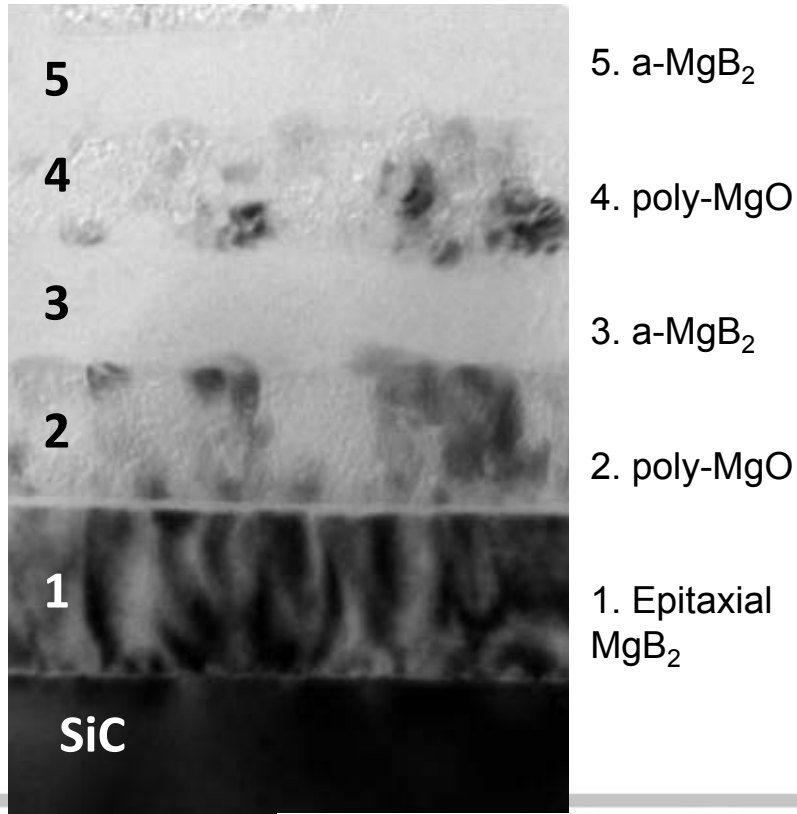
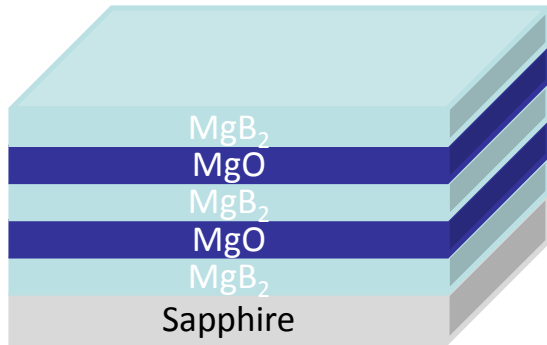
H_{c1} -NbN-based-Multilayer ~ 220 mT!

H_{c1} -bulk Nb = 170 mT



"Magnetic Shielding Larger than the Lower Critical Field of Niobium in Multilayers" W. M. Roach, D. B. Beringer, Z. Li, C. Clavero, and R. A. Lukaszew, *IEEE Trans. Appl. Supercond.* **23**, 8600203 (2013).

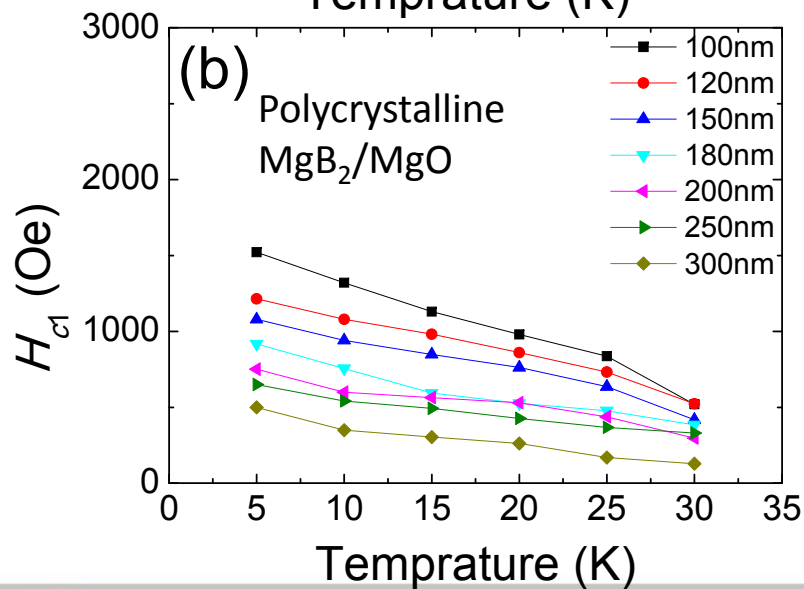
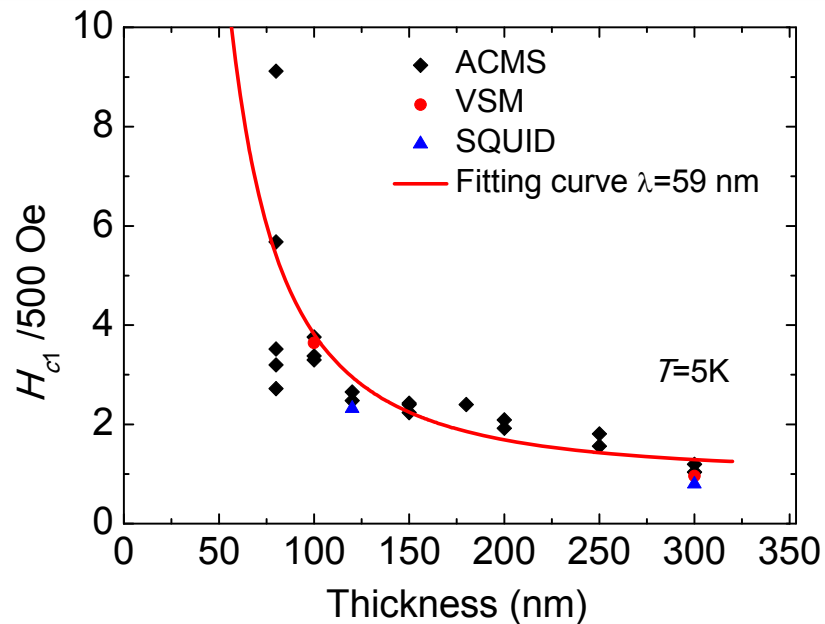
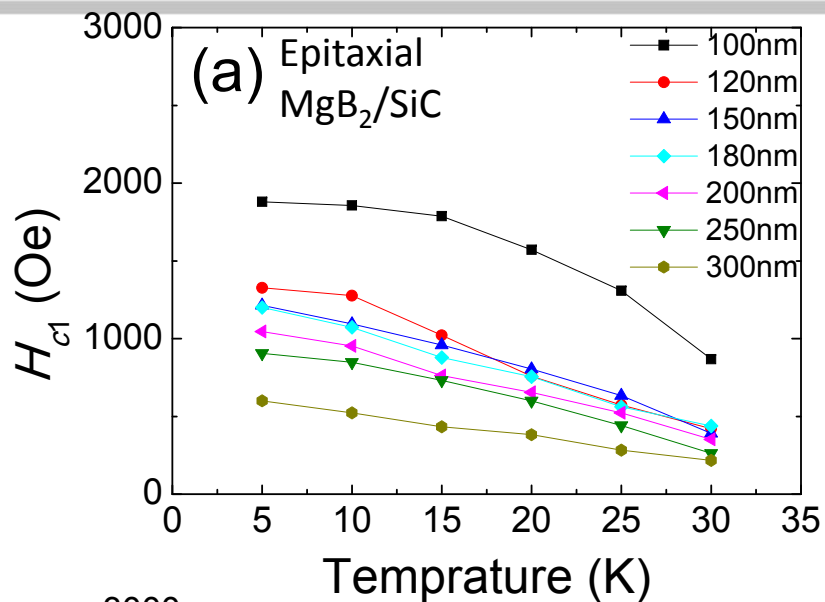
MgB₂-MgO Multilayer Films



Alternating MgB₂-insulator structures have been fabricated on sapphire substrate. Sputtering MgO are used as insulating layer.

Top MgB₂ layers amorphous.

Epitaxial and Polycrystalline Films: H_{c1} vs Thickness



$$\frac{H_{c1}}{H_{c1b}} = \left\{ 1 + \int_0^\infty \frac{(\tanh \sqrt{k^2 + (d/\lambda)^2} - 1) dk}{(\ln \kappa + 0.5) \sqrt{k^2 + (d/\lambda)^2}} \right\} / \left(1 - \operatorname{sech} \frac{d}{2\lambda} \right)$$

Epitaxial and polycrystalline MgB_2 films both show increase in $H_{c1}(0)$ with decreasing film thickness.

References

- Experimental evidences of the enhancement of the parallel H_{c1} in thin films
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- Increasing the high-field performance and reduction of R_s by a NbN overlayer
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CONCLUDING REMARKS

- ✓ Progress in energetic condensation deposition techniques opens the way to high quality Nb/Cu films and the reduction of the Q-slope of Nb/Cu cavities may be in sight.
- ✓ Over the years, some attempts have been made to study alternative materials to Nb for applications to SRF cavities.
- ✓ Most of the sample/cavities using alternative materials have been produced by reactive magnetron sputtering or thermal diffusion.
- ✓ Use of Energetic Condensation Techniques like Vacuum Arc Deposition, ECR, or ALD? (production of very dense films with nm-scale roughness...Some trials with Vacuum Arc @ INFN-Rome, non conclusive)

CONCLUDING REMARKS

- ✓ The multilayer approach opens the door to further potential improvement for SRF cavities, taking benefits from the advantages of higher T_c superconductors without the penalty of lower field onset for vortex penetration and increase the accelerating gradient and Q.
- ✓ Strong reduction of the BCS resistance for superconducting layers with higher Δ (Nb_3Sn , MgB_2 , BKBO, NbN ... s-wave, fully gapped superconductors)
- ✓ New non-magnetic members of the oxypnictide family with $15K < T_c < 55K$ or $BaO_{0.6}K_{0.4}BiO_3$ with $T_c \approx 30K$: a possibility to greatly increase Q if TF coating can be developed.

CONCLUDING REMARKS cont.

- ✓ Possibility to move from 2K to 4.2K: huge cost saving on refrigeration in LINACS
- ✓ Higher $-T_c$ s-wave materials are often multi-gap superconductors, which do not always have better SRF performance despite higher H_c and T_c
- ✓ First experimental evidence of field enhancement with NbN/MgO/.../Nb, NbTiN/AlN/Nb/Cu samples with SQUID measurement with $H \parallel$ sample plane. Field penetration delayed for NbN , NbTiN layered samples compared to Nb sample.

CONCLUDING REMARKS cont.

- ✓ Many long-standing problems of condensed matter physics and non-equilibrium superconductivity will have to be addressed to understand nonlinear surface resistance under strong rf fields
- ✓ Multi-parameter materials optimization is required to unravel the full SRF performance potential.
- ✓ The effort for new materials research for SRF cavities application has been very limited so far...but interest has been regained. There is still a lot of work ahead!